Black Phosphorus—Monolayer MoS$_2$ van der Waals Heterojunction p–n Diode

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**ABSTRACT** Phosphorene, a elemental 2D material, which is the monolayer of black phosphorus, has been mechanically exfoliated recently. In its bulk form, black phosphorus shows high carrier mobility ($\sim 10,000$ cm$^2$/V·s) and a $\sim 0.3$ eV direct band gap. Well-behaved p-type field-effect transistors with mobilities of up to $1000$ cm$^2$/V·s, as well as phototransistors, have been demonstrated on few-layer black phosphorus, showing its promise for electronics and optoelectronics applications due to its high hole mobility and thickness-dependent direct band gap. However, p–n junctions, the basic building blocks of modern electronic and optoelectronic devices, have not yet been realized based on black phosphorus. In this paper, we demonstrate a gate-tunable p–n diode based on a p-type black phosphorus/n-type monolayer MoS$_2$ van der Waals p–n heterojunction. Upon illumination, these ultrathin p–n diodes show a maximum photodetection responsivity of 418 mA/W at the wavelength of 633 nm and photovoltaic energy conversion with an external quantum efficiency of 0.3%. These p–n diodes show promise for broad-band photodetection and solar energy harvesting.

**KEYWORDS:** black phosphorus · phosphorene · MoS$_2$ · p–n diode · van der Waals heterojunction · photodetection · solar cell

The successful isolation of graphene from graphite has led to its extensive study in physics, materials, and nanoengineering due to its extraordinary electrical and mechanical properties.$^1$–$^4$ However, a lack of a band gap limits its potential for electronic device applications and has inspired the exploration of other 2D-layered materials.$^5$–$^7$ Among them, transition metal dichalcogenides (TMDCs), such as MoS$_2$, are the most studied materials.$^8$–$^{11}$ Recently, phosphorene, the monolayer form of black phosphorus, has been successfully isolated.$^{12}$ Analogous to graphite and graphene, black phosphorus is a stack of phosphorene monolayers, bound together by van der Waals interactions.$^{12,13}$ Bulk black phosphorus shows a $\sim 0.3$ eV direct band gap and a mobility of up to $\sim 10,000$ cm$^2$/V·s.$^{14}$–$^{17}$ Its band gap increases as its thickness decreases and is predicted to have a $> 1$ eV direct band gap in its monolayer form.$^{12,13}$ Well-behaved p-type field-effect transistors with mobilities of up to $1000$ cm$^2$/V·s, as well as inverters, have been demonstrated on few-layer black phosphorus.$^{12,13,18}$ Based on its direct band gap, few-layer black phosphorus phototransistors have been demonstrated with a responsivity of 4.8 mA/W.$^{19}$ These results indicate that black phosphorus is a promising candidate for both high-performance electronics and optoelectronics applications due to its ultrathin 2D nature, high hole mobility, and narrower direct band gap compared to most of TMDCs.

The p–n junctions are the basic building blocks of modern semiconductor devices, including diodes, bipolar transistors, photodiodes, light-emitting diodes, and solar cells. In the conventional p–n homojunction, the p- and n-type regions are formed by chemically doping a bulk semiconductor, creating a graded junction region. The p–n heterojunctions can be realized by epitaxially growing an n-type semiconductor on another p-type semiconductor or vice versa and can form abrupt p–n junctions. Relying on the van der Waals interactions, atomically...
sharp 2D heterostructures can be achieved without the problems caused by lattice mismatches between materials. Moreover, the ultrathin nature of 2D heterostructures allows for electrical modulation of the band structure in the vertical direction. This creates a new avenue for the realization of novel 2D electronic and optoelectronic devices.21–33

Here, we report an electrically tunable black phosphorus–monolayer MoS2 van der Waals heterojunction p–n diode. To the best of our knowledge, this is the first 2D heterostructure demonstrated using black phosphorus. The p-type black phosphorus and n-type monolayer MoS2 form an atomically sharp type II heterointerface through van der Waals interactions.13,34 The p–n diode exhibits gate-tunable current-rectifying characteristics. Upon illumination, the p–n diode can be used as photodetector with a maximum responsivity of 418 mA/W, which is nearly 100 times higher than that reported for few-layer black phosphorus phototransistors and 26 times higher than that recently reported for WSe2 p–n diodes.19,35,36 The photovoltaic power generation in the diode reaches a peak external quantum efficiency of ∼0.3%. Furthermore, photocurrent mapping confirms that the current is generated throughout the entire overlapped p–n junction region, showing its feasibility for use in large area solar cells and photodetectors.

RESULTS AND DISCUSSION

In the experiments, monolayer MoS2 was synthesized by chemical vapor deposition (CVD) on a 285 nm SiO2/p+-doped Si substrate using MoO3 and sulfur. More information about the synthesis method can be found in our previous papers.37,38 Few-layer black phosphorus was mechanically exfoliated using adhesive tape from bulk material onto a SiO2 substrate on which CVD monolayer MoS2 had been synthesized. Due to van der Waals interactions, 2D black phosphorus–monolayer MoS2 heterojunctions can be formed at the overlapped regions. Electron-beam lithography was used to define the contact patterns. Finally, 20/60 nm of Ni/Au were deposited as metal contacts. Figure 1a,b shows the schematics of the device structure and the electrical connections. A voltage (Vd) was applied across the diode, and a voltage (Vg) was applied to the p+-doped silicon as a back gate, which can electrically modulate the device. The optical image of the fabricated device is shown in Figure 1c, and all the following results were obtained from this particular device unless otherwise specified. To achieve better current collection and reduce series resistance, the MoS2 electrode is designed to surround the black phosphorus flake.

The thicknesses of the few-layer black phosphorus flake and monolayer MoS2 are ∼11 and ∼0.9 nm, respectively, as determined from the atomic force microscopy measurements shown in Figure 2a. The monolayer nature of the MoS2 can be further confirmed by the peak position (∼1.8 eV) in its photoluminescence spectrum, shown in Figure 2b.39 The Raman spectra from the black phosphorus, monolayer MoS2, and the heterojunction regions are presented in Figure 2c. The observed Raman-active modes of black phosphorus and MoS2 are consistent with previously reported data12,37 The peaks of both MoS2 and black phosphorus can be observed in the overlapped region with very little shift from their positions in the non-overlapped regions, indicating good quality of thin
films in the junction region after exfoliation and device fabrication.

Next, we studied the electrical characteristics of the fabricated device. In this paper, all the measurements were performed in ambient atmosphere. Figure 3a shows the gate-tunable $I-V$ characteristics of the p–n diode, and the inset (1) shows the $I-V$ curves on a semilog plot. The current-rectifying characteristics can be modulated by the back gate voltage, shown in the inset (2) of Figure 3a. The rectification ratio, defined as the ratio of the forward/reverse current, increases as the back gate voltage decreases. By using −30 V back gate voltage, a rectification ratio of $\sim 10^5$ is obtained at $V_d = -2/2$ V. Additionally, the ideal factor of the p–n diode increases as the back gate voltage increases (Supporting Information Figure S8) and achieves a minimum value of 2.7 with a back gate voltage of −30 V. These strong current-rectifying characteristics indicate that a good van der Waals p–n heterojunction formed between p-type black phosphorus and n-type MoS$_2$. The modulation effect of the back gate voltage can also be seen in the transfer curves (Figure 3b). Both forward and reverse currents can be substantially increased by increasing the back gate voltage. These results can be explained by a simplified model describing the current transport of the p–n diode (see Supporting Information for more details). For simplicity, the total resistance of the device can be roughly divided into three parts: the resistance of the p–n junction near the interface, the sheet resistances of MoS$_2$ and few-layer black phosphorus, and the contact resistances of metal/MoS$_2$ and metal/black phosphorus. The current-rectifying characteristics come from the p–n interface region, which may be modeled as many parallel p–n diodes (Figure S2a). The band alignment of MoS$_2$ and black phosphorus at the p–n junction interface can be modulated by the back gate voltage, as shown in Figure S3. This can explain why the rectification ratio increases as the back gate voltage decreases, and the ideal factor decreases as the back gate voltage decreases. Moreover, the back gate voltage can also modulate the sheet resistance and the contact resistance. As the device size is relatively large, the sheet resistances are more significant than the contact resistances. Because the sheet resistance of CVD monolayer MoS$_2$ is much larger than the 11 nm black phosphorus, the overall device resistance is strongly impacted by the resistance of MoS$_2$. By increasing the back gate voltage ($V_{bg}$), both the resistance of MoS$_2$ and its contact resistance are reduced,

Figure 2. (a) Atomic force microscopy image of part of the junction region. It shows that the black phosphorus and MoS$_2$ films are $\sim 11$ and $\sim 0.9$ nm thick, respectively. Scale bar, 2 μm. (b) Photoluminescence spectrum of the MoS$_2$ thin film. It shows a strong peak close to 1.8 eV, confirming that the MoS$_2$ is a monolayer. (c) Raman spectra of the MoS$_2$, black phosphorus, and the overlapped regions. Both of the black phosphorus and MoS$_2$ Raman peaks can be observed in the overlapped region.
which can efficiently boost the total current under both forward and reverse bias (see the Supporting Information for more details).

As good current-rectifying characteristics were achieved, the optoelectronic characteristics of these 2D p–n diodes were then explored. A 633 nm He–Ne laser was used to illuminate the device. The spot size of the laser was controlled to be much smaller than the junction area in order to exclude the photoresponse from the non-overlapped region. Figure 4a shows the I–V curves of the p–n diode under various incident laser power, and the inset shows the details of the negative Vg bias region. The photocurrent Ipht, is defined as \( I_{\text{ph}} = I_{\text{illumination}} - I_{\text{dark}} \), where \( I_{\text{illumination}} \) and \( I_{\text{dark}} \) are the \( I_d \) with and without laser illumination. Under reverse bias, the photocurrent has a strong dependence on the incident power (Figure 4b). By increasing the back gate voltage, the photocurrent is substantially increased. This can be attributed to the reduction of the MoS2 sheet resistance and the MoS2/metal contact resistance (see Supporting Information for more information). On the other hand, by applying a strong negative back gate voltage (−40 V) to modulate the band alignment of MoS2 and black phosphorus (Figure S3), the ratio of the illumination current over the dark current \( (I_{\text{illumination}}/I_{\text{dark}}) \) can be increased to \( 3 \times 10^4 \), though the illumination current decreases due to the increasing of sheet and contact resistance (see Supporting Information for more details). This is \( \sim 100 \) times larger than that of a MoS2 photodetector due to a better suppression of \( I_{\text{dark}} \) by the reversed bias p–n diode. Compared with the phototransistor, it is beneficial for the device to detect signal from noise.41

With these results, the photodetection responsivity, \( R \), defined as \( I_{\text{ph}}/P_{\text{laser}} \), where \( P_{\text{laser}} \) is the incident power of the laser, can be calculated for both forward (as photoconductor) and reverse bias (as photodiode) as shown in Figure 4c. The responsivity of both the forward and the reverse bias regions decreases as the laser power increases, which is due to the saturation of electron–hole pair generation at higher incident light intensity and the increased surface recombination in either the junction region or the underlying SiO2.41 Therefore, to further improve the responsivity, it is necessary to improve the film quality of the CVD MoS2 and black phosphorus, as well as the dielectric interface quality. In this device, the responsivity reaches 1.27 A/W and 11 mA/W for forward and reverse bias, respectively. The maximum responsivity in our devices is up to 3.54 A/W (at \( V_d = +2 \) V) and 418 mA/W (at \( V_d = -2 \) V) under 1 μW laser power (Figure S5). As a photodiode, this is nearly 100 times higher than that of the recently reported black phosphorus phototransistor and 4.8 times higher than that of the carbon nanotube–MoS2 p–n diode with a much smaller \( V_{oc} \).19,32 By using this electrically tunable 2D p–n diode, it is possible to realize sensitive and broad-band photodetection due to the high mobility and relatively small direct band gap of few-layer black phosphorus.

In addition to the photodetection, this p–n diode is also usable for photovoltaic energy conversion. Under laser illumination, the short circuit current \( (I_{\text{sc}}) \), which is the current through the device with 0 V bias, is shown in Figure 5a. The open circuit voltage \( (V_{oc}) \) can be obtained by examining the \( V_g \) axis intercepts in the plot. As seen in Figure 5c,d, both the \( I_{\text{sc}} \) and \( V_{oc} \) increase as the laser power increases. However, the back gate modulates the \( I_{\text{sc}} \) and \( V_{oc} \) in different ways. Increasing the back gate voltage boosts the \( I_{\text{sc}} \) while reducing the \( V_{oc} \). This can be attributed to the opposite modulation effect of band alignment of MoS2/black phosphorus at the junction interface and the sheet and contact resistance of MoS2 (see Supporting Information for details). With these data, the power generated by the p–n diode can be calculated by \( P_d = I_{\text{sc}}V_{oc} \), as shown in Figure 5b. Then, the fill factor, which is defined as \( FF = P_d/(I_{\text{sc}}V_{oc}) \), can be obtained. The maximum FF is \( \sim 0.5 \) in this 2D p–n diode. Furthermore, the external quantum efficiency (EQE) can be calculated by \( EQE = \frac{I_{\text{ph}}}{P_{\text{laser}}(hc/e\lambda)} \), where \( h \) is the Planck constant, \( e \) is the electron charge, \( c \) is the velocity of light, and \( \lambda \) is the
wavelength of the incident light. The peak EQE in this device is 0.3%. To the best of our knowledge, this is the first time that efficient photovoltaic energy conversion has been demonstrated using black phosphorus. By further reducing the thickness of the black phosphorus to bilayer, it is possible to essentially enhance the EQE to up to 18%, according to theoretical predictions. Moreover, by stacking the 2D p–n diode vertically, a stacked solar cell structure can be realized. Since the direct band gap of black phosphorus changes with its thicknesses, the efficiency may be further improved by tuning its thickness to better utilize the photon energies of different wavelengths of the solar spectrum.

To further understand the photoresponse process, photocurrent mapping was performed in order to determine the spatial distribution of the photocurrent generation. A device with a larger junction region was chosen to get more spatial information, as shown in Figure 6a. The thickness of the black phosphorus flake is ∼22 nm (Figure S6). A 633 nm He–Ne laser was used with ∼8 μW power and ∼1 μm spot size. The current was measured under 0 V bias. As seen in the Figure 6b, the current is generated throughout the junction region, which indicates the charge separation mainly in the heterojunction region. However, the photoreponse across the junction is not uniform. This can be attributed to the non-uniform contacts between the two flakes. Moreover, the part of the junction region that is near the metal contact on the MoS2 exhibits stronger photocurrent. This is related to the lateral sheet resistance of MoS2 and black phosphorus. Under illumination, the charges are separated near the vertical interface p–n diode region. For simplicity, we treat the interface p–n diode (see Figure S4) at a certain position as a voltage source under illumination. Here we assume the value of this voltage source is uniform across the junction if different positions are under laser illumination. The photocurrent flows though MoS2 and black phosphorus to reach the metal contact (Figure S4). Because the sheet resistance of CVD monolayer MoS2 is much larger than the black phosphorus, the resistance of the MoS2 makes a much stronger impact. The part of the junction region closer to the metal contact (B in Figure S4, compared with A) on MoS2 shows a lower MoS2 sheet resistance, resulting in the stronger photocurrent near the contact of MoS2 (see Supporting Information for details). To get a more uniform spatial photoresponse, a transparent contact material, such as graphene, could be used as top electrodes. This can increase the contact area and yield better carrier collection. Furthermore, graphene can also reduce the contact resistance as an interlayer between the metals and the TMDCs, further improving the device efficiency. This scheme makes it possible for the black

Figure 4. (a) I–V characteristics of the p–n diode under various incident laser powers. The inset shows the details in the reverse bias region. (b) Photocurrent as a function of incident laser power. Increasing the back gate voltage can increase the photocurrent. The inset shows the ratio of illumination/dark. (c) Photodetection responsivity (R) calculated as a function of incident power. The responsivity decreases as the power increases.
phosphorus/MoS2 heterojunction to be used as large area transparent and flexible solar cells, as well as photodetectors.

CONCLUSION

In conclusion, for the first time, we demonstrate a gate-tunable black phosphorus–/p-MoS2 monolayer MoS2 van der Waals heterojunction p–n diode. These 2D p–n diodes exhibit strong gate-tunable current-rectifying I–V characteristics, which is investigated using a simplified device model. As a photodetector, the p–n diode shows a photodetection responsivity of 418 mA/W, which is nearly 100 times higher than the reported black phosphorus phototransistor. Due to the small band gap and high mobility of few-layer black phosphorus, it is very suitable for broad-band and sensitive photodetection. Our p–n diodes also show photovoltaic power generation.
with a peak external quantum efficiency of 0.3%, which can be further improved by reducing the thickness of black phosphorus. Furthermore, photocurrent mapping is performed to study the spatial distribution of photosresponse, which reveals the importance of large area transparent electrodes for 2D p–n diodes to achieve large area optoelectronics applications.

METHODS
First, monolayer MoS$_2$ was synthesized using chemical vapor deposition method on heavily doped silicon wafer capping with 285 nm SiO$_2$. The black phosphorus was mechanically exfoliated from the bulk black phosphorus onto the wafer. The black phosphorus—monolayer MoS$_2$ heterojunction was achieved through the van der Waals interactions. Then the standard e-beam lithography process was used to define contact patterns and pads. A 20/60 nm Ni/Au was deposited as metal contacts. The electrical measurements were performed using Keithley 4200 semiconductor characterization system. All the electrical and optical measurements were performed in the ambient atmosphere.

Conflict of Interest: The authors declare no competing financial interest.

Acknowledgment. We thank Prof. Jong-Hyun Choi, Haorong Chen, and Te-Wei Wen from Purdue University for help with AFM characterization. This material is based upon work partly supported by NSF under Grant CMMI-1120577 and SRC under Task 2396.

Supporting Information Available: I–V characteristics of various devices, analysis of current transport in these p–n diodes, photodetection responsivity of a p–n diode, and thickness data of the few-layer black phosphorus flake for photocurrent mapping measurements. This material is available free of charge via the Internet at http://pubs.acs.org.

REFERENCES AND NOTES


