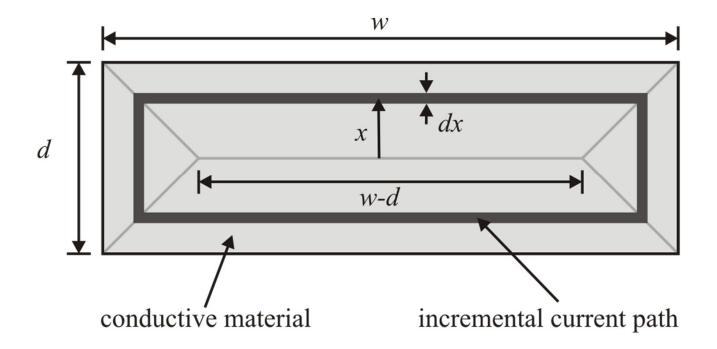
Power Magnetic Devices: A Multi-Objective Design Approach

Chapter 6: Magnetic Core Loss

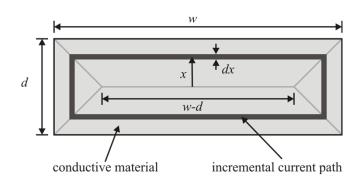
- Eddy current lost is a resistive power loss associated with induced currents
- Let's consider a rectangular conductor



Defining

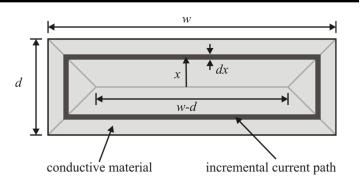
$$k_1 = \min(w, d)$$

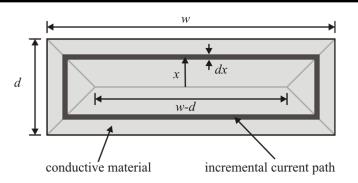
$$k_2 = |w - d|$$

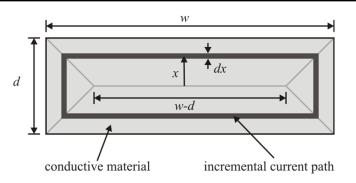


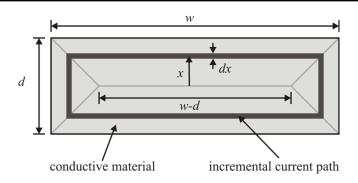
We can show

$$P = \frac{\sigma l}{128} \left(4k_1^4 + 8k_2k_1^3 + 2k_2^2k_1^2 - 2k_2^3k_1 + k_2^4 \ln\left(1 + \frac{2k_1}{k_2}\right) \right) \frac{1}{T} \int_0^T \left(\frac{dB}{dt}\right)^2 dt$$









• Example 6.1A. Eddy current loss in a toroid

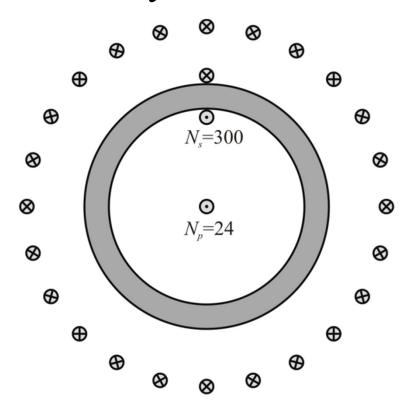


Table 6.1A-1 Aluminum alloy test samples.

	Material	σ (MS/m)	μ_r	r_i (mm)	r_o (mm)	w (mm)	<i>d</i> (mm)
	6061T6	25.3	1.23	140	150	10	12.7
S.D. Sudhoff, Power Mag	gnetic 60:1:3 : A Mu	lti-Obj <mark>ecil</mark> ve <mark>Design</mark> 2	$_{4pprodch}20$	140	150	10	27.2

Example 6.1A. Normalized power loss

$$p_{n} = \frac{P}{V \frac{1}{T} \int_{0}^{T} \left(\frac{dB}{dt}\right)^{2} dt}$$

• Thus, with our expression for eddy current loss, we have

$$p_n = \frac{\sigma \left(4k_1^4 + 8k_2k_1^3 + 2k_2^2k_1^2 - 2k_2^3k_1 + k_2^4 \ln\left(1 + \frac{2k_1}{k_2}\right)\right)}{128wd}$$

• Example 6.1A. Measured power loss. We can show

$$P = \frac{N_p}{N_s} \frac{1}{T} \int_0^T v_s i_p dt$$

• Example 6.1A. Derivation of average power.

• Example 6.1A. Derivation of average power.

• Example 6.1A. Now we measured power, we can find the measured normalized power loss density

$$p_n = \frac{P}{V \frac{1}{T} \int_0^T \left(\frac{dB}{dt}\right)^2 dt}$$

• Clearly – we need to know B. There are two approaches.

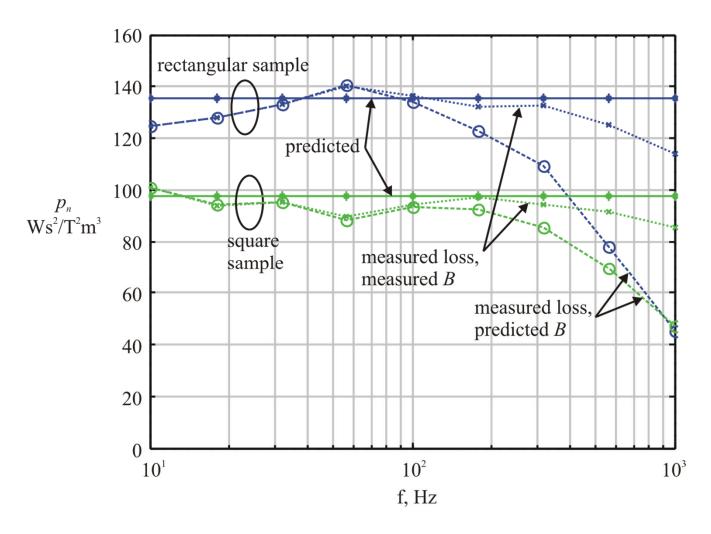
- Example 6.1A. Approach 1 to finding B (predicted)
- We use

$$B = \frac{\mu_0 \mu_r i_p}{2\pi r}$$

- Example 6.1A. Approach 1 to finding B (measured)
- We use

$$\frac{dB}{dt} = \frac{v_s}{wdN_s}$$

• Example 6.1A. Results:



Sinusoidal excitation

$$B = B_{pk} \cos(\omega_e t) \qquad \frac{1}{T} \int_0^T \left(\frac{dB}{dt}\right)^2 dt = \frac{1}{2} \omega_e^2 B_{pk}^2$$

Periodic excitation

$$B = \sum_{k=1}^{K} B_{a,k} \cos(k\omega_e t) + B_{b,k} \sin(k\omega_e t)$$

$$\frac{1}{T} \int_{0}^{T} \left(\frac{dB}{dt} \right)^{2} dt = \frac{1}{2} \omega_{e}^{2} \sum_{k=1}^{K} k^{2} \left(B_{a,k}^{2} + B_{b,k}^{2} \right)$$

- Thin laminations. Suppose $k_2 >> k_1$
- Then we can show

$$P = \frac{\sigma l \left(4k_1^4 + 8k_2k_1^3 + 2k_2^2k_1^2 - 2k_2^3k_1 + k_2^4 \ln\left(1 + \frac{2k_1}{k_2}\right)\right)}{128} \frac{1}{T} \int_0^T \left(\frac{dB}{dt}\right)^2 dt$$

Can be approximated as

$$P = \frac{\sigma l}{12} k_2 k_1^3 \frac{1}{T} \int_0^T \left(\frac{dB}{dt}\right)^2 dt$$

• Derivation (1/2)

• Derivation (2/2)

It follows that

$$p = \frac{\sigma w^2}{12} \frac{1}{T} \int_0^T \left(\frac{dB}{dt}\right)^2 dt$$

• To show this, we start with

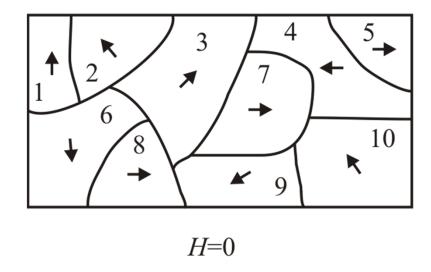
$$P = \frac{\sigma l}{12} k_2 k_1^3 \frac{1}{T} \int_0^T \left(\frac{dB}{dt}\right)^2 dt$$

• Continuing ...

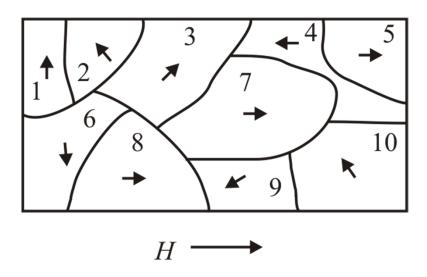
Importance of

$$p = \frac{\sigma w^2}{12} \frac{1}{T} \int_0^T \left(\frac{dB}{dt}\right)^2 dt$$

Domain wall motion

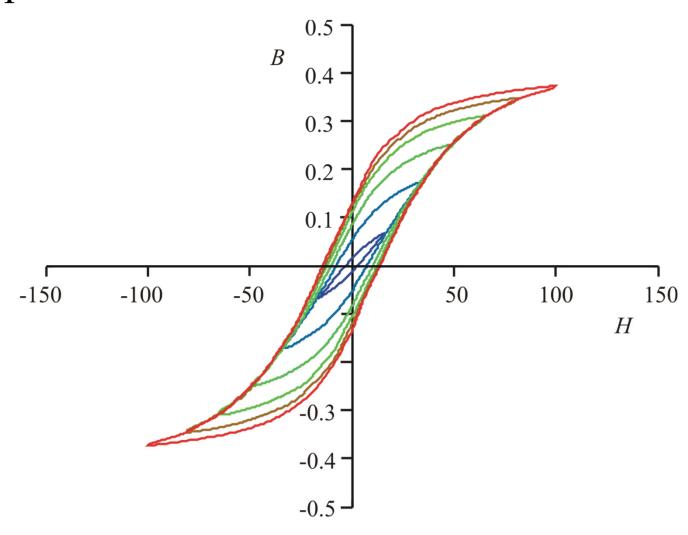


(a) domains without applied field



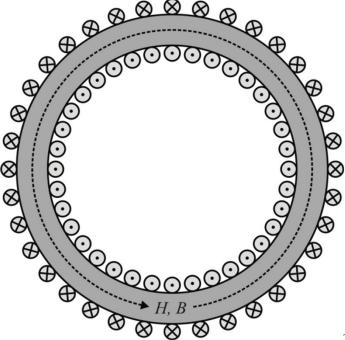
(b) domains with applied field

• Sample B-H characteristics

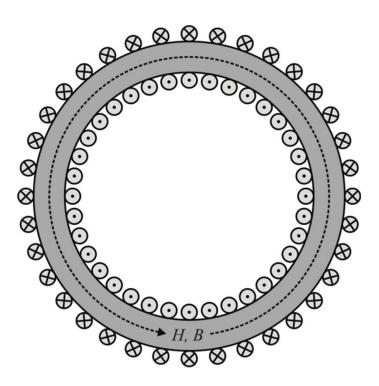


• The area of a B-H trajectory represent energy loss.

- The first step to do this is to show $e = \int_{B_0}^{1} HdB$
- To do end, consider the toroid



• Proceeding ...



• Now let us consider a closed path

$$e = \int_{B_0}^{B_1} HdB + \int_{B_0}^{B_2} HdB$$

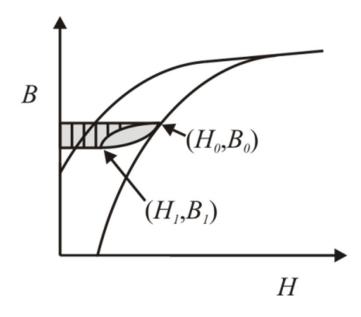
$$= \int_{B_0}^{B_0} HdB + \int_{B_0}^{B_0} HdB$$

$$= \int_{B_0}^$$

Minor loop loss

$$e = \int_{B_0}^{B_1} HdB + \int_{B_0}^{B_0} HdB$$

$$= \int_{B_0}^{B_0} HdB + \int_{B_0}^{B_0} HdB$$
Term 1 Term 2



• Now that we have established the energy loss per cycle, the average power loss may be expressed

$$p_h = fe_{BH}$$

The Steinmetz Equation

$$e_{BH} = \frac{k_h}{f_b} \left(\frac{f}{f_b}\right)^{\alpha - 1} \left(\frac{B_{pk}}{B_b}\right)^{\beta}$$

$$p_h = k_h \left(\frac{f}{f_b}\right)^{\alpha} \left(\frac{B_{pk}}{B_b}\right)^{\beta}$$

- The Modified Steinmetz Equation (MSE)
 - Motivated by observation that localized eddy current due to domain movement being tied rate of change
- MSE equivalent frequency

$$\Delta B = B_{mx} - B_{mn}$$

$$\mathcal{B} = \frac{1}{\Delta B} \int_{0}^{T} \left(\frac{dB}{dt}\right)^{2} dt$$

$$f_{eq} = \frac{2}{\Delta B^2 \pi^2} \int_0^T \left(\frac{dB}{dt}\right)^2 dt$$

MSE loss equation

$$e_{BH} = \frac{k_h}{f_b} \left(\frac{f_{eq}}{f_b}\right)^{\alpha - 1} \left(\frac{\Delta B}{2B_b}\right)^{\beta}$$

$$p_h = k_h \left(\frac{f_{eq}}{f_b}\right)^{\alpha - 1} \left(\frac{\Delta B}{2B_b}\right)^{\beta} \frac{f}{f_b}$$

• MSE with sinusoidal flux density

• Example 6.3A. MSE with triangular excitation

- Generalized Steinmetz Equation (GSE)
 - Assumes instantaneous loss is a function of rate of change and flux density value

$$p_i = f\left(\frac{dB}{dt}, B\right)$$

- A suggested choice is

$$p_{i} = k \left| \frac{1}{f_{b}B_{b}} \frac{dB}{dt} \right|^{\alpha} \left| \frac{B}{B_{b}} \right|^{\beta - \alpha}$$

The GSE then becomes

$$p_{h} = \frac{1}{T} k \int_{0}^{T} \left| \frac{1}{f_{b} B_{b}} \frac{dB}{dt} \right|^{\alpha} \left| \frac{B}{B_{b}} \right|^{\beta - \alpha} dt$$

where we choose

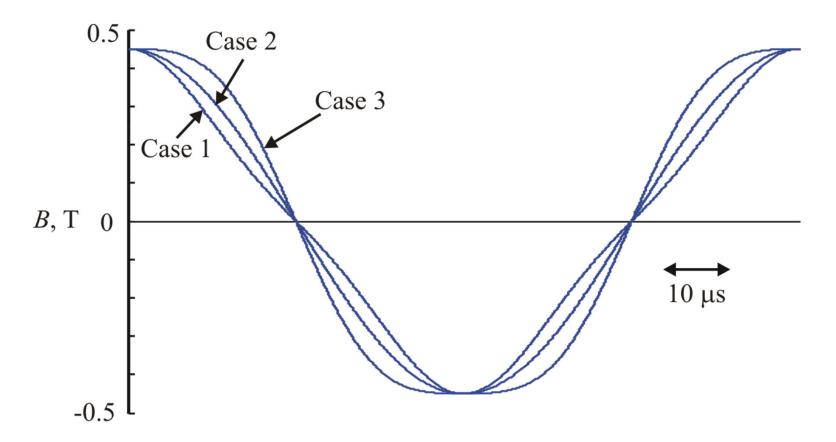
$$k = \frac{k_h}{(2\pi)^{\alpha-1} \int_{0}^{2\pi} |\cos \theta|^{\alpha} |\sin \theta|^{\beta-\alpha} d\theta}$$

- Example 6.3B. Consider MN60LL ferrite with α =1.034, β =2.312, k_h = 40.8 W/m³
- Let us compare MSE and GSE models for

$$B = B_1 \cos(2\pi ft) + B_3 \cos(6\pi ft)$$

- Cases
 - Case 1: B_1 =0.5 T, B_3 =-0.05T
 - Case 2: B_1 = 0.45T, B_3 =0 T
 - Case 3: B_1 = 0.409 T, B_3 =0.0409 T

• Example 6.3 flux density waveforms



Results

- MSE model: $87.4, 86.5, 86.2 \text{ kW/m}^3$

- GSE model: 86.9, 86.5, 86.8 kW/m³

Comments on MSE and GSE models

- Combined loss modeling
 - Eddy current loss

$$p_e = k_e \frac{1}{T} \int_0^T \left(\frac{dB}{dt}\right)^2 dt$$

- Hysteresis loss (MSE)

$$p_h = k_h \left(\frac{f_{eq}}{f_b}\right)^{\alpha - 1} \left(\frac{\Delta B}{2B_b}\right)^{\beta} \frac{f}{f_b}$$

Combined loss model

$$p_{t} = k_{h} \left(\frac{f_{eq}}{f_{b}}\right)^{\alpha} \left(\frac{\Delta B}{2B_{b}}\right)^{\beta} f + k_{e} \frac{1}{T} \int_{0}^{T} \left(\frac{dB}{dt}\right)^{2} dt$$

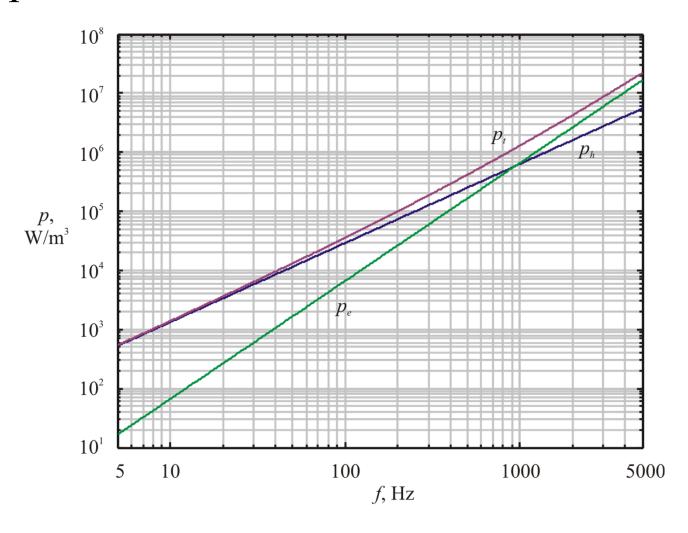
• Example 6.3C. Lets look at losses in M19 steel with k_h =50.7 W/m³, α =1.34, β =1.82, and k_e =27.5·10⁻³ Am/V. We will assume

$$B = B_{pk} \cos(2\pi f t)$$

• Since the flux density is sinusoidal

$$p_e = 2\pi^2 k_e B_{pk}^2 f^2 \qquad p_h = k_h \left(\frac{f}{f_b}\right)^{\alpha} \left(\frac{B_{pk}}{B_b}\right)^{\beta}$$

• Example 6.3C results



- These models can capture effects such as
 - waveforms with dc offset
 - aperiodic excitation
 - waveforms with minor loops
- Two approaches
 - Jiles-Atherton
 - Praisach

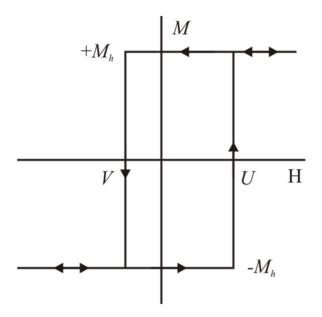
Jiles-Atherton model

$$B = \mu_0(H + M)$$

$$M = M_{irr} + M_{rev}$$

$$\begin{split} M &= M_{irr} + M_{rev} \\ \frac{dM_{irr}}{dt} &= \left(\frac{M_{an}(H) - M_{irr}}{k \delta - \alpha(M_{an}(H) - M_{irr})}\right) \frac{dH}{dt} \\ M_{rev} &= c(M_{an}(H) - M_{irr}) \\ \delta &= \begin{cases} 1 & \frac{dH}{dt} > 0 \\ 0 & \frac{dH}{dt} = 0 \\ -1 & \frac{dH}{dt} < 0 \end{cases} \\ \text{S.D. Sudhoff, Power Magnetic Devices: A Multi-Objective Design Approach} \end{split}$$

Preisach model is based on behavior of a hysteron



Total magnetization based on

$$M = \int_{-\infty}^{\infty} \int_{-\infty}^{U} Q(U, V) P(Q, V) dV dU$$

Saturated magnetization

$$M_s = \int_{-\infty}^{\infty} \int_{-\infty}^{U} P(U, V) dV dU$$

Normalized magnetization

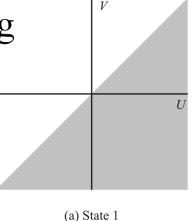
$$m = \frac{M}{M_s}$$

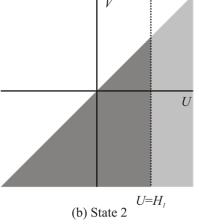
Normalized hysteron density

$$p(U,V) = \frac{1}{M_s} P(U,V)$$

- Lets consider behavior of model
- Start in State 1 all hysterons neg
- State 1 to State 2

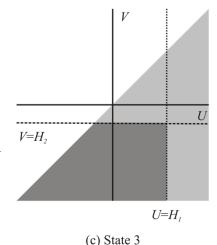
$$m_{s2} = m_{s1} + 2 \int_{-\infty}^{H_1} \int_{-\infty}^{U} p(U, V) dV dU$$

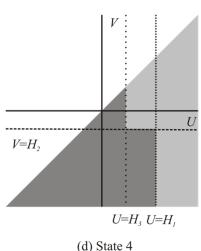




State 2 to State 3

$$m_{s3} = m_{s2} - 2 \int_{H_2}^{H_1} \int_{V}^{H_1} p(U,V) dU dV$$





State 3 to State 4

$$m_{s4} = m_{s3} + 2 \int_{H_2}^{H_3} \int_{H_2}^{U} p(U,V) dV dU$$
(a) State 1
(b) State 2
(c) V
(e) V
(f) V
(f)

(c) State 3

(d) State 4

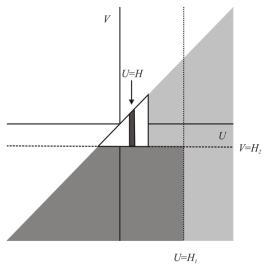
Reconsider State 3 to State 4

$$m = m_{s3} + 2 \int_{H_2}^{H} \int_{H_2}^{U} p(U, V) dV dU$$
$$\frac{dm}{dH} = 2 \int_{H_2}^{H} p(H, V) dV$$

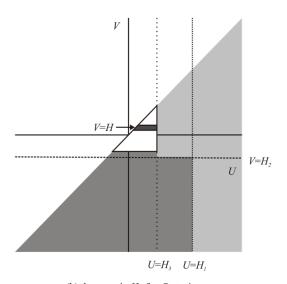
Suppose field decreases after State 4

$$m = m_{s4} - 2 \int_{H}^{H_3} \int_{V}^{H_3} p(U, V) dU dV$$

$$\frac{dm}{dH} = 2 \int_{H}^{H_3} p(U, H) dU$$



(a) transition from State 3 to State 4



(b) decrease in H after State 4

Completing the model

$$\frac{dm}{dt} = \frac{dm}{dH} \frac{dH}{dt}$$

$$M = M_s m$$

$$B = \mu_0 H + M$$

• Relative advantages and disadvantages of approaches