Simulating Buoyancy-Driven Airflow in Buildings by Coarse-Grid Fast Fluid Dynamics

Mingang Jin[1], Wei Liu[2,1], Qingyan Chen[2,1*]

[1]School of Mechanical Engineering, Purdue University, West Lafayette, IN 47907, USA
[2]School of Environmental Science and Engineering, Tianjin University, Tianjin 300072, China

[*]Tel. (765)496-7562, Fax (765)494-0539, Email: yanchen@purdue.edu

Abstract

Fast fluid dynamics (FFD) is an intermediate model between multi-zone models and computational fluid dynamics (CFD) models for indoor airflow simulations. The use of coarse grids is preferred with FFD in order to increase computing speed. However, by using a very large mesh cell to represent a heat source that could have a much smaller physical size than the cell, coarse-grid FFD would under-predict the thermal plume and thermal stratification. This investigation integrated a thermal plume model into coarse-grid FFD. The integration first used the plume model to calculate a source for the momentum equations and then corrected the temperature at the plume cell. The integration enabled coarse-grid FFD to correctly predict the plumes. When applied to displacement ventilation, coarse-grid FFD with the plume model can accurately predict the mean air temperature stratification in rooms as compared with experimental data from the literature. The improved model has also been used to calculate the ventilation rate for buoyancy-driven natural ventilation. The calculated ventilation rates agree well with the experimental data or predictions by CFD and analytical models. Coarse-grid FFD with the plume model used only a small fraction of the computing time required by fine-grid FFD, while the associated errors for the two grid sizes were comparable.

Keywords: Heat sources, plume model, computational fluid dynamics, experimental validation, analytical model

Research highlights

- This study integrated a thermal plume model into coarse-grid FFD.
- The integrated model was tested with buoyancy-driven airflow in buildings.
- The plume model can improve representing the heat source with very coarse cell.
- The integrated model used less computing time but maintained reasonable accuracy.

1. Introduction

Whole-building airflow simulations are required in applications such as natural ventilation design, coupled building airflow and energy simulation, smoke control, and air quality diagnosis in a building. These simulations generally use multi-zone models [1]. However, the models can provide only very limited airflow information because of the assumption that a room within a building can be treated as a single homogeneous node. Computational Fluid Dynamics (CFD) models, on the other hand, can perform detailed airflow simulations, but the use of CFD for whole-building airflow simulations is too computationally expensive [2]. Between CFD models and multi-zone models, researchers have also developed intermediate models for whole-building airflow simulations. Zonal models [3] are typical intermediate models that can achieve a balance between reduced computing costs and the level of detail required in airflow simulations. Additionally, by using very coarse grids, coarse-grid CFD models [4] can provide more detailed airflow simulations at a competitive computing speed with respect to zonal models, and they are expected to replace zonal models in the future [2]. Fast Fluid dynamics (FFD), a recently developed intermediate model that can provide reliable simulations of indoor airflows at a speed that is about 15 times faster than CFD models, currently has great potential for performing whole-building airflow simulations [5, 6]. Because FFD is also a grid-based model, reducing the grid number can further enhance the computing speed of FFD simulations. Coarse-grid FFD would be an ideal tool for performing whole-building airflow simulations at a greatly reduced computing cost.

Although a coarse grid could significantly reduce the computing time of FFD simulations, it may cause problems in the representation of the boundary conditions encountered in building airflow simulations. For example, many heat sources in buildings are of small physical size, such as computers, desk lamps, occupants, etc. Using a very large mesh cell to represent a small-sized heat source would result in the prediction of lower energy intensity in the cell and smaller buoyancy forces from the heat source. Coarse-grid FFD would thus tend to under-predict the plume flow generated by the heat source and would not accurately predict buoyancy-driven ventilation and room air temperature distribution. Because buoyancy-driven ventilation is a major feature of high-performance building systems, such as displacement ventilation and buoyancy-driven natural ventilation systems, correct prediction of buoyancy-driven ventilation and room air temperature distribution is essential with coarse-grid FFD.

Thus it is necessary to improve the representation of small heat sources with large cells. In CFD models, simulations are normally performed on fine grids, which allows CFD models to avoid the aforementioned problem. Instead, to reduce the complexities of representing heat sources in simulations, it is still necssary for CFD models to apply simple heat source geometries or use replacement boundary conditions [7]. As intermediate models, zonal models also have the same problem of reprezenting heat sources as coarse-grid FFD does. Thus the approaches applied in zonal models to model thermal plumes could also be a potential solution for coarse-grid FFD.

71 Extensive research has been conducted into the characteristics of thermal plumes [8]. Morton et
72 al. [9] proposed a theoretical model to describe the physics of thermal plumes, and this model
73 has been adapted for studying a wide variety of thermal plumes. Kofoed [10] experimentally
74 studied thermal plumes generated by indoor heat sources in ventilated rooms and proposed a
75 model coefficient to account for the influence of enclosing walls. Trzeciakiewicz [11]
76 experimentally investigated the characteristics of thermal plumes in response to objects of
77 varying shape, such as computers, desk lamps, and light bulbs. The investigation revealed that
78 the experimentally determined model of a plume above a point heat source could be used to
79 characterize the thermal plumes in displacement ventilation. Zukowska et al. [12] investigated
80 the characteristics of the thermal plume generated by a sitting person using four different
81 geometries and found that a rectangular box could correctly simulate the enthalpy flux and
82 buoyancy flux generated by the person. Craven and Settles [13] performed a computational and
83 experimental investigation to characterize the thermal plume from a person and concluded that
84 the room temperature stratification had a significant effect on plume behavior. The
85 aforementioned research into the characteristics of indoor thermal plumes has provided a wealth
86 of information.

87 On the basis of these studies of thermal plume physics and analytical plume models, simple
88 models have been developed to quantify ventilation and temperature distributions in buildings
89 [14, 15]. In addition, plume models have been used to improve the performance of other models
90 for simulating buoyancy-driven airflows in buildings. Inard et al. [16] integrated a wall thermal
91 plume model into a zonal model to improve the latter's performance in simulating the
92 temperature distribution in a room. Musy et al. [17] integrated a plume model with a zonal model
93 to obtain a better simulation of natural convection in a room with a radiative-convective heater.
94 Stewart and Ren [18] used a plume model to improve the simulation accuracy of a zonal model
95 for studying the airflow rising from a cooking plate. It has been shown that plume models
96 effectively enhance the performance of zonal models in simulating buoyancy-driven airflows in
97 buildings.

98 The integration of plume models with zonal models suggests that plume models could also be
99 integrated with FFD for improving the performance of coarse-grid FFD simulations for room
100 airflows driven by heat sources. This study therefore developed a method of implementing a
101 plume model in FFD when the mesh cell is much larger than the heat source. The proposed
102 integration method was also tested and evaluated.

103

2. Research Method

2.1 Fast fluid dynamics

106 FFD simulates an airflow by numerically solving a set of partial differential equations
107 representing the transport phenomena in the airflow, Eq. (1)-(3), which are derived on the basis

108   of the conservation of mass, momentum (Navier-Stokes equations), and scalar transport
109   quantities (such as energy and species), respectively.

$$\frac{\partial U_i}{\partial x_i} = 0, \tag{1}$$

$$\frac{\partial U_i}{\partial t} + U_j \frac{\partial U_i}{\partial x_j} = -\frac{1}{\rho}\frac{\partial p}{\partial x_i} + v\frac{\partial^2 U_i}{\partial x_j \partial x_j} + \frac{1}{\rho}F_i, \tag{2}$$

$$\frac{\partial \phi}{\partial t} + U_j \frac{\partial \phi}{\partial x_j} = \Gamma \frac{\partial^2 \phi}{\partial x_j \partial x_j} + S, \tag{3}$$

113   where i or j = 1, 2, 3; U_i is the ith component of the velocity vector, x_i the ith direction of
114   coordinate, t time, p pressure, ρ density, v the kinetic viscosity, F_i the ith component of the body
115   forces, ϕ the scalar variables, Γ the transport coefficient for ϕ, and S the source term. In each
116   time step, FFD solves this set of transport equations sequentially. To enhance computational
117   efficiency, a time-splitting scheme [19] was applied to solve the transport equations. For
118   example, FFD splits the scalar transport equation (3) into an advection equation (4) and a
119   diffusion equation (5),

$$\frac{\phi^{(1)} - \phi^n}{\Delta t} = -U_j \frac{\partial \phi^n}{\partial x_j}, \tag{4}$$

$$\frac{\phi^{n+1} - \phi^{(1)}}{\Delta t} = \Gamma \frac{\partial^2 \phi^{n+1}}{\partial x_j^2} + S, \tag{5}$$

122   where ϕ^n and ϕ^{n+1} represent the variable at the current and next time steps, respectively, and $\phi^{(1)}$
123   represents the intermediate variables solved by the advection equation. The advection equation
124   (4) is first solved with the conservative semi-Lagrangian scheme [20] to obtain the intermediate
125   value $\phi^{(1)}$, and then FFD is implicitly solved using the diffusion equation (5) to update the scalar
126   distribution at the next time step. To effectively resolve the coupling between the momentum
127   equations and the continuity equation, a pressure projection [21] is performed force the velocity
128   field to satisfy continuity.

129   *2.2 Integration with plume model*

130   As a result of natural convection, air surrounded by a heat source in a room can form a buoyancy
131   plume. As the plume rises, it induces the surround air into the plume flow, and flow carries the
132   heat from the heat source to the upper part of the room. To describe the features of buoyancy-
133   driven airflows, the plume flow rate (V_p) and excess temperature (ΔT_p) in the plume region are
134   usually the two most important parameters [22], and they vary with the heat generation rate, heat
135   source geometry, heat source location, etc. When the mesh cell size used in coarse-grid FFD is

136 much larger than the physical size of a heat source, the two parameters predicted by FFD may
137 not be accurate. In order to improve the performance of coarse-grid FFD in predicting the
138 thermal plume, this study used an analytical plume model with empirical coefficients to calculate
139 the plume flow rate and the excess temperature in the plume region. The two calculated
140 parameters were then integrated into FFD in order to represent the heat source.

141 By assuming axisymmetric temperature and velocity profiles and by using the similarity rule and
142 constant entrainment ratio in the fully developed plume region, Eqs. (6) and (7) were developed
143 to describe the excess temperature and plume flow rate, respectively, for a thermal plume
144 generated by a point heat source [9].

$$\Delta T_p = k_T P_c^{2/3} (z + z_v)^{-5/3} \tag{6}$$

146
$$V_p = k_v P_c^{1/3} (z + z_v)^{5/3} \tag{7}$$

147 where P_c is the convective heat generation rate from the heat source, z the distance between a
148 cross-section in the plume and the top surface of the heat source, z_v the distance between the
149 virtual start point of the plume and the top surface of the heat source, k_v the flow rate coefficient,
150 and k_T the excess temperature coefficient, as shown in Figure 1. Although the heat sources in
151 buildings are generally not point sources and the generated themal plumes are not always fully
152 developed, researches found that it was possible to apply Eqs. (6) and (7) to simulate the plumes
153 in a room with buoyancy-driven ventilation system [11]. To characterize themal plumes in
154 buildings, empirical coefficients based on the investigation of heat sources in buildings were
155 applied in Eqs. (6) and (7), where k_v is 0.0051 $m^{4/3}/(W^{1/3} s)$ [10] and k_T is 0.4 K $m^{5/3}/W^{2/3}$ [22].
156 Noted that the value of these two coeffients are only applicable for free plume generated by
157 single heat source, it is necessary to adjust the coefficients accordingly for merging thermal
158 plumes from multiple heat sources or thermal plume attached to solid wall [23].
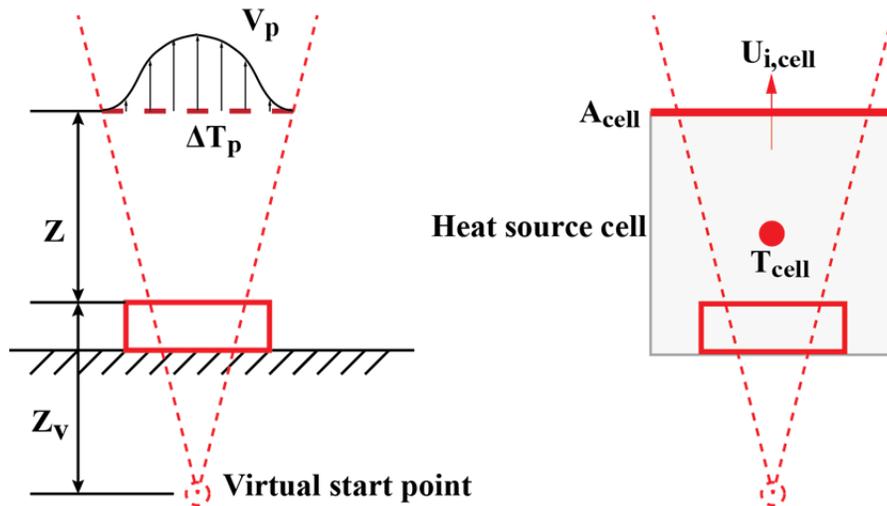
159

160 Figure 1 Schematic of the plume model and the heat source cell

5

161 The integration of the plume model with FFD was divided into two parts: integration with the
162 momentum equations and with the energy equation. Because coarse-grid FFD cannot correctly
163 predict the plume flow rate driven by the heat source, Eq. (7) was used to estimate the airflow
164 rate of the plume at the heat source cell. Dividing the flow rate by the area of the horizontal cell
165 surface A_{cell} provided an estimate of the velocity component in the plume direction $U_{i,\mathrm{plume}}$ for
166 the heat source cell,

$$U_{i,\mathrm{plume}} = V_p / A_{cell} . \tag{8}$$

168 The plume velocity $U_{i,\mathrm{plume}}$ was compared with the vertical velocity component $U_{i,cell}^n$ of the heat
169 source cell at the previous time step in order to determine the amount of correction required for
170 adjusting the velocity at the heat source cell. On the basis of the difference between $U_{i,\mathrm{plume}}$ and
171 $U_{i,cell}^n$, the momentum source or sink was calculated and was incorporated into the momentum
172 equations as a source term, as shown in Equation (9),

$$F_{i,cell} = F_{i,cell} + \rho(U_{i,\mathrm{plume}} - U_{i,cell}^n)/\Delta t, \tag{9}$$

174 where $F_{i,cell}$ is the body force term in the heat source cell and Δt the time step size. When the

175 momentum source term in the cell is incorporated the plume velocity at the cell can be corrected

176 to be the same as that in the analytical model.

177 For integration with the energy equation, the plume model was used to adjust the air temperature
178 at the heat source cell because FFD underestimates the heat source temperature when a very
179 coarse grid is used. To minimize the impact of the temperature correction on the energy
180 conservation in the domain, the correction was applied in the advection process and before the
181 energy conservation correction of the semi-Lagragian scheme. Therefore, a three-step approach
182 was applied to solve the advection equation for energy transport:

$$T_{cell}^* = T^n(\vec{x}_{cell} - \Delta t \vec{U}_{cell}), \tag{10}$$

$$T_{cell}^{**} = T_{cell}^* + (\Delta T_p + T_{ambient}^n - T_{cell}^n), \tag{11}$$

$$T_{cell}^{(1)} = T_{cell}^{**} + \eta_{cell}\beta, \tag{12}$$

186 where, $\vec{x}_{cell}$ represents the coordinates of the heat source cell, $\vec{U}_{cell}$ the velocity vector at the
187 heat source cell, Δt the time step size, T^n the temperature at the previous time step, $T^{(1)}$ the
188 temperature solved by the advection equation, $T_{ambient}$ the air temperature away from the plume
189 region, η the correction weighting factor, and β the energy imbalance rate. To solve the
190 advection equation for temperature at the cell, the standard semi-Lagrangian scheme was first

191 applied to obtain intermediate temperature T^* from Eq. (10). Next, on the basis of the excess
192 temperature calculated by the plume model, a correction was conducted for the temperature at
193 the cell in order to obtain corrected intermediate temperature T^{**} by using Eq. (11). FFD then
194 applied the energy conservation correction (Eq. (12)) to obtain the cell temperature after
195 advection.

196 When the plume model has been integrated into the momentum and energy equations according
197 to the procedure described above, FFD with coarse grids can correctly predict the airflow and
198 temperature in the heat source cell. Furthermore, the simulation is stable, and the airflow is
199 conservative.

200 ## 3. Results

201 Using the integrated plume model, this study applied coarse-grid FFD to simulate building
202 airflows driven by heat sources, including displacement ventilation in a chamber with a heated
203 box, displacement ventilation in an occupied office, buoyancy-driven single-sided natural
204 ventilation, and buoyancy-driven natural ventilation in an atrium. To evaluate the performance of
205 the plume model in representing heat sources in the coarse-grid simulations, the mean vertical air
206 temperature distribution and the ventilation rate predicted by FFD with and without the plume
207 model were compared with the corresponding experimental data or analytical solution. In
208 addition, this study compared the simulation accuracy of FFD with coarse and fine grids.

209 *3.1 Displacement ventilation in a chamber with a heated box*

210 This study first applied FFD to simulate displacement ventilation with a single heat source and
211 compared the simulated results with the experimental data obtained by Li et al. [24]. Figure
212 2shows the test chamber with dimensions of 4.2 m × 3.6 m × 2.75 m inside a laboratory. Air at a
213 temperature of 18 °C and ventilation rate of 125 m³/h was supplied by a diffuser with dimensions
214 of 0.5 m × 0.45 m located on the side wall of the test chamber at floor level. An air exhaust with
215 dimensions of 0.525 m × 0.22 m was located on the front wall at ceiling level. A 300-W cubic
216 heat source with dimensions of 0.4 m × 0.3 m × 0.3 m was placed on the floor in the center of the
217 test room. Table 1 shows the measured temperatures of the chamber's interior surfaces.

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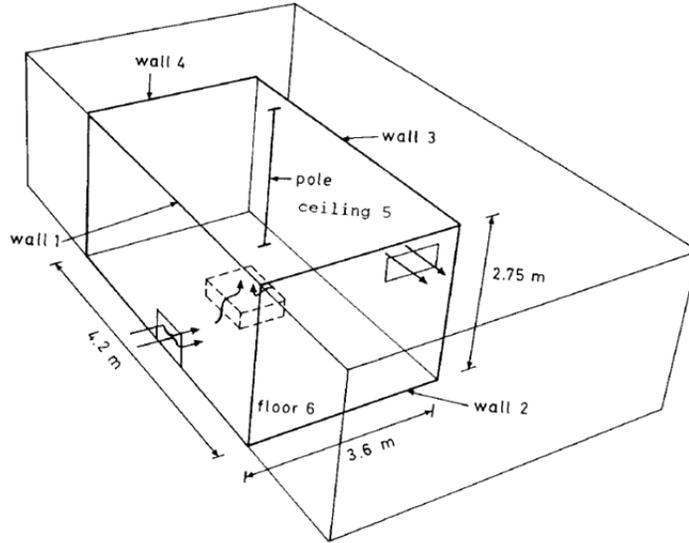
Figure 2 Geometry of the test chamber with displacement ventilation inside a laboratory [24]

221

Table 1 Temperatures of the chamber's interior surfaces [24]

Surface	Floor	Side walls at different heights					Ceiling
		0.5 m	1.0 m	1.5 m	2.0 m	2.5 m	
Temperature (K)	295.1	295	295.7	296.5	296.8	296.8	298.2

223

In the coarse-grid FFD simulations, the test chamber was represented by a total grid number of 5 × 5 × 5. Because the mesh cells were very large, the dimensions of the mesh cell containing the heat source were 0.72 m × 0.55 m × 0.84 m. The cell size was almost ten times the physical size of the heat source. This study simulated the temperature distribution in the test chamber with and without the use of the plume model to represent the heat source. In addition, fine-grid FFD with a grid number of 20 × 20 × 20 was applied to simulate the plume generation and temperature stratification in the test chamber.

Figure 3presents the temperature distribution and velocity field at the vertical mid-plane of the chamber as predicted by FFD with and without the plume model. Without the plume model, as shown in Figure 3(b), FFD predicted a smaller thermal plume. This is because the large cell used for the heat source would have a lower mean air temperature than that in the actual plume. Because of the reduced air temperature, the plume flow rate was also lower, and there was less air entrainment from the surroundings. However, with the plume model as depicted in Figure 3(a), the predicted air entrainment was much greater. Because the energy intensity was reduced in the large heat source cell, FFD without the plume model predicted an unrealistically low air temperature in the plume region, as illustrated by Figure 3b). For example, the air temperature in the plume region was even lower than that at ceiling level. The model correctly predicted the

8

241 high air temperature near the ceiling that was caused by the rising thermal plume, as shown in
242 Figure 3(a).

244   (a)   (b)
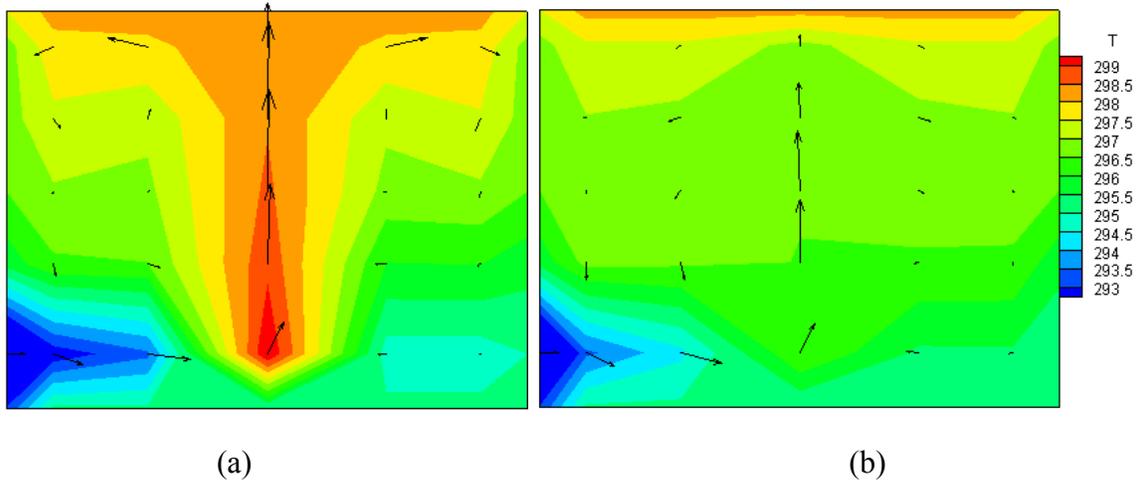
245 Figure 3 The air temperature and velocity distribution at the vertical mid-plane of the test
246 chamber as predicted by FFD (a) with the plume model and (b) without the plume model

248 This study further examined the mean air temperature distribution at different heights in the test
249 chamber as predicted by coarse-grid FFD. Figure 4 compares the temperature profile predicted
250 by course-grid FFD with those from fine-grid FFD and the data measured by Li et al. [24]. FFD
251 without the plume model under-predicted the temperature stratification in the chamber, and it
252 predicted temperatures that were too low in the upper part of the chamber and too high at floor
253 level. Coarse-grid FFD with the plume model produced a temperature profile that agrees much
254 better with the experimental data, except in the region near the floor. Because the temperature
255 gradient near the floor was high, the grid resolution used by coarse-grid FFD was too low to
256 accurately reflect the large temperature variation in this region. Fine-grid FFD provided a
257 slightly better prediction of the temperature profile near the floor. Due to the averaging of the
258 temperature at the same hight, the discrepency between the temperature profiles predicted by two
259 models was not as apprant as observed in Figure 3. But it was still obvious that the plume model
260 improved FFD's prediction of thermal plume [24]. Coarse-grid FFD with the plume model was
261 able to simulate the temperature stratification with an acceptable accuracy for engineering
262 applications.
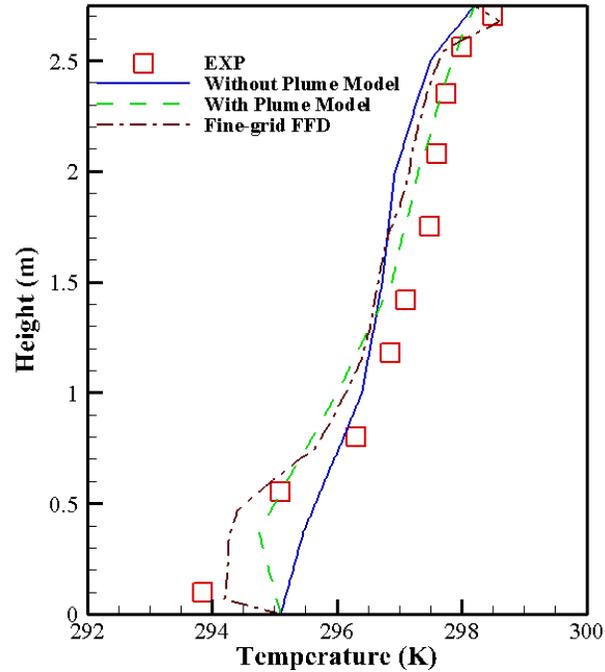
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264 Figure 4 Comparison of the mean vertical air temperature profiles in the chamber as predicted by
265 coarse-grid FFD with and without the plume model and by fine-grid FFD, as well as the
266 measured profile.

267

268 *3.2 Displacement ventilation in an occupied office*

269 To evaluate the performance of the plume model for improving the accuracy of coarse-grid FFD
270 in simulating thermal plumes, this study next investigated airflows in an occupied office space
271 with displacement ventilation. An office mock-up with dimensions of 5.16 m × 3.65 m × 2.43 m
272 was used by Yuan et al. [25] for experimental measurements, as shown in Figure 5(a). Air at a
273 temperature of 17.0 $^{\circ}$C was supplied through a displacement diffuser on the side wall at floor
274 level at a ventilation rate of 183.1 m^3/h. The air was exhausted through an outlet located in the
275 center of the ceiling. Two heated dummies with dimensions of 0.4 m × 0.4 m × 0.7 m to simulate
276 occupants, and two heated boxes with dimensions of 0.4 m × 0.4 m × 0.4 m to simulate
277 computers, were placed in the chamber as heat sources. These heat sources generated thermal
278 plumes that reached the upper part of the room. The experiment by Yuan et al. [25] measured the
279 air temperature along nine vertical poles distributed in the streamwise center plane (P1–P5) and
280 the cross-sectional center plane (P6-P9), as shown in Figure 5(b).

10

281

282   (a)   (b)
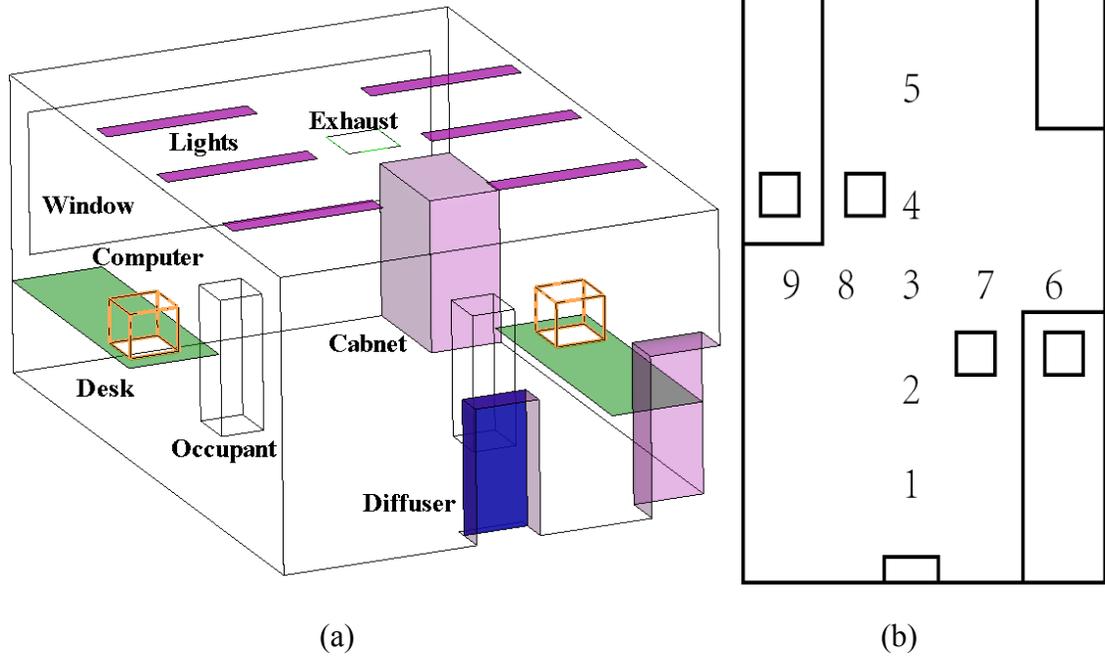
283   Figure 5 (a) Schematic of the office with displacement ventilation and (b) the measurement
284   locations

285

286   The coarse-grid FFD simulations used $5 \times 5 \times 8$ mesh cells to represent the office, as shown in
287   Figure 6. Since the mesh was very coarse, two of the cells contained heat sources, and each
288   contained an occupant and a nearby computer on a desk. Each heat source cell had dimensions of
289     $0.7 \text{ m} \times 0.75 \text{ m} \times 0.75 \text{ m}$. Once again, the simulations were performed by coarse-grid FFD with
290   and without the plume model and by fine-grid FFD with a cell number of $20 \times 20 \times 20$.
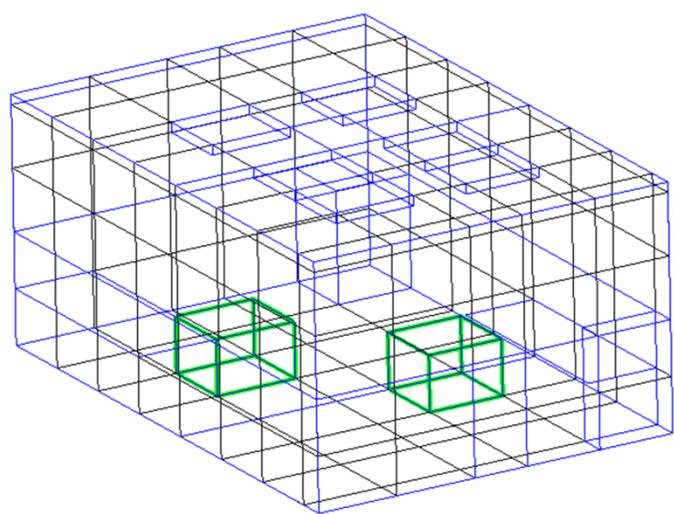
291

292   Figure 6 Coarse grid used to represent an office with two heat sources

11

293

294   Figure 7 compares the vertical air temperature distribution at five locations (PI-P5) predicted by
295   coarse-grid and fine-grid FFD with the experimental data. The results at the other measured
296   locations (P6-P9) were similar, but they are not presented here because of the limited space
297   available in this paper. Coarse-grid FFD without the plume model could not correctly predict the
298   air temperature profiles as compared with the experimental data. Because the thermal plumes
299   were artificially weakened by the use of a large cell to represent the heat sources, the amount of
300   energy transported to the upper part of the room was significantly reduced. Thus, in the upper
301   part of the office the predicted air temperature was lower than that measured experimentally.
302   However, the use of the plume model to represent heat sources improved the accuracy of the
303   coarse-grid FFD in simulating the buoyancy flows driven by the heat sources and thus provided a
304   better prediction of the vertical air temperature profiles. In addition, the simulation results proved
305   the viability of combining several adjacent heat sources into a single heat source so that the
306   plume model could be used. Interestingly, coarse-grid FFD can predict the air temperature
307   profiles with accuracy similar to that of fine-grid FFD.
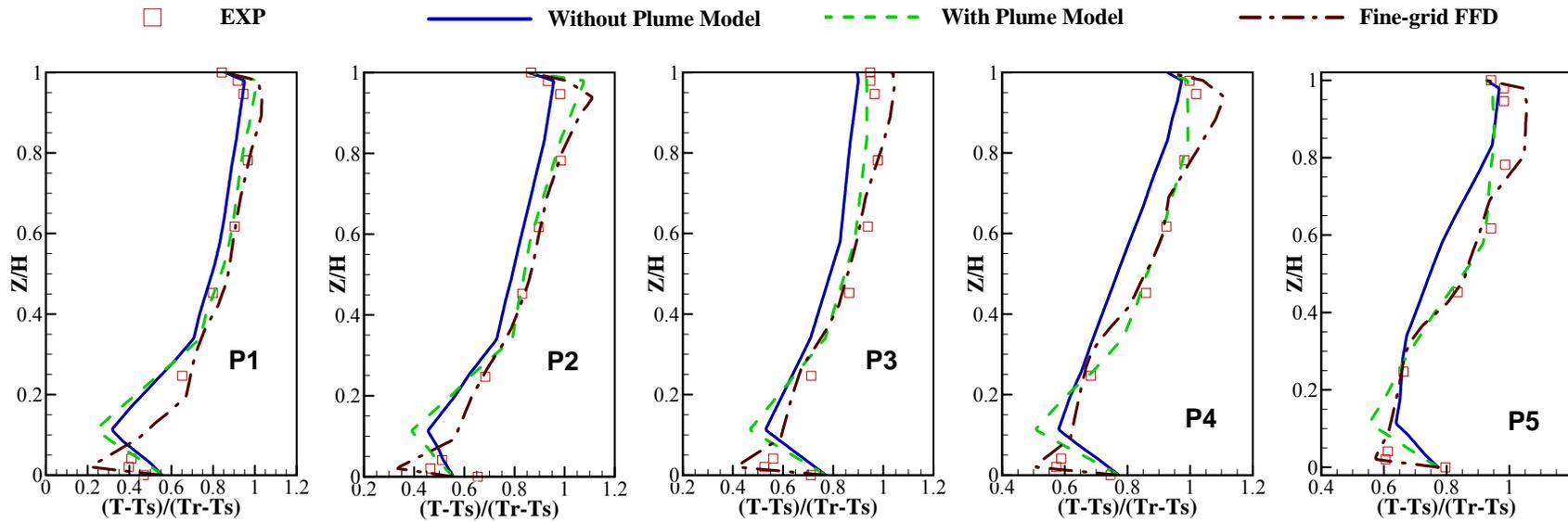
308

Figure 7 Comparison of the vertical air temperature profiles at five positions (P1-P5) as predicted by various FFD models and according to experimental data from Yuan et al. [25]

309   *3.3 Buoyancy-driven single-sided natural ventilation*

310   Following the successes in the displacement ventilation cases, this study applied the coarse-grid
311   FFD to more challenging cases, such as buoyancy-driven natural ventilation, where flow rate is
312   determined by thermal plumes. The case selected for buoyancy-driven natural ventilation in this
313   study was taken from Jiang and Chen [26], who used a test chamber with dimensions of 5.34 m
314   × 3.57 m × 2.46 m inside a laboratory to simulate the indoor environment, and the surrounding
315   laboratory space to simulate the outdoor environment, as shown in Figure 8(a). The thermal
316   plume was generated by a 1500-W baseboard heater with dimensions of 0.16 m × 0.2 m × 0.7 m
317   that was placed on the floor in close proximity to the interior wall. An open door was used to
318   simulate the single-sided opening where natural ventilation occurred. Because the walls of the
319   test chamber and environmental chamber were of very high insulation value, they were
320   considered to be adiabatic. Jiang and Chen measured the vertical air temperature profiles at five
321   different locations, as shown in Figure 8(b), as well as the ventilation rates in the test chamber.
322   Table 2 shows the enclosure surface temperatures of the laboratory as measured by Jiang and
323   Chen.

324   Table 2 Enclosure surface temperatures of the laboratory

	Ceiling	Floor	North wall	South wall	East wall	West wall
Temperature (K)	296.11	295.11	296.01	295.90	293.94	295.83

325

326

327

328   (a)   (b)
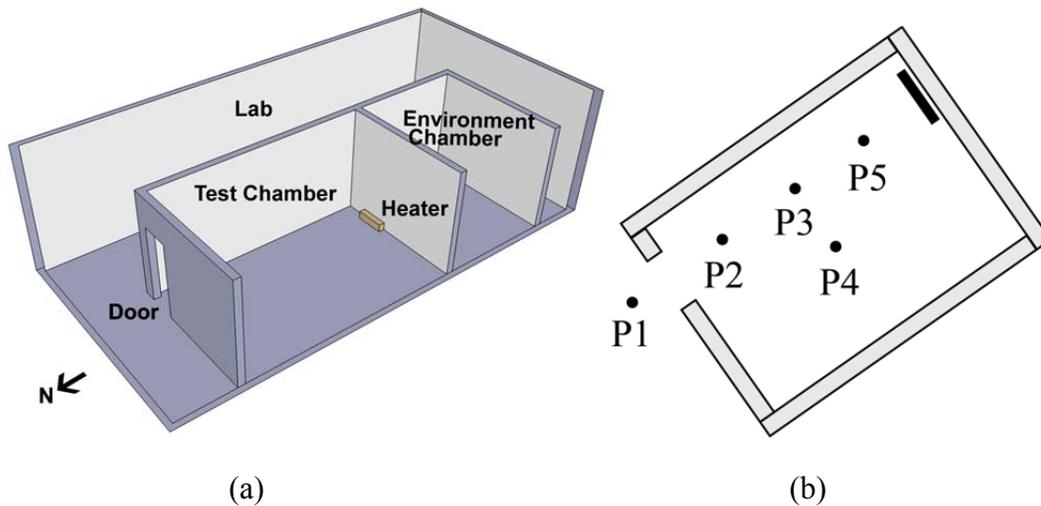
329   Figure 8 (a) Sketch of experimental setup and (b) measurement positions inside and outside the
330   test chamber

331

332   Our coarse-grid FFD used 18 × 10 × 10 mesh cells to represent the entire laboratory, while 6 × 6
333   × 6 cells were used for the test chamber, and a single cell of 0.3 m × 0.64 m × 0.7 m was used for
334   the baseboard heater. Note that the latter cell was much larger than the physical size of the
335   heater. In addition, the entire laboratory in this case was simulated by FFD with a fine grid of 30
336   × 20 × 20.

337   Figure 9 compares the vertical air temperature profiles simulated by coarse-grid FFD with and
338   without the plume model and by fine-grid FFD, with the experimental data at positions P1, P2,
339   P3, and P5. The temperature profiles predicted by coarse-grid FFD with the plume model and by
340   fine-grid FFD were in very good agreement with the measured profiles. Coarse-grid FFD without
341   the plume model resulted in temperature profiles that did not agree as well with the experimental
342   data. The model predicted a lower air temperature in the upper part of the room because of the
343   weaker–than-actual thermal plume predicted by the cell that was much larger than the heater.

344   Because the buoyancy flow generated by the heater was the primary driving force for the airflow
345   in the chamber, correct simulation of the thermal plume was crucial for the prediction of the
346   ventilation rate through the door opening. Table 3 compares the ventilation rates computed by
347   coarse-grid FFD with and without the plume model and by fine-grid FFD, with the experimental
348   data. Some uncertainties were observed in the measured data for this case. They resulted from
349   instabilities in the natural ventilation system, which caused the ventilation rate to vary from 9.18
350   to 12.6 ACH. Because of the reduced heat source intensity in the large heat source cell, a small
351   plume was predicted by the coarse-grid FFD without the plume model. As a result, the predicted
352   ventilation rate was low in comparison with the lowest measured ventilation rate. Both the
353   coarse-grid FFD with the plume model and the fine-grid FFD can accurately predict the
354   ventilation rate caused by the thermal plume from the heater in the chamber. Coarse-grid FFD
355   with the plume model predicted a ventilation rate that was equal to the mean value in the
356   experiment.

357   Table 3 Ventilation rates for the buoyancy-driven, single-side natural ventilation case obtained
358   from various models and the experiment by Jiang and Chen [26].
359

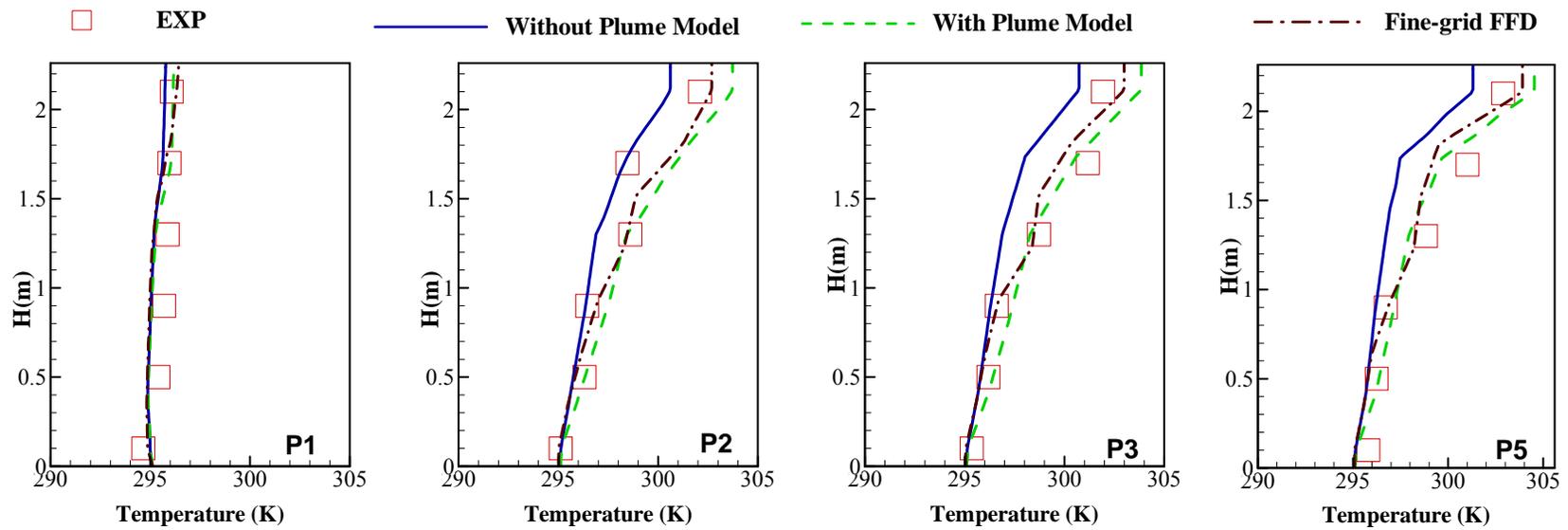| | Experiment | Fine-grid FFD | Coarse-grid FFD | |
			Without plume model	With plume model
Ventilation rate (ACH)	9.18-12.6	9.36	7.4	10.6

Figure 9 Comparison of the vertical air temperature profiles predicted by various FFD models with the experimental data from Jiang and Chen [26].

360   *3.4 Buoyancy-driven natural ventilation in an atrium*

361   This study also applied FFD to the simulation of buoyancy-driven natural ventilation in an
362   atrium, which is a major feature of modern natural ventilation designs for buildings. A tool that
363   predicts the ventilation rate of buoyancy-driven airflow through an atrium with reasonable
364   accuracy would effectively help architects improve natural ventilation designs. To test the
365   performance of coarse-grid FFD in simulating atrium-assisted natural ventilation, this study
366   simulated buoyancy-driven natural ventilation flows in a single-storey space connected to an
367   atrium [27].

368   Figure 10 shows a case in which a scaled building shell of 2.5 m (W) × 6.5 m (L) × 1.58 m (H)
369   was connected to an atrium with a height of 3.71 m (M). A heat source of 100 W was placed in
370   the center of the building to generate a buoyancy force. The building was ventilated through a
371   top-down chimney (TDC) located near the left wall. An opening was constructed in the ceiling of
372   the building shell connected to the atrium so that air could travel between the building and the
373   atrium. Because of the thermal plume generated by the heat source, ambient air flowed through
374   the TDC to the floor of the building and was entrained by the thermal plume. The plume
375   transported the air to ceiling level, and then the air travelled towards the atrium. Finally, the air
376   flowed outside through the outlet in the ceiling of the atrium. The ventilation rate would vary
377   with the area of the opening (A_u). With a constant effective single-storey opening area of 0.044
378   m^2 [27], this study varied the area of the atrium outlet in order to examine the performance of
379   coarse-grid FFD in simulating the natural ventilation rate (Q_v). A grid size of 13 × 8 × 5 was
380   used for the entire building.

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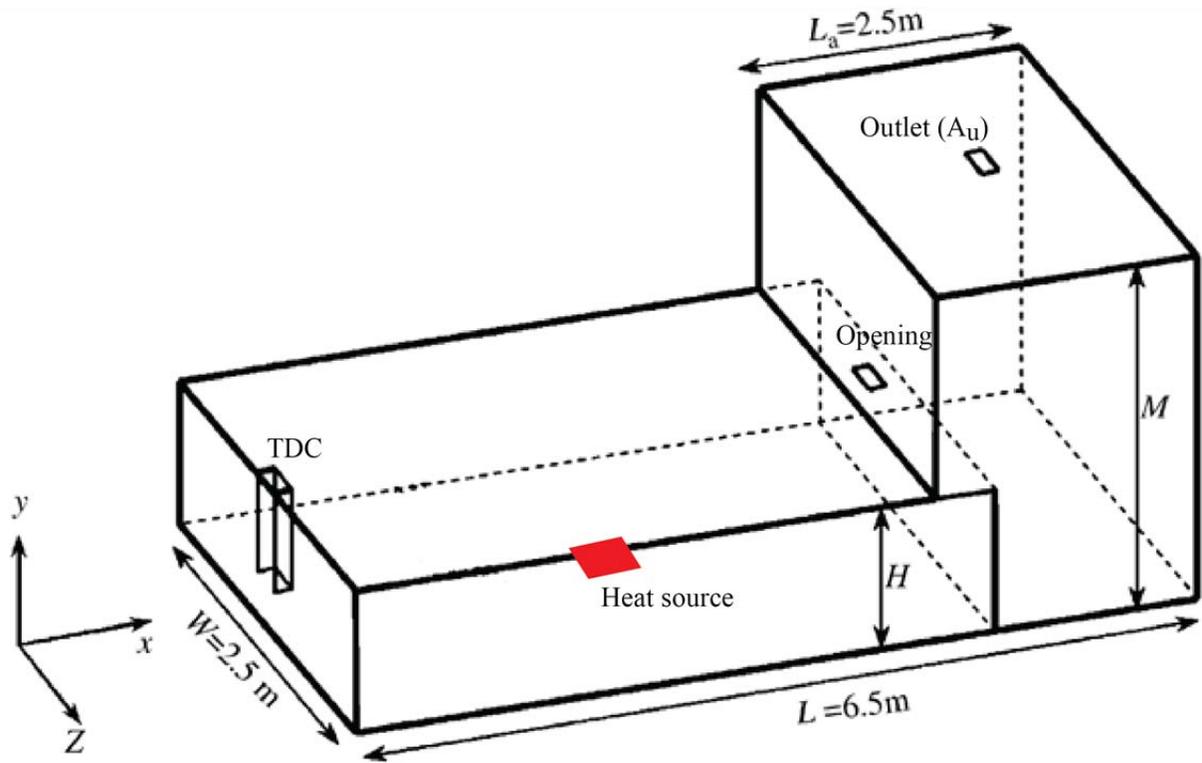
Figure 10 Geometry and opening locations of the scaled building with an atrium [27]

383

Figure 11 shows the relationship between the ventilation rate and the atrium outlet size as predicted by coarse-grid FFD with and without the plume model. Results obtained by CFD [27] and an analytical model [15] are also shown for the purpose of comparison. The analytical model was developed from experimental data that should be quite reliable. The ventilation rate (Q_v) was normalized by the volume flow rate (Q_H) of a pure thermal plume at ceiling height, and the atrium outlet size (A_u) was normalized by the squared height (H^2) of the building. As the size of the atrium outlet increased, the natural ventilation rate also increased. For coarse-grid FFD without the plume model, the predicted ventilation rate was significantly lower than that calculated by the analytical model. Because the heat source was represented by a very large cell, the predicted temperature at the heat source cell were likely to have been much lower, and the buoyancy force much smaller, than they would be in reality. As a result, coarse-grid FFD without the plume model could not predict the natural ventilation rate well. Coarse-grid FFD with the plume model was able to predict a flow rate that was close to that predicted by CFD. The FFD-predicted ventilation rate also agreed with that from the analytical model. Therefore, the plume model can effectively improve the accuracy of coarse-grid FFD for simulating buoyancy-driven natural ventilation.
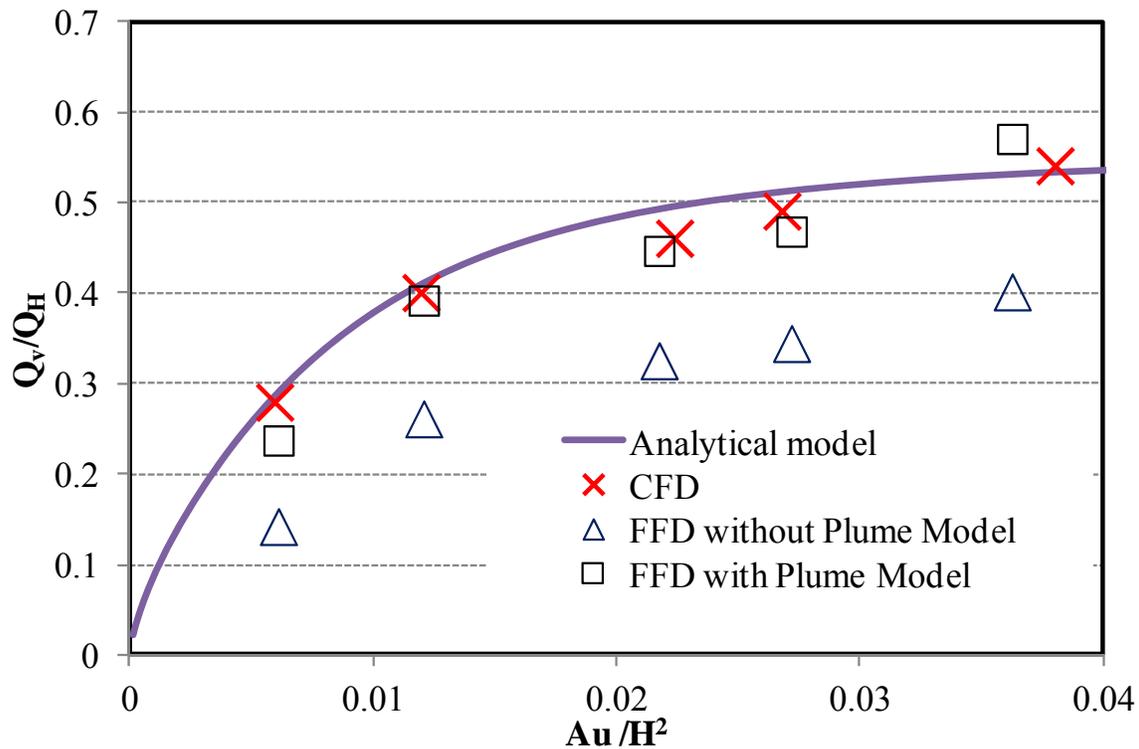
Figure 11 Variation in ventilation rate with outlet size for the scale building with an atrium as predicted by different models

4. Discussion

The tests described above demonstrated that the effectivness of the plume model. It helps coarse-grid FFD achieve better performance for simulating temperature stratification in rooms with buoyancy-driven ventilation and ventilation rate for buoyancy-driven natural ventilation. Through tuning the model coefficients for individual cases, it is possible to further improve the simulation results, as the development of thermal plume is different from case to case. However, because the purpose of FFD is not to obtain as accurate simulation as CFD does, the investigation on the impact of coefficients on improving the accuracy of coarse-grid FFD is limited in this study.

To examine the overall benefits of the plume model for FFD simulations, this study further investigated the computing time and accuracy of coarse-grid FFD with and without the plume model and fine-grid FFD for predicting temperature distributions. The computing time was recorded for each model by using the elapsed time when running the simulation for 1000 time steps. In addition, the root mean square errors based on the differences between the simulated

418 results and the experimental data were used to quantify the accuracy. The computing time and
419 errors in each case were normalized with those of the fine-grid FFD.

420 Table 4 shows that the coarse grid could significantly reduce computing time, by 5 to 50 times
421 according to the coarseness of the grid. On the other hand, the coarse grid produced slightly large
422 errors, but they do not seem to be critical for our applications. In a decision as to whether or not
423 to use FFD, computing time is a much more critical factor than computing accuracy. Coarse-grid
424 FFD with the plume model had errors that were similar to those of fine-grid FFD. Without the
425 plume model, the errors could be quite large, as shown in Table 4.

426 Table 4 Comparison of computing time and simulation error with the use of different FFD
427 models

Case	Model	Grid number (X × Y × Z)	Normalized computing time	Normalized error
1	Fine-grid FFD	20 × 20 × 20	1	1
	Coarse-grid FFD w/o plume model	5 × 5 × 5	0.019	2.43
	Coarse-grid FFD w/ plume model	5 × 5 × 5	0.019	1.67
2	Fine-grid FFD	25 × 18 × 16	1	1
	Coarse- grid FFD w/o plume model	5 × 5 × 8	0.032	1.44
	Coarse-grid FFD w/ plume model	5 × 5 × 8	0.033	1.03
3	Fine-grid FFD	30 × 20 × 20	1	1
	Coarse- grid FFD w/o plume model	18 × 10 × 10	0.226	3.03
	Coarse-grid FFD w/ plume model	18 × 10 × 10	0.226	1.46

428

Conclusions

430 This study proposed the integration of a thermal plume model into coarse-grid FFD in order to
431 more accurately simulate airflows driven by thermal plumes. The integration used the plume
432 flow rate calculated by the plume model to estimate the source for the momentum equations in
433 FFD. The integration also used the plume model to predict the air temperature at the heat source
434 cell for correcting the prediction by the energy equation in FFD.

435 This study applied coarse-grid FFD with the plume model to simulate airflows in buildings
436 driven by thermal buoyancy forces. Displacement ventilation and buoyancy-driven natural
437 ventilation cases were used. Experimental data and CFD results from the literature were used to
438 compare with the predicted results by the integrated FFD model. Coarse-grid FFD with the
439 plume model greatly improved FFD performance in predicting air velocity and temperature in
440 the plume, which led to a more accurate prediction of mean air temperature stratification in the
441 displacement ventilation cases. The results were comparable to those of fine-grid FFD.

Coarse-grid FFD with the plume model was also able to accurately predict the ventilation rate in buoyancy-driven natural ventilation cases. The accuracy was comparable to that of fine-grid FFD, CFD, or an analytical model developed from experimental data.

The computing time for coarse-grid FFD with and without the plume model was a small fraction of that for fine-grid FFD for all the cases tested. The coarse-grid number was actually the minimum necessary for representing room air distribution. At the same time, the significant reduction in grid number did not lead to a much larger error than that of fine-grid FFD, especially for coarse-grid FFD with the plume model.

Acknowledgement

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