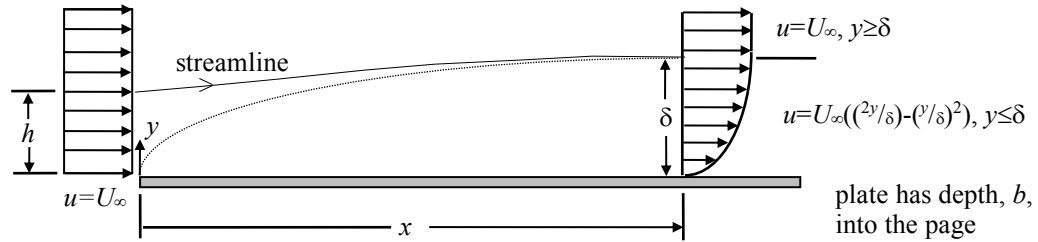


An incompressible, viscous fluid with density, ρ , flows past a solid flat plate which has a depth, b , into the page. The flow initially has a uniform velocity U_∞ , before contacting the plate. After contact with the plate at a distance x downstream from the leading edge, the flow velocity profile is altered due to the no-slip condition. The velocity profile at location x is estimated to have a parabolic shape, $u=U_\infty((2y/\delta)-(y/\delta)^2)$, for $y \leq \delta$ and $u=U_\infty$ for $y \geq \delta$ where δ is termed the “boundary layer thickness.”



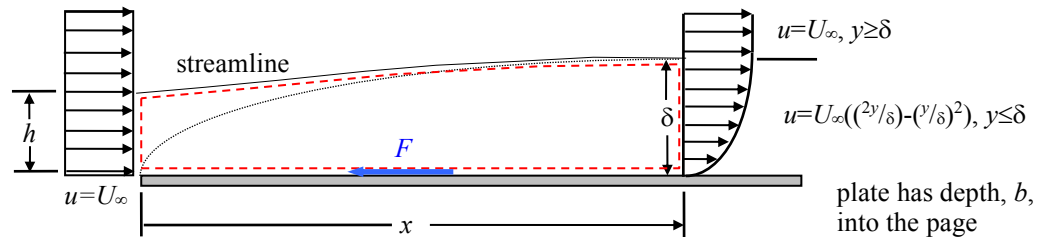
1. Determine the upstream height from the plate, h , of a streamline which has a height, δ , at the downstream distance x . Express your answer in terms of δ .
2. Determine the force the plate exerts on the fluid over the distance x . Express your answer in terms of ρ , U_∞ , b , and δ . You may assume that the pressure everywhere is p_∞ . The force the drag exerts on the plate is called the “skin friction” drag.

BRIEF SOLUTION:

1. Apply conservation of mass to a control volume that is adjacent to the plate, crosses perpendicularly to the stream at the leading edge of the plate, follows a streamline, and crosses perpendicularly to the stream at the location where the boundary layer has thickness, δ . Note that there is no flow across a streamline.
2. Apply the linear momentum equation to the same control volume used in Step 1. Be sure to include the force the plate exerts on the control volume.

DETAILED SOLUTION:

Apply conservation of mass to the fixed control volume shown below.



$$\frac{d}{dt} \int_{CV} \rho dV + \int_{CS} \rho \mathbf{u}_{rel} \cdot d\mathbf{A} = 0$$

where

$$\frac{d}{dt} \int_{CV} \rho dV = 0 \quad (\text{steady flow})$$

$$\begin{aligned} \int_{CS} \rho \mathbf{u}_{rel} \cdot d\mathbf{A} &= -\rho U_{\infty} h b + \int_{y=0}^{y=\delta} \rho U_{\infty} \left[2 \frac{y}{\delta} - \frac{y^2}{\delta^2} \right] dy b = -\rho U_{\infty} h b + \rho U_{\infty} \left(\delta - \frac{1}{3} \delta \right) b \\ &= -\rho U_{\infty} h b + \frac{2}{3} \rho U_{\infty} \delta b \end{aligned}$$

(Note that there is no flow across the streamline.)

Substitute into conservation of mass and solve for h .

$$\boxed{h = \frac{2}{3} \delta} \quad (1)$$

Now apply the linear momentum equation in the x -direction on the same control volume.

$$\frac{d}{dt} \int_{CV} u \rho dV + \int_{CS} u (\rho \mathbf{u}_{rel} \cdot d\mathbf{A}) = F_{B,x} + F_{S,x}$$

where

$$\frac{d}{dt} \int_{CV} u \rho dV = 0 \quad (\text{steady flow})$$

$$\begin{aligned}
 \int_{CS} u(\rho \mathbf{u}_{rel} \cdot d\mathbf{A}) &= -\rho U_{\infty}^2 hb + \int_{y=0}^{y=\delta} \rho U_{\infty}^2 \left[2\frac{y}{\delta} - \frac{y^2}{\delta^2} \right]^2 dyb \\
 &= -\rho U_{\infty}^2 hb + \rho U_{\infty}^2 b \int_0^{\delta} \left[4\frac{y^2}{\delta^2} - 4\frac{y^3}{\delta^3} + \frac{y^4}{\delta^4} \right] dy \\
 &= -\rho U_{\infty}^2 hb + \rho U_{\infty}^2 b \left[\frac{4}{3}\delta - \delta + \frac{1}{5}\delta \right] \\
 &= -\rho U_{\infty}^2 hb + \frac{8}{15} \rho U_{\infty}^2 b\delta
 \end{aligned}$$

$$F_{B,x} = 0$$

$$F_{S,x} = -F \quad (\text{the pressure everywhere is } p_{\infty})$$

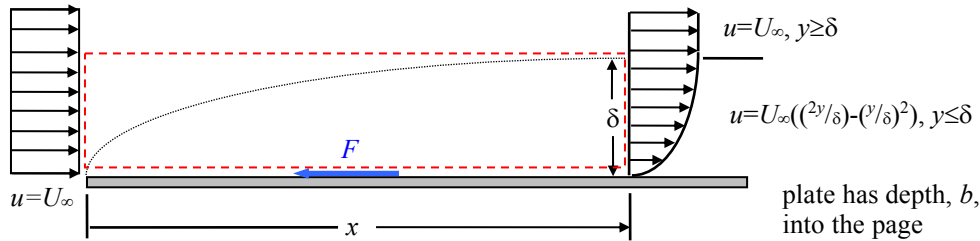
Substitute and simplify, making use of Eqn. (1).

$$-\rho U_{\infty}^2 \left(\frac{2}{3} \delta \right) b + \frac{8}{15} \rho U_{\infty}^2 b\delta = -F$$

$$\boxed{F = \frac{2}{15} \rho U_{\infty}^2 b\delta}$$

(2)

We could have also determined the force using a different control volume as shown below.



Determine the mass flow rate out of the control volume through the top using conservation of mass.

$$\frac{d}{dt} \int_{CV} \rho dV + \int_{CS} \rho \mathbf{u}_{rel} \cdot d\mathbf{A} = 0$$

where

$$\frac{d}{dt} \int_{CV} \rho dV = 0 \quad (\text{steady flow})$$

$$\int_{CS} \rho \mathbf{u}_{rel} \cdot d\mathbf{A} = -\rho U_{\infty} \delta b + \int_{y=0}^{y=\delta} \rho U_{\infty} \left[2\frac{y}{\delta} - \frac{y^2}{\delta^2} \right] dyb + m_{top} = -\frac{1}{3} \rho U_{\infty} \delta b + m_{top}$$

Substitute and solve for the mass flow rate.

$$m_{top} = \frac{1}{3} \rho U_{\infty} \delta b$$

(3)

Now apply the linear momentum equation in the x -direction to the same control volume.

$$\frac{d}{dt} \int_{CV} u \rho dV + \int_{CS} u (\rho \mathbf{u}_{rel} \cdot d\mathbf{A}) = F_{B,x} + F_{S,x}$$

where

$$\frac{d}{dt} \int_{CV} u \rho dV = 0 \quad (\text{steady flow})$$

$$\begin{aligned} \int_{CS} u (\rho \mathbf{u}_{rel} \cdot d\mathbf{A}) &= -\rho U_{\infty}^2 \delta b + \int_{y=0}^{y=\delta} \rho U_{\infty}^2 \left[2 \frac{y}{\delta} - \frac{y^2}{\delta^2} \right]^2 dy b + m_{top} U_{\infty} \\ &= -\frac{7}{15} \rho U_{\infty}^2 \delta b + m_{top} U_{\infty} \end{aligned}$$

(Note that the horizontal component of the velocity at the top is U_{∞} since it's outside of the boundary layer.)

$$F_{B,x} = 0$$

$$F_{S,x} = -F \quad (\text{the pressure everywhere is } p_{\infty})$$

Substitute and simplify making use of Eqn. (3).

$$-\frac{7}{15} \rho U_{\infty}^2 \delta b + \left(\frac{1}{3} \rho U_{\infty} \delta b \right) U_{\infty} = -F$$

$$F = \frac{2}{15} \rho U_{\infty}^2 \delta b \quad (\text{This is the same answer as before!})$$