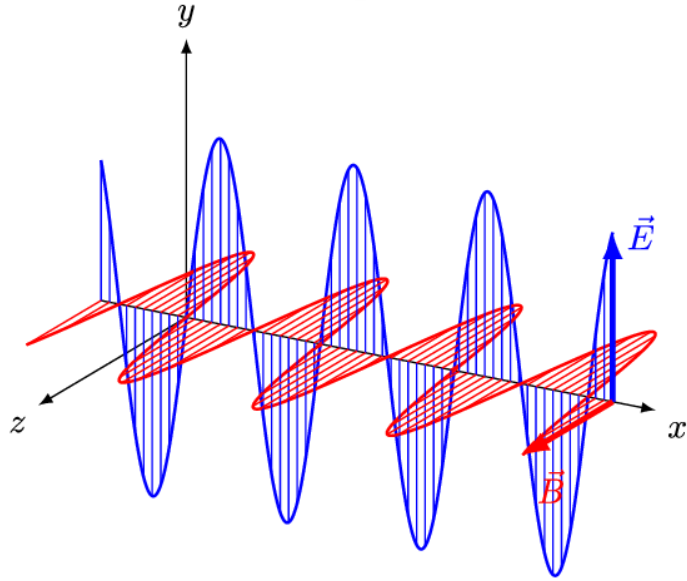


Notes on Fourier Optics

- How does light propagate through free space?
- What do lenses do?
- How does a directional antenna work?
- Four important ideas:
 - Wave equation
 - Optical plane wave and direction of transmission/arrival
 - Fresnel approximation: Assumes light is traveling through an aperture
 - Fraunhofer approximation: Assumes light is far from the aperture

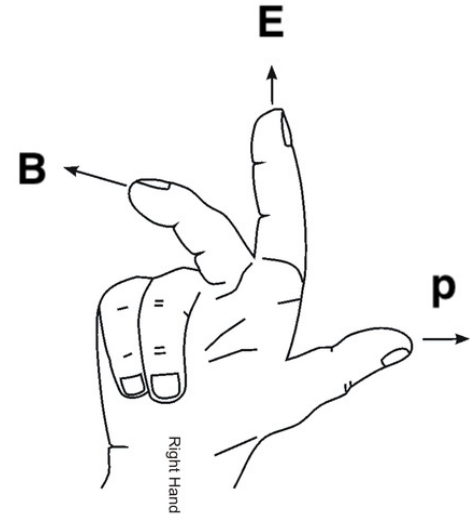
Electro-Magnetic Waves



- Plane linearly polarized wave

- EM waves are vector fields

- E and B are traveling vector fields
- E is perpendicular to B
- $\vec{p} = E \times B$ where \vec{p} is the direction of travel



The Wave Equation

- From Maxwell's equations:

- We can show that

$$\nabla^2 \vec{E} - \mu\epsilon \frac{\partial^2 \vec{E}}{\partial t^2} = 0$$

This is the wave equation in 3 dimensions.

- For simplicity, we consider a single component of the electric field vector:

$$\nabla^2 E - \mu\epsilon \frac{\partial^2 E}{\partial t^2} = 0$$

- where μ is the permeability; and ϵ is the dielectric constant; and $c = \frac{1}{\sqrt{\mu\epsilon}}$ where c is the speed of light in the medium.

Solution to the Wave Equation in Time

- The wave equation is given by

$$\nabla^2 E - \frac{1}{c^2} \frac{\partial^2 E}{\partial t^2} = 0$$

- Separation of variables gives solutions of the form

$$\begin{aligned} E_t(x, y, z, t) &= E_o \exp[j2\pi(v_x x + v_y y + v_z z)] \exp[-j2\pi f_o t] \\ &= E_o \exp[j2\pi(v_x x + v_y y + v_z z - f_o t)] \end{aligned}$$

– where

$$\sqrt{v_x^2 + v_y^2 + v_z^2} = v_o = \frac{1}{\lambda} = \frac{f_o}{c}$$

Wavelength (m) $\rightarrow \lambda = \frac{1}{v_o} = \frac{c}{f_o}$ \leftarrow Velocity (m/sec)

spatial frequency (m⁻¹) $\rightarrow \frac{1}{\lambda} = v_o = \frac{f_o}{c}$ \leftarrow Frequency (Hz)

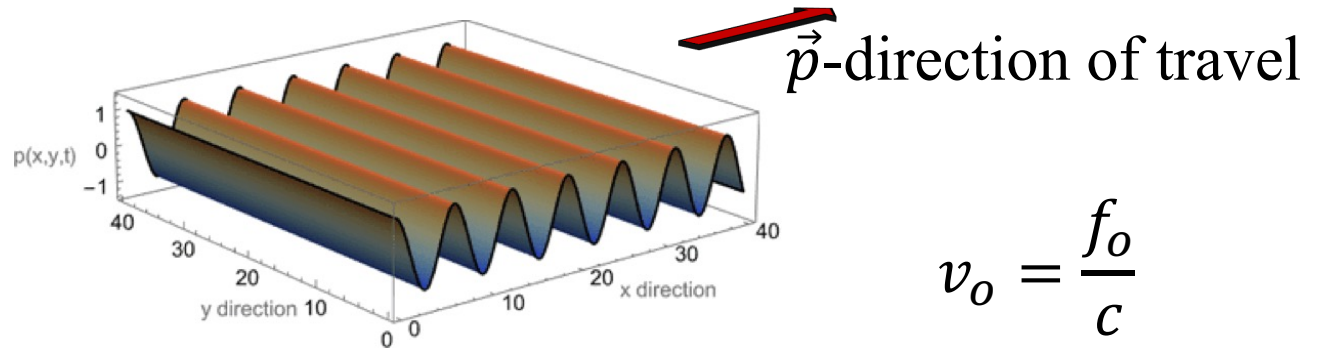
- We will assume that f_o , v_o , and c are all fixed.

Plane Waves

- Plane waves have the form

$$\begin{aligned} E_t(x, y, z, t) &= E_o \exp[j2\pi(v_x x + v_y y + v_z z)] \exp[-j2\pi f_o t] \\ &= E_o \exp[j2\pi(v_o(\vec{p} \cdot \vec{r}) - f_o t)] \end{aligned}$$

- where $\vec{r} = [x, y, z]$, $\vec{v} = [v_x, v_y, v_z]$, $v_o = \|\vec{v}\|$ and $\vec{p} = \frac{\vec{v}}{\|\vec{v}\|}$.



Plane Wave Interpretation

- Solution has the form:

$$E_t(x, y, z, t) = E_o \exp[j2\pi(v_x x + v_y y + v_z z)] \underbrace{\exp[-j2\pi f_o t]}_{\text{carrier wave}}$$

- Electric field is the real component:

$$\text{Electric Field} = \text{Re}\{E_t(x, y, z, t)\}$$

- So, we can always recover the space-time solution from:

$$E(x, y, z) = E_o \exp[j2\pi(v_x x + v_y y + v_z z)]$$

– where

$$\sqrt{v_x^2 + v_y^2 + v_z^2} = v_o = \frac{1}{\lambda}$$

1D Plane Wave Interpretation

- Solution has the form:

$$\begin{aligned} E_t(x, t) &= E_o \exp[j2\pi(v_o x)] \exp[-j2\pi f_o t] \\ &= E_o \exp \left[j2\pi v_o \left(x - \frac{v_o}{f_o} t \right) \right] \end{aligned}$$

Velocity (m/sec)

- Then the electric field is given by

$$\text{Electric Field} = \text{Re}\{E_t(x, y, z, t)\}$$

$$\begin{aligned} &= \text{Re} \left\{ E_o \exp \left[j2\pi v_o \left(x - \frac{v_o}{f_o} t \right) \right] \right\} \\ &= A_o \text{Re} \left\{ \exp \left[j2\pi v_o \left(x - \frac{v_o}{f_o} t \right) + j\theta \right] \right\} \\ &= A_o \cos \left\{ 2\pi v_o \left(x - \frac{v_o}{f_o} t \right) + \theta \right\} \end{aligned}$$

$$E_o = A_o e^{j\theta}$$

- So the plane wave is completely specified by

$$E(x, t) = E_o \exp[j2\pi(v_o x)]$$

2D Solutions

- In 3D, we have that

$$E(x, y, z) = E_o \exp[j2\pi(v_x x + v_y y + v_z z)]$$

- where

$$\sqrt{v_x^2 + v_y^2 + v_z^2} = v_o = \frac{1}{\lambda}$$

- So, if $v_x = 0$, we have that

$$E(y, z) = E_o \exp[j2\pi(v_y y + v_z z)]$$

- where

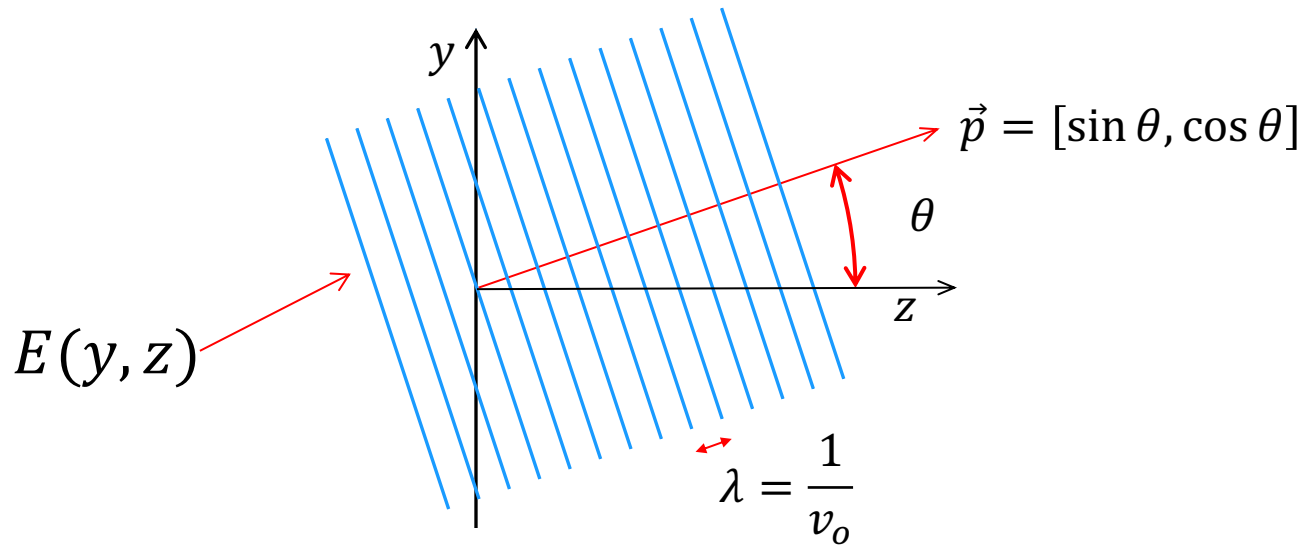
$$v_y = v_o \sin \theta$$

$$v_z = v_o \cos \theta$$

Direction of Propagation for 2D Solution

- If $v_x = 0$,

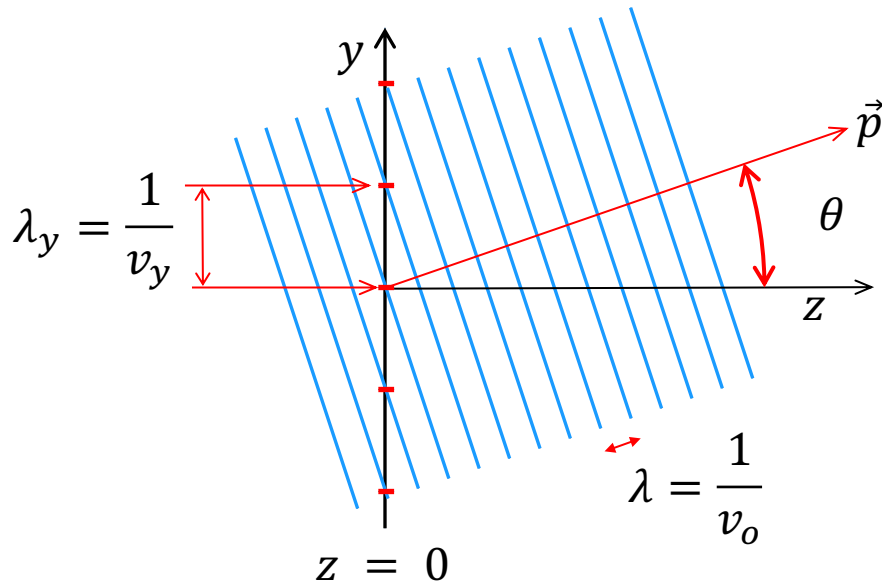
$$E(y, z) = E_o \exp[j2\pi(y v_o \sin \theta + z v_o \cos \theta)]$$



Intuition from 2D Solution

- The E field on the image plane ($z = 0$) is given by:

$$\begin{aligned}u_0(y) &= E_o \exp[j2\pi(y v_o \sin \theta + 0 v_o \cos \theta)] \\ &= E_o \exp[j2\pi(y v_o \sin \theta)]\end{aligned}$$



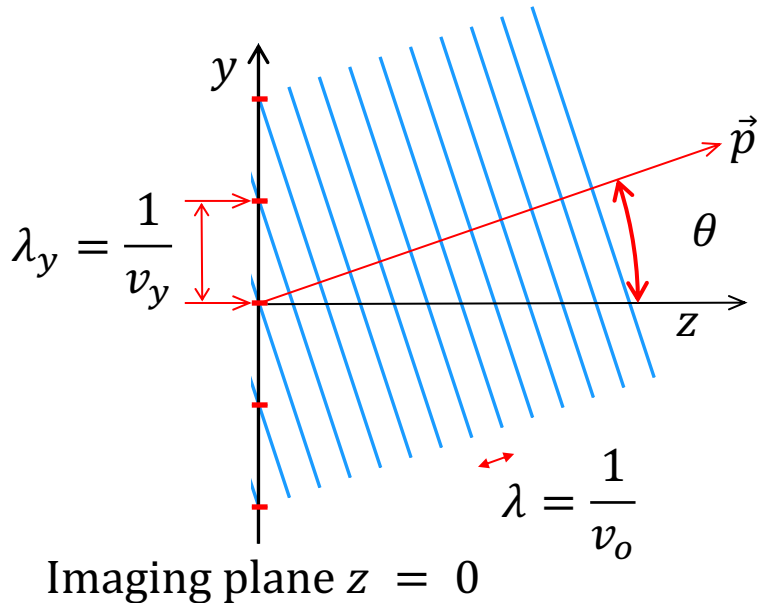
$$v_y = v_o \sin \theta$$

$$\lambda_y = \frac{\lambda}{\sin \theta}$$

Direction of Transmission in 2D

- A plane wave arriving at angle θ creates an electric field on the image plane given by

$$u_0(y) = E_o \exp[j2\pi(y \underbrace{v_o \sin \theta}_{v_y})]$$



- Angle of departure:

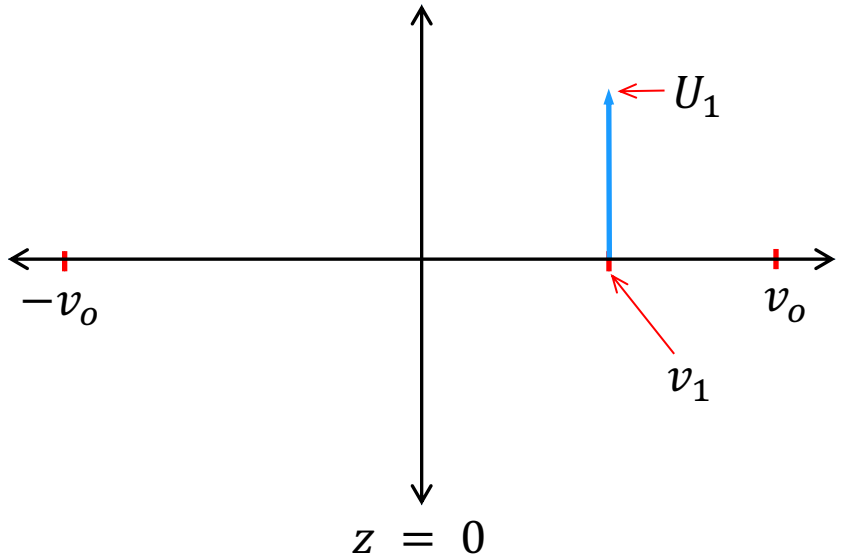
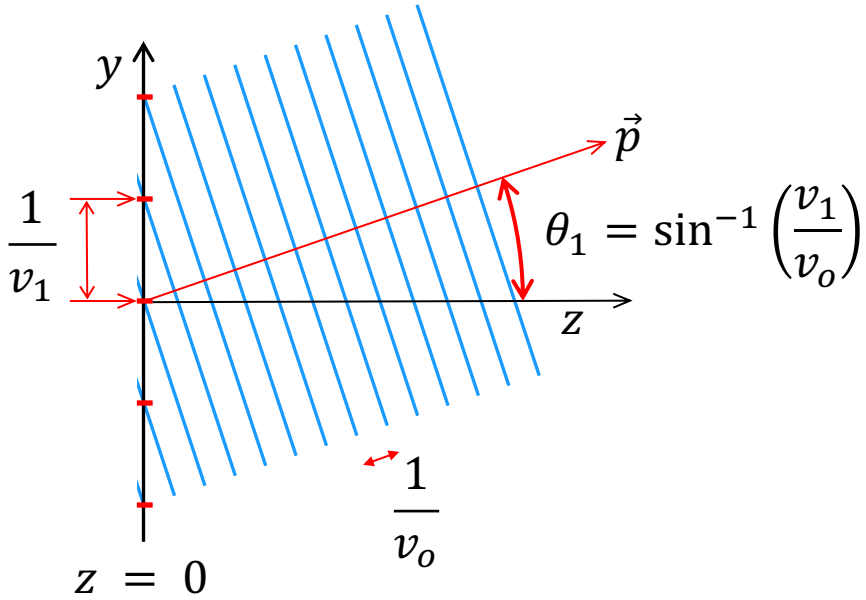
$$\theta = \sin^{-1} \left(\frac{v_y}{v_o} \right)$$

Beam Forming: One transmitted signal

- Excite a signal at $z = 0$ of the form:

$$u_0(y) = U_1 \exp[j2\pi(y v_1)]$$

- This “forms a beam” in the direction $\theta_1 = \sin^{-1} \left(\frac{v_1}{v_o} \right)$

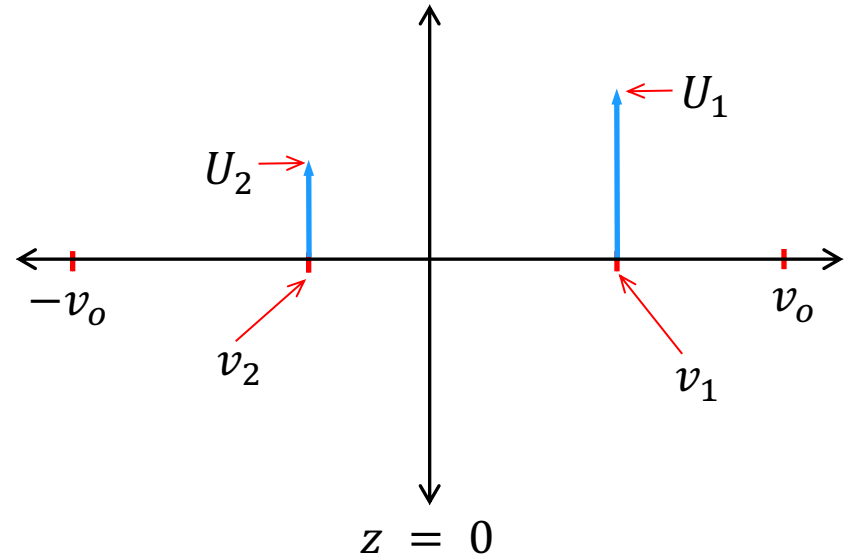
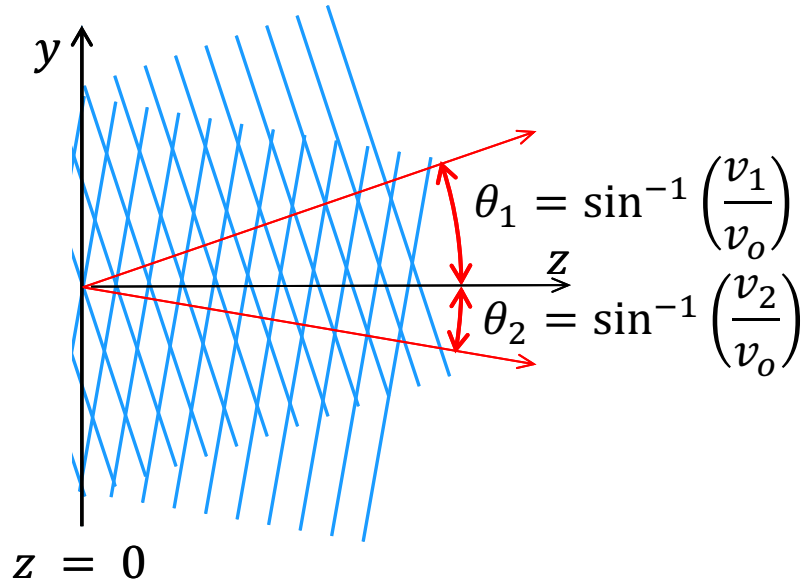


Beam Forming: Two transmitted signals

- Excite a signal at $z = 0$ of the form:

$$u_0(y) = U_1 \exp[j2\pi(y v_1)] + U_2 \exp[j2\pi(y v_2)]$$

- This “forms a beam” in the directions $\theta_1 = \sin^{-1} \left(\frac{v_1}{v_o} \right)$ and $\theta_2 = \sin^{-1} \left(\frac{v_2}{v_o} \right)$



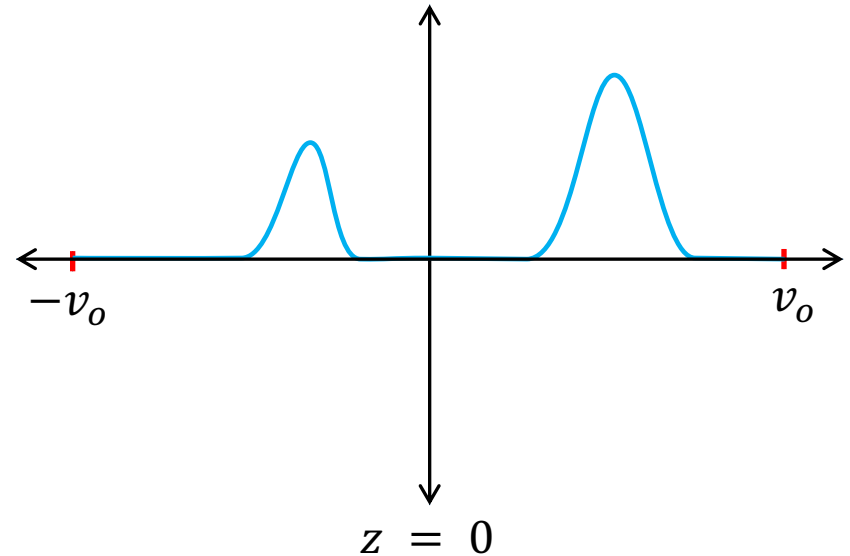
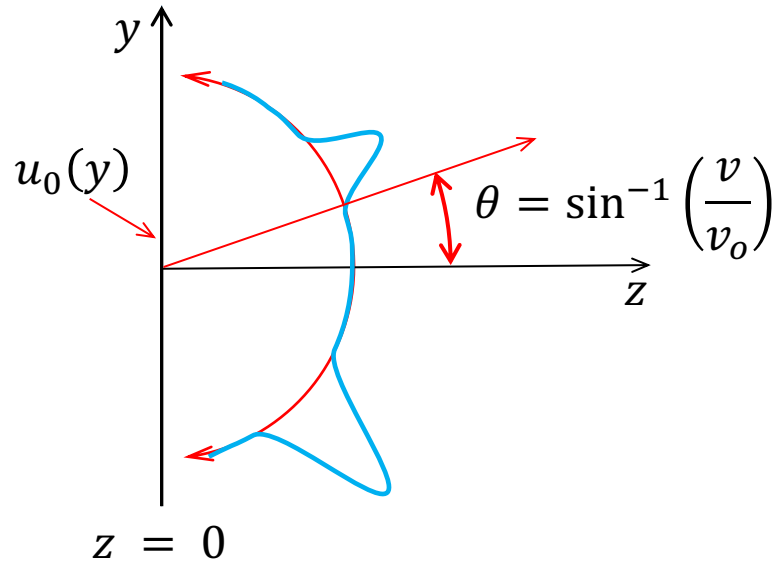
Beam Forming: Arbitrary transmitted signals

- Excite a signal at $z = 0$ of the form:

$$u_0(y) = \int_{\Re} U_0(v) \exp[j2\pi(y v)] dv$$

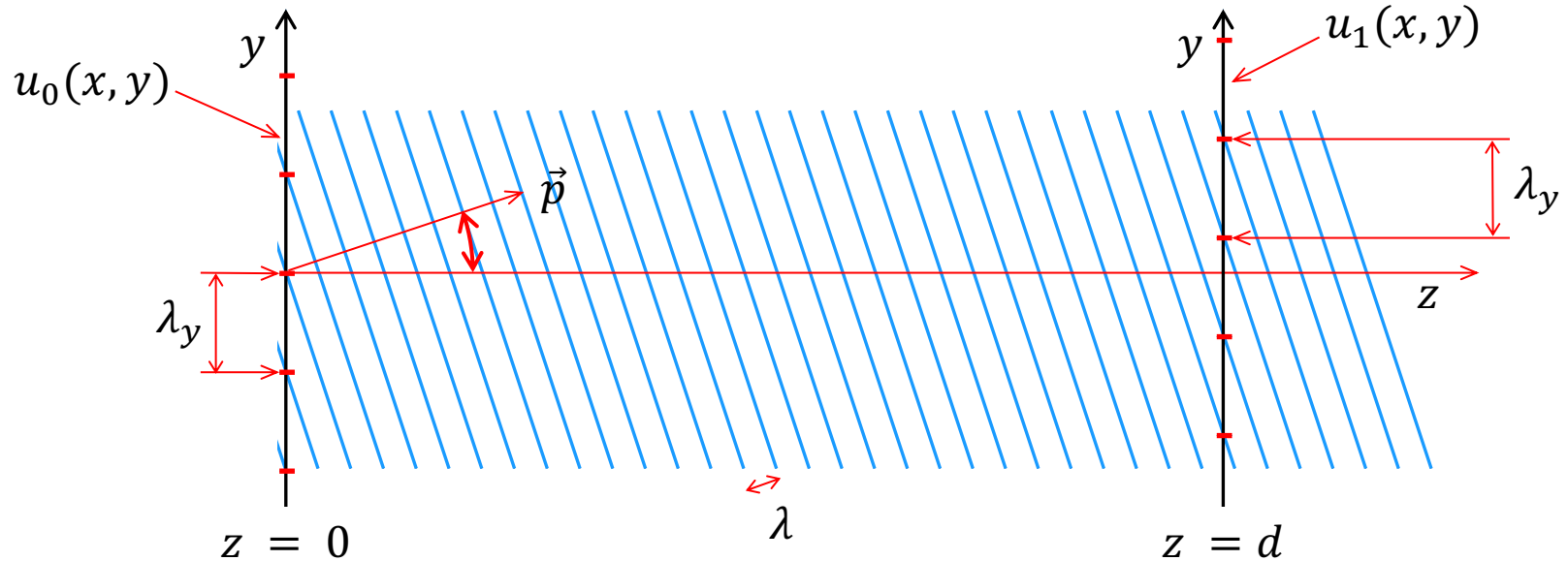
- where $U_0 = \mathcal{F}[u_0]$ is the Fourier transform of u_0 .

- This “forms a beam” in all directions given by



Propagation between Planes

- For a wave of the form: $E(x, y, z) = \exp[j2\pi(v_x x + v_y y + v_z z)]$
- Wave on the $z = 0$ plane is: $u_0(x, y) = \exp[j2\pi(v_x x + v_y y)]$
- Wave on the $z = d$ plane is: $u_d(x, y) = \exp[j2\pi(v_x x + v_y y + v_z d)]$



Phase Modulation between Planes

- For a wave of the form:

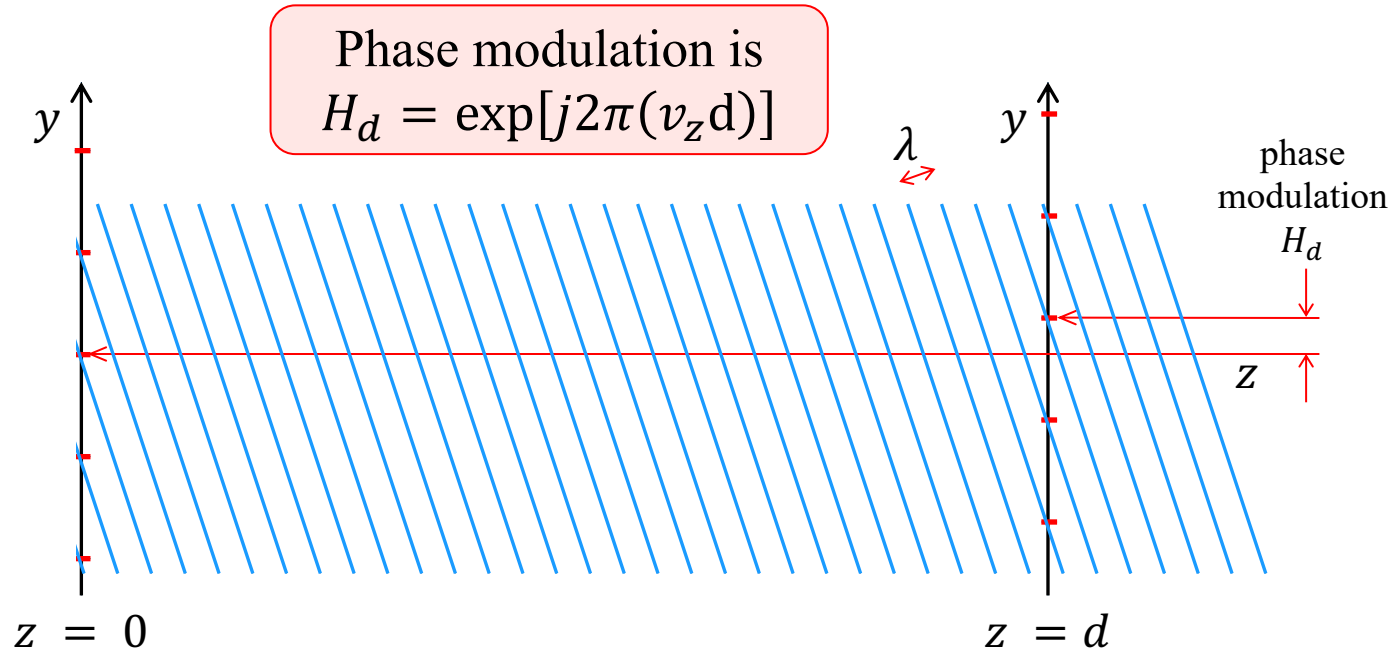
$$E(x, y, z) = \exp[j2\pi(v_x x + v_y y + v_z z)]$$

- Wave on the $z = 0$ plane of:

$$u_0(x, y) = \exp[j2\pi(v_x x + v_y y)]$$

- Wave on the $z = d$ plane of:

$$\begin{aligned} u_d(x, y) &= \exp[j2\pi(v_x x + v_y y)] \exp[j2\pi(v_z d)] \\ &= \exp[j2\pi(v_z d)] u_0(x, y) = H_d u_0(x, y) \end{aligned}$$



Deriving the Fresnel Phase Modulation Function

- So then we have that

$$\begin{aligned}\exp[j2\pi v_z d] &= \exp\left[j2\pi(v_o^2 - [v_x^2 + v_y^2])^{\frac{1}{2}}d\right] \\ &= \exp\left[j2\pi d v_o \left(1 - \frac{v_x^2 + v_y^2}{v_o^2}\right)^{\frac{1}{2}}\right] \\ &\approx \exp\left[j2\pi d v_o \left(1 - \frac{v_x^2 + v_y^2}{2v_o^2}\right)\right] \\ &= \exp[j2\pi d v_o] \exp\left[-j2\pi d v_o \left(\frac{v_x^2 + v_y^2}{2v_o^2}\right)\right] \\ &= \exp[j2\pi d v_o] \exp\left[-j2\pi d \left(\frac{v_x^2 + v_y^2}{2v_o}\right)\right]\end{aligned}$$

Useful Facts:

$$v_o^2 = v_x^2 + v_y^2 + v_z^2$$

$$\text{For } |\beta| \ll 1, \sqrt{1 - \beta} \approx 1 - \frac{\beta}{2}$$

Assumes that:

$$\frac{v_x^2 + v_y^2}{v_o^2} \ll 1$$

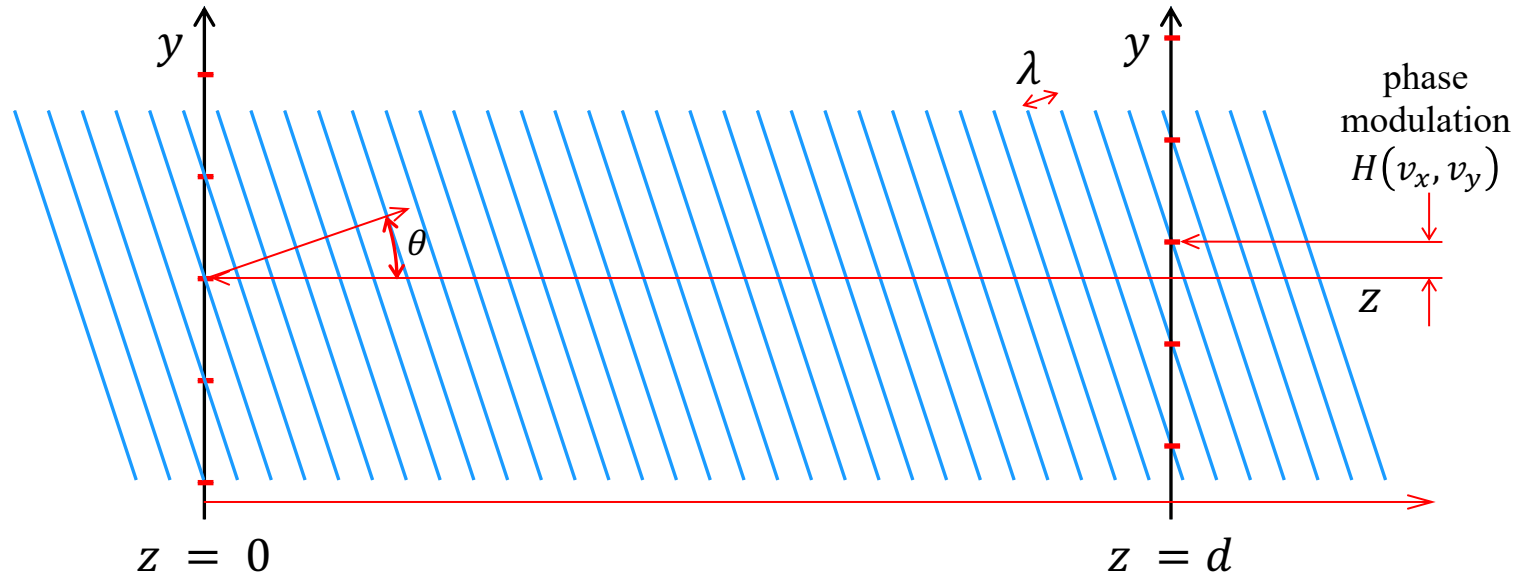
$$\sin^2(\theta) \ll 1$$

Interpretation of Fresnel Phase Modulation

- The transfer function is given by

$$H(v_x, v_y) = \exp[j2\pi d v_o] \exp \left[-j2\pi d v_o \left(\frac{v_x^2 + v_y^2}{2v_o^2} \right) \right]$$

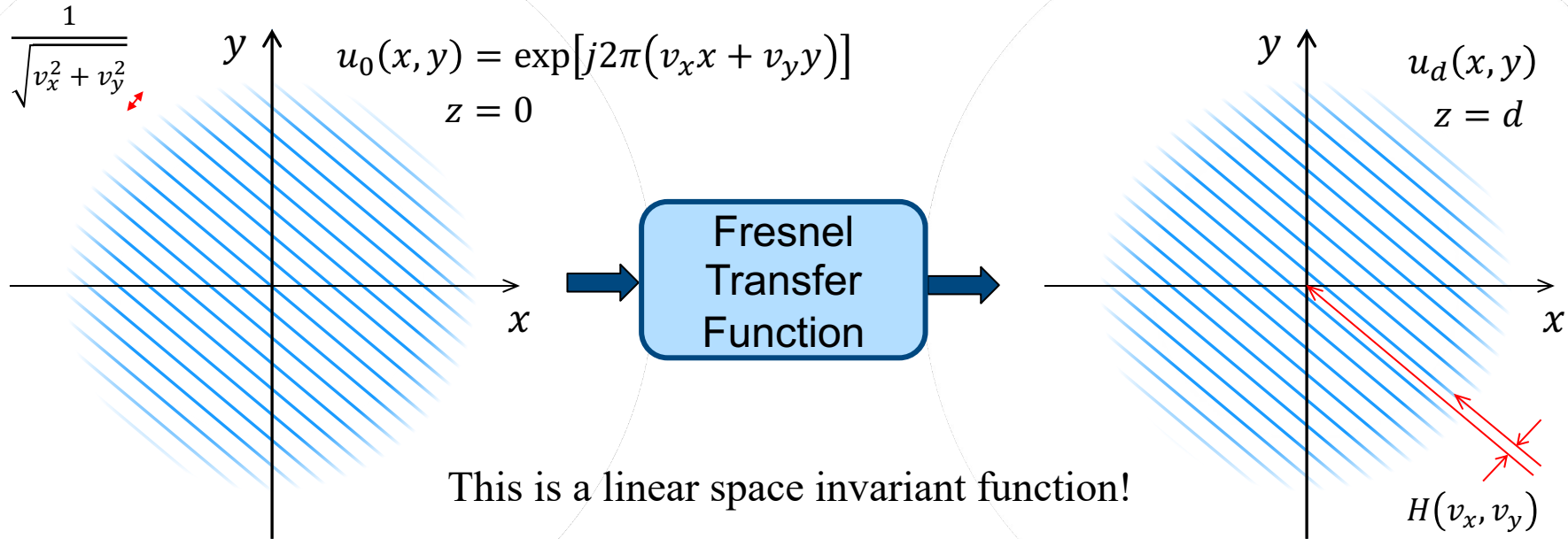
- Pure phase shift; no amplitude; \Rightarrow conserves energy
- Phase shift depends on angle $\frac{1}{2} \sin^2 \theta$ where $\frac{\sqrt{v_x^2 + v_y^2}}{v_o} = \sin \theta = \sin(\text{tilt angle})$ in 3D.



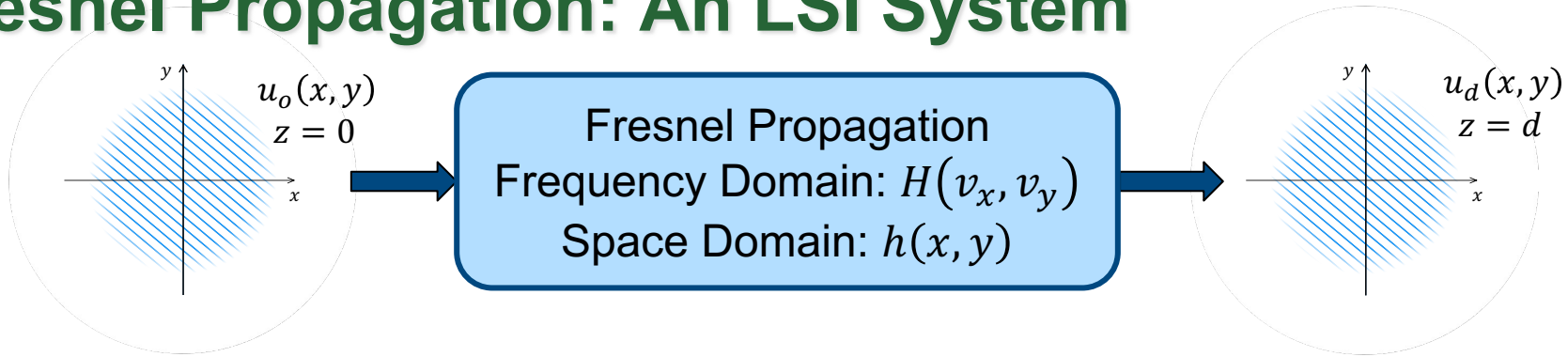
3D Interpretation of Fresnel Phase Modulation

- We can visualize the plane wave on the two planes.

$$u_d(x, y) = H(v_x, v_y) \exp[j2\pi(v_x x + v_y y)]$$



Fresnel Propagation: An LSI System



- Frequency domain transfer function:

$$H(v_x, v_y) = \exp[j2\pi d v_o] \exp \left[-j2\pi d v_o \left(\frac{v_x^2 + v_y^2}{2v_o^2} \right) \right]$$

- Space domain point-spread function (psf):

$$h(x, y) = h_o \exp \left[j\pi \left(\frac{x^2 + y^2}{\lambda d} \right) \right]$$

(from a table
of transform pairs)

$$h_o = \frac{1}{j\lambda d} \exp \left[j2\pi \frac{d}{\lambda} \right]$$

$$\lambda = \frac{1}{v_o}$$

Fresnel Transfer Function

Assumes that:

$$\frac{v_x^2 + v_y^2}{v_o^2} \ll 1$$

$$\sin^2(\theta) \ll 1$$

- In the frequency domain:

$$U_d(v_x, v_y) = H(v_x, v_y) U_o(v_x, v_y)$$

- where

$$H(v_x, v_y) = \exp[j2\pi d v_o] \exp\left[-j2\pi d v_o \left(\frac{v_x^2 + v_y^2}{2v_o^2}\right)\right]$$

$$U_o(v_x, v_y) = \mathcal{F}[u_o] = \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} u_o(x, y) \exp\{-j2\pi(v_x x + v_y y)\} dx dy$$

$$U_1(v_x, v_y) = \mathcal{F}[u_1] = \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} u_1(x, y) \exp\{-j2\pi(v_x x + v_y y)\} dx dy$$

Fresnel Convolution

- In the space domain:

$$u_d(x, y) = h(x, y) * u_0(x, y)$$

$$u_d(x, y) = \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} u_0(x', y') h(x - x', y - y') dx' dy'$$

- where

$$h(x, y) = h_o \exp \left[j\pi \frac{d}{\lambda} \left(\frac{x^2 + y^2}{d^2} \right) \right]$$

$$h_o = \frac{\exp \left[j2\pi \frac{d}{\lambda} \right]}{j\lambda d}$$

Assumes that:

$$\frac{x^2 + y^2}{d^2} \ll 1$$

$$\sin^2(\theta) \ll 1$$

Fresnel Convolution/PSF Example

- In the space domain:

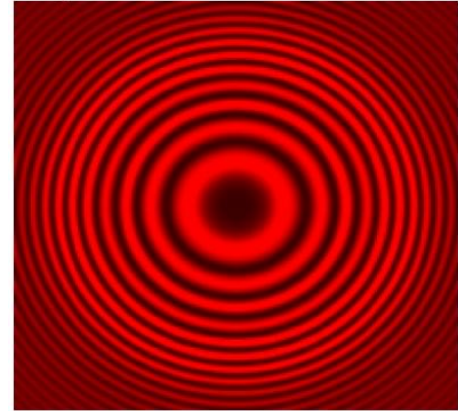
$$u_d(x, y) = h(x, y) * u_0(x, y)$$

- where

$$h(x, y) = h_o \exp \left[j\pi \frac{d}{\lambda} \left(\frac{x^2 + y^2}{d^2} \right) \right]$$

$$h_o = \frac{\exp \left[j2\pi \frac{d}{\lambda} \right]}{j\lambda d}$$

Fresnel diffraction of circular aperture"



$$h(x, y) * \text{circ}(x, y)$$

Fresnel Transformation

Derivation

$$\begin{aligned}u_d(x, y) &= \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} u_0(x', y') h(x - x', y - y') dx' dy' \\&= h_o \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} u_0(x', y') \exp \left\{ j\pi \frac{(x - x')^2 + (y - y')^2}{\lambda d} \right\} dx' dy' \\&= h_o \exp \left\{ j\pi \frac{x^2 + y^2}{\lambda d} \right\} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} u_0(x', y') \exp \left\{ j\pi \frac{x'^2 + y'^2}{\lambda d} \right\} \exp \left\{ -j2\pi \frac{xx' + yy'}{\lambda d} \right\} dx' dy' \\&= h_o \exp \left\{ j\pi \frac{x^2 + y^2}{\lambda d} \right\} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \tilde{u}_0(x', y') \exp \left\{ -j2\pi \frac{xx' + yy'}{\lambda d} \right\} dx' dy' \\&= h_o \exp \left\{ j\pi \frac{x^2 + y^2}{\lambda d} \right\} \tilde{U}_0 \left(\frac{x}{\lambda d}, \frac{y}{\lambda d} \right)\end{aligned}$$

Assumes that:

$$\frac{x^2 + y^2}{d^2} \ll 1$$

$$\sin^2(\theta) \ll 1$$

Fresnel Transformation

Derivation

$$\tilde{u}_0(x, y) = \exp\left\{j\pi \frac{x^2 + y^2}{\lambda d}\right\} u_0(x, y)$$

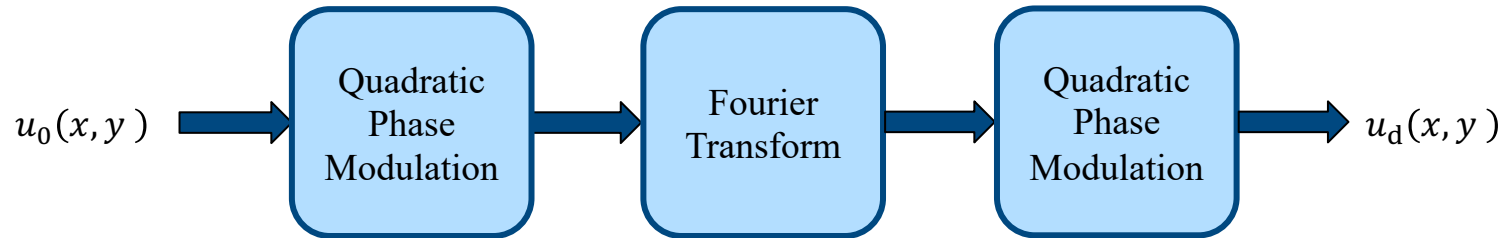
$$\tilde{U}_0(v_x, v_y) = \mathcal{F}[\tilde{u}_0(x, y)]$$

$$u_d(x, y) = h_o \exp\left\{j\pi \frac{x^2 + y^2}{\lambda d}\right\} \tilde{U}_0\left(\frac{x}{\lambda d}, \frac{y}{\lambda d}\right)$$

Assumes that:

$$\frac{x^2 + y^2}{d^2} \ll 1$$

$$\sin^2(\theta) \ll 1$$



Fraunhofer Propagation

Assumes that:

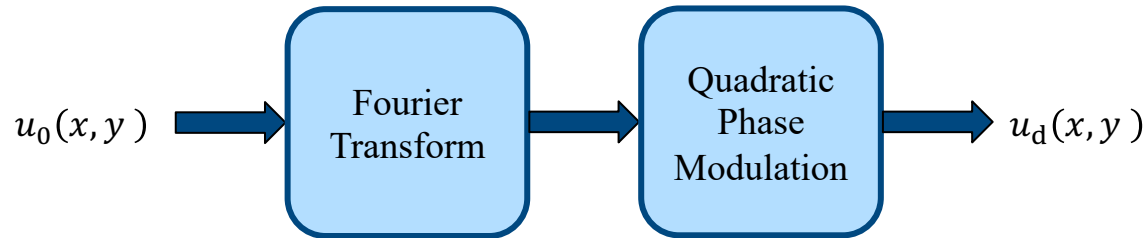
If $u_0 = 0$ for $\|[x, y]\| > a$, then $\frac{a^2}{\lambda d} \ll 1$

- Assume aperture is small, then $\tilde{u}_0(x, y) \approx u_0(x, y)$ and

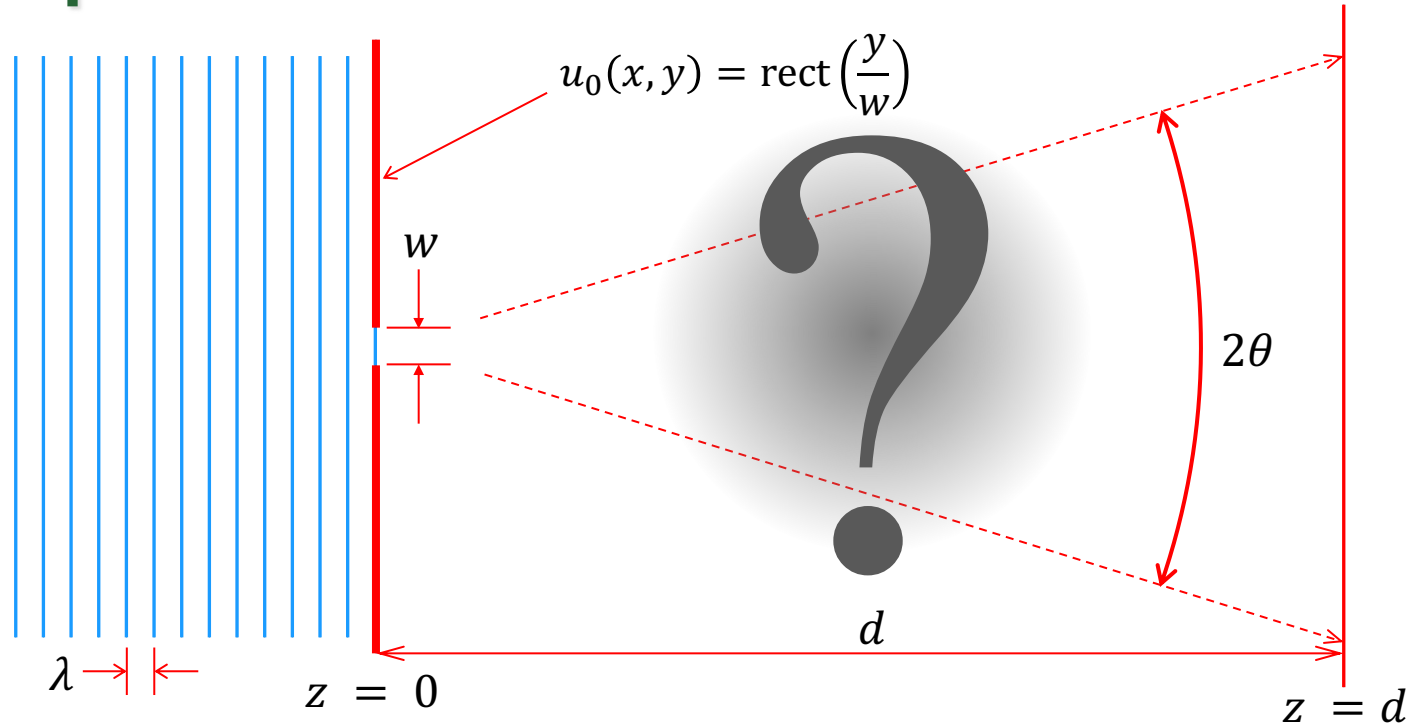
$$u_d(x, y) = h_o \exp\left\{j\pi \frac{x^2 + y^2}{\lambda d}\right\} U_0\left(\frac{x}{\lambda d}, \frac{y}{\lambda d}\right)$$

– where

$$U_0(v_x, v_y) = \mathcal{F}[u_0(x, y)] = \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} u_0(x, y) \exp\{-j2\pi(v_x x + v_y y)\} dx dy$$



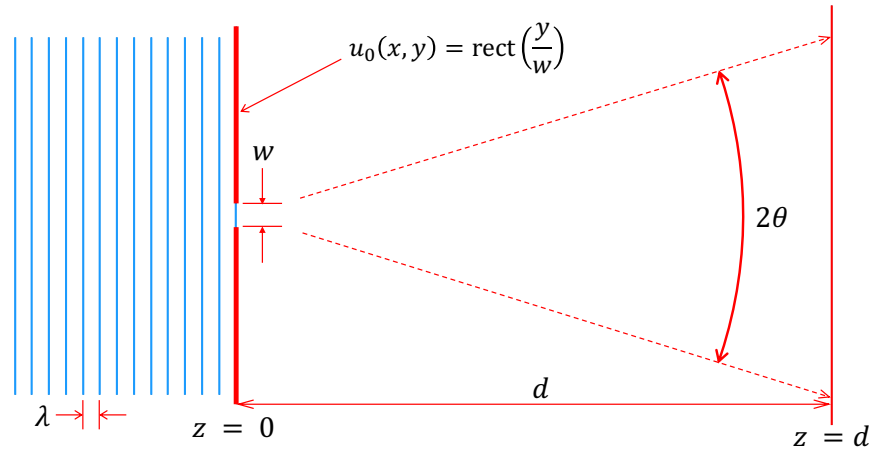
Slit Experiment



■ Questions:

- How does wave propagate?
- What is $u_d(x, y) = ?$

Slit Experiment: Which approximation holds?



■ Fresnel:

- Holds if $\sin^2(\theta) \ll 1$
- If $\sin^2(\theta) = 0.05$, then $|\theta| \leq 13^\circ$

■ Fraunhofer:

- Holds if $w^2/\lambda d \ll 1$; or equivalently, $d \gg w \left(\frac{w}{\lambda}\right)$
- If $\lambda = 1\mu\text{m}$ and $w = 1\text{mm}$, then $d \gg 1\text{m}$
- If $\lambda = 1\mu\text{m}$ and $w = 1\text{cm}$, then $d \gg 100\text{m}$

2D Thin Slit Analysis

■ Which approximation:

- $d \gg w \left(\frac{w}{\lambda}\right) = w \Rightarrow \underline{\text{Yes!}} \Rightarrow \text{Fraunhofer}$

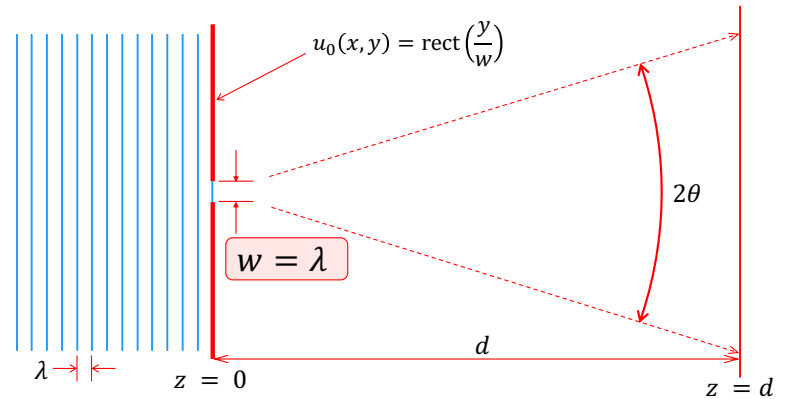
■ Fraunhofer Analysis

- Field at slit: $u_0(y) = \text{rect}\left(\frac{y}{w}\right)$
- Field at distance d :

$$\begin{aligned}u_d(y) &= h_o \exp\left\{j\pi \frac{y^2}{\lambda d}\right\} U_0\left(\frac{y}{\lambda d}\right) \\&= h_o \exp\left\{j\pi \frac{y^2}{\lambda d}\right\} w \text{sinc}\left(\frac{wy}{\lambda d}\right) \\&= wh_o \exp\left\{j\pi \frac{y^2}{\lambda d}\right\} \text{sinc}\left(\frac{y}{d}\right)\end{aligned}$$

- Field at distance d :

$$\begin{aligned}\|u_d(y)\|^2 &= \left\|U_0\left(\frac{y}{\lambda d}\right)\right\|^2 \\&= w^2 \text{sinc}^2\left(\frac{y}{d}\right)\end{aligned}$$

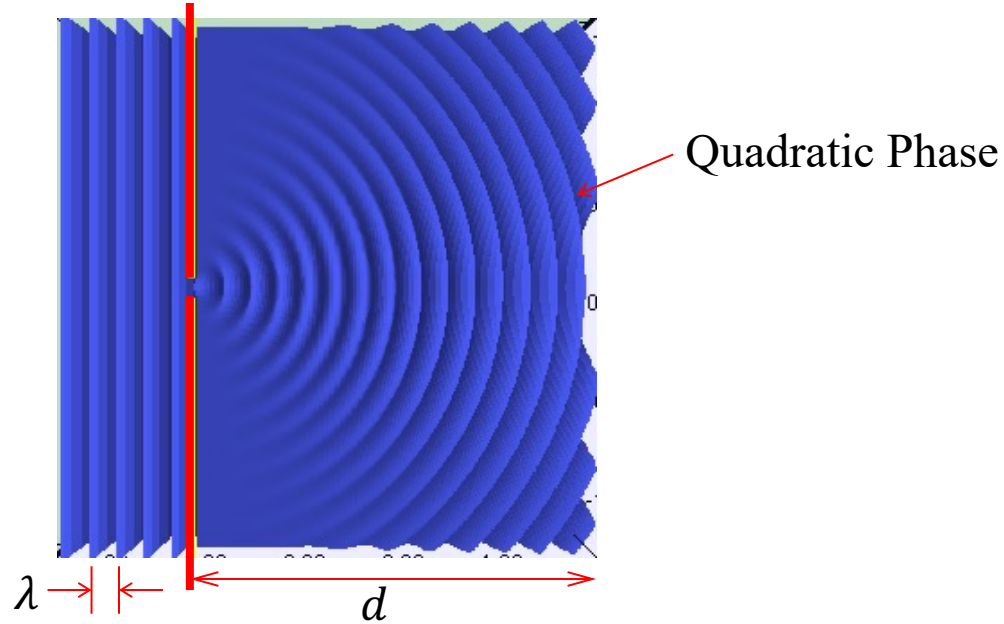


Thin Slit Result

Example:

Thin slit

$$\frac{w}{\lambda} = 1; \frac{d}{\lambda} = 15$$



▪ Field at slit:

$$u_0(x, y) = \text{rect}\left(\frac{y}{w}\right)$$

▪ Field at distance d :

$$u_d(x, y) = wh_0 \exp\left\{j\pi \frac{y^2}{\lambda d}\right\} \text{sinc}\left(\frac{y}{d}\right)$$

▪ Power at distance d :

$$\|u_d(x, y)\|^2 = w^2 \text{sinc}^2\left(\frac{y}{d}\right)$$

2D Wider Slit Analysis

Which approximation:

- $d \gg w \left(\frac{w}{\lambda}\right) = 4w \Rightarrow \underline{\text{Yes!}} \Rightarrow \text{Fraunhofer}$

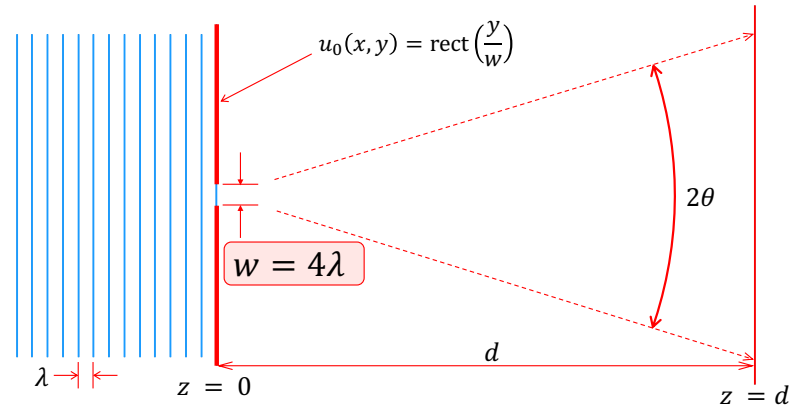
Fraunhofer Analysis

- Field at slit: $u_0(y) = \text{rect}\left(\frac{y}{w}\right)$
- Field at distance d :

$$\begin{aligned}u_d(y) &= h_o \exp\left\{j\pi \frac{y^2}{\lambda d}\right\} U_0\left(\frac{y}{\lambda d}\right) \\&= h_o \exp\left\{j\pi \frac{y^2}{\lambda d}\right\} w \text{sinc}\left(\frac{wy}{\lambda d}\right) \\&= wh_o \exp\left\{j\pi \frac{y^2}{\lambda d}\right\} \text{sinc}\left(4 \frac{y}{d}\right)\end{aligned}$$

- Field at distance d :

$$\begin{aligned}\|u_d(y)\|^2 &= \left\|U_0\left(\frac{y}{\lambda d}\right)\right\|^2 \\&= w^2 \text{sinc}^2\left(\frac{wy}{\lambda d}\right) = w^2 \text{sinc}^2\left(4 \frac{y}{d}\right)\end{aligned}$$



First null occurs at:

$$\frac{wy}{\lambda d} = 1 \Rightarrow \frac{y}{d} = \frac{\lambda}{w}$$

$$y_{\text{null}} = d \left(\frac{\lambda}{w}\right)$$

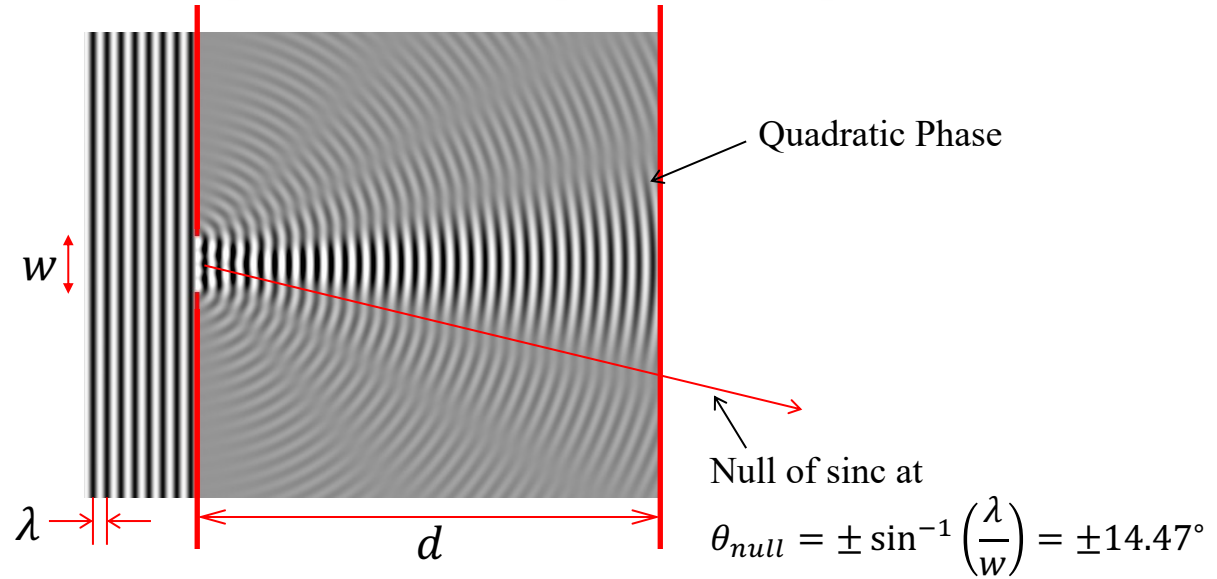
$$\theta_{\text{null}} = \sin^{-1}\left(\frac{\lambda}{w}\right)$$

Fresnel Propagation: Light through a large slit

Example:

Thin slit

$$\frac{w}{\lambda} = 4; \frac{d}{\lambda} = 32$$



- Field at slit: $u_0(x, y) = \text{rect} \left(\frac{y}{w} \right)$
- Field at distance d : $u_d(x, y) = w h_0 \exp \left\{ j\pi \frac{x^2}{\lambda d} \right\} \text{sinc} \left(\frac{yw}{\lambda d} \right)$
- Power at distance d : $\|u_d(x, y)\|^2 = w^2 \text{sinc}^2 \left(\frac{yw}{\lambda d} \right)$