

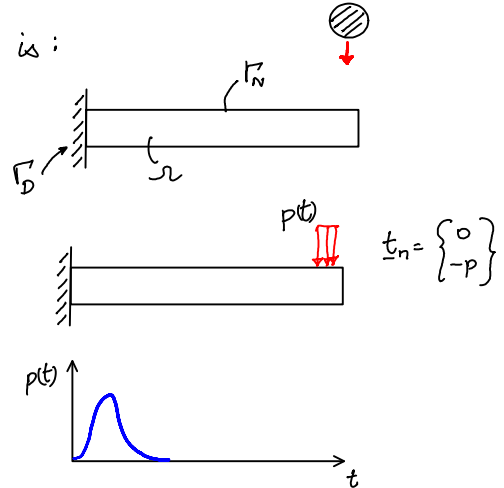
Dynamics

For dynamic problems the equation of motion is:

$$\text{div } \underline{\underline{\sigma}} + \underline{\underline{b}} = \rho \underline{\underline{u}} \ddot{\underline{\underline{u}}} \quad \text{in } \Omega$$

$$\underline{\underline{u}}(\underline{\underline{x}}, t) = \underline{\underline{u}}_0(\underline{\underline{x}}, t) \quad \text{on } \Gamma_D$$

$$\underline{\underline{t}}_n(\underline{\underline{x}}, t) = \underline{\underline{t}}_0(\underline{\underline{x}}, t) \quad \text{on } \Gamma_N$$



Find $\underline{\underline{u}}(\underline{\underline{x}}, t)$

Weak form: Method of weighted residuals

$$G(\underline{\underline{u}}, \underline{\underline{u}}) \equiv \int_{\Omega} \underline{\underline{u}} (\text{div } \underline{\underline{\sigma}} + \underline{\underline{b}} - \rho \underline{\underline{u}} \ddot{\underline{\underline{u}}}) d\Omega$$

Integration by parts:

$$G(\underline{\underline{u}}, \underline{\underline{u}}) = - \left[\int_{\Omega} (\underline{\underline{u}} \cdot \rho \underline{\underline{u}} \ddot{\underline{\underline{u}}} + \nabla \underline{\underline{u}} : \underline{\underline{\sigma}} - \underline{\underline{u}} \cdot \underline{\underline{b}}) d\Omega - \int_{\Gamma_N} \underline{\underline{u}} \underline{\underline{t}}_n d\Gamma \right]$$

Discretization + Finite Element Approximation:

$$\int_{\Omega} \cdot = \sum_{e=1}^M \int_{\Omega^e} \cdot$$

$$\underline{\underline{u}}(\underline{\underline{x}}, t) \approx \underline{\underline{N}}(\underline{\underline{x}}) \underline{\underline{d}}(t) \Rightarrow \ddot{\underline{\underline{u}}} \approx \underline{\underline{N}} \ddot{\underline{\underline{d}}}$$

$$\underline{\underline{\epsilon}} \approx \underline{\underline{B}} \underline{\underline{d}}$$

$$\underline{\underline{\sigma}} \approx \underline{\underline{D}} \underline{\underline{B}} \underline{\underline{d}}$$

$$\tilde{G}^h(\underline{\underline{d}}, \underline{\underline{d}}) =$$

$$- \sum_{e=1}^M \underline{\underline{d}}^e \left[\underbrace{\left(\int_{\Omega^e} \underline{\underline{N}}^T \underline{\underline{N}} d\Omega \right)}_{\underline{\underline{M}}^e} \ddot{\underline{\underline{d}}}^e + \underbrace{\left(\int_{\Omega^e} \underline{\underline{B}}^T \underline{\underline{D}} \underline{\underline{B}} d\Omega \right)}_{\underline{\underline{K}}^e} \underline{\underline{d}}^e - \underbrace{\left(\int_{\Omega^e} \underline{\underline{N}}^T \underline{\underline{b}} d\Omega - \int_{\Gamma_N^e} \underline{\underline{N}}^T \underline{\underline{t}}_n d\Gamma \right)}_{\underline{\underline{f}}^e} \right] = 0$$

$$\text{Assembly } \underline{\underline{A}} \Rightarrow \tilde{G}^h(\underline{\underline{d}}, \underline{\underline{d}}) = - \underline{\underline{d}}^g \left[\underline{\underline{M}}^g \ddot{\underline{\underline{d}}}^g(t) + \underline{\underline{K}}^g \underline{\underline{d}}^g(t) - \underline{\underline{f}}^g(t) \right] = 0 \quad \forall \underline{\underline{d}}$$

This gives us the semi-discrete equations of motion.

This is a 2nd order system of ordinary differential equations (ODEs).

Sometimes artificial physical damping is also included:

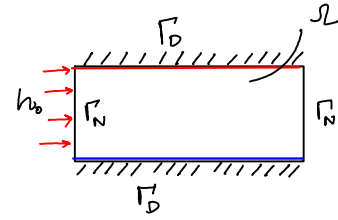
$$\underline{\underline{M}}^g \ddot{\underline{\underline{d}}}^g + \underline{\underline{C}}^g \dot{\underline{\underline{d}}}^g + \underline{\underline{K}}^g \underline{\underline{d}}^g - \underline{\underline{f}}^g = \underline{\underline{0}}$$

where $\underline{\underline{C}} = \alpha \underline{\underline{M}} + \beta \underline{\underline{K}}$. This is called Rayleigh Damping.

Transient Heat Conduction

Governing Equation :

$$\left. \begin{aligned} \rho c \dot{\theta} + \text{div}(\underline{q}) &= f \\ \underline{q} &= -\underline{\kappa} \nabla \theta \end{aligned} \right\} \text{ in } \Omega$$



BCs:

$$\begin{aligned} \theta(\underline{x}, t) &= \theta_0(\underline{x}, t) && \text{on } \Gamma_D \\ \underline{q}(\underline{x}, t) \cdot \underline{n} &= h_0(\underline{x}, t) && \text{on } \Gamma_N \end{aligned}$$

Weak Form + Discretization + FE approximation

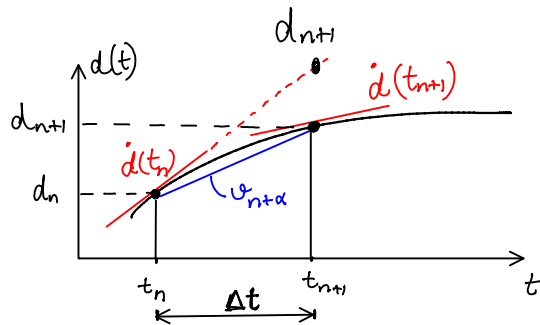
$$\Rightarrow \tilde{G}^h(\underline{d}, \underline{d}) = -\underline{d}^T \left[\underline{M}^g \dot{\underline{d}}^g(t) + \underline{K}^g \underline{d}^g(t) - \underline{f}^g \right] = 0 \quad \forall \underline{d}$$

$$\begin{aligned} \text{where } \underline{M}^g &= \sum_{e=1}^M \underline{M}^e && \text{and } \underline{M}^e = \int_{\Omega^e} \rho c \underline{N}^T \underline{N} \, d\Omega \\ \underline{K}^g &= \sum_{e=1}^M \underline{K}^e && \text{and } \underline{K}^e = \int_{\Omega^e} \underline{B}^T \underline{\kappa} \underline{B} \, d\Omega \end{aligned}$$

Time Integration:

For 1st order systems:

$$\begin{aligned} \underline{d}_n &\approx \underline{d}(t_n) \\ \underline{v}_n &\approx \dot{\underline{d}}(t_n) \end{aligned}$$



Euler's Method:

$$\begin{aligned} \underline{v}_{n+\alpha} &= (1-\alpha) \underline{v}_n + \alpha \underline{v}_{n+1} && - (1) \\ \underline{d}_{n+1} &= \underline{d}_n + \Delta t \underline{v}_{n+\alpha} && - (2) \end{aligned}$$

$\alpha = 0 \Rightarrow$ Forward Euler (Explicit)
 $\alpha = 1 \Rightarrow$ Backward Euler (Implicit)
 $\alpha = 1/2 \Rightarrow$ Mid point rule
 Trapezoidal "
 Crank-Nicolson "
 (Implicit)

Enforce governing equation at t_{n+1} :

$$\underline{M} \underline{v}_{n+1} + \underline{K} \underline{d}_{n+1} - \underline{f}_{n+1} = 0$$

Substitute (1) & (2) :-

$$\left[\underline{M} + \alpha \Delta t \underline{K} \right] \underline{v}_{n+1} = \left\{ \underline{f}_{n+1} - \underline{K} \left(\underline{d}_n + \Delta t (1-\alpha) \underline{v}_n \right) \right\}$$

Solve for $\underline{v}_{n+1} \Rightarrow$ find \underline{d}_{n+1} from (1) & (2)

Repeat for next time step.

Time Integration for Structural Dynamics

To integrate 2nd order ODEs:

$$\underline{\underline{M}} \ddot{\underline{d}} + \underline{\underline{C}} \dot{\underline{d}} + \underline{\underline{K}} \underline{d} - \underline{f} = \underline{0}$$

We can recast them as:

(system of 1st order ODEs)
(However, size is doubled)

$$\underline{\underline{M}} \dot{\underline{v}} + \underline{\underline{C}} \underline{v} + \underline{\underline{K}} \underline{d} - \underline{f} = \underline{0}$$

$\underline{\underline{d}} = \underline{v}$

Alternatively, use Newmark Family of Integrators:

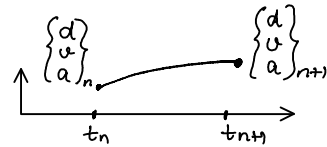
Approximate $\underline{a}_n \approx \ddot{d}(t_n)$ (Parameters: γ, β)
 $\underline{v}_n \approx \dot{d}(t_n)$
 $\underline{d}_n \approx d(t_n)$

Method:

$$\underline{v}_{n+1} = \underline{v}_n + \Delta t \left[(1-\gamma) \underline{a}_n + \gamma \underline{a}_{n+1} \right] \quad \text{--- (1)}$$

$$\underline{d}_{n+1} = \underline{d}_n + \Delta t \underline{v}_n + \Delta t^2 \left[\left(\frac{1}{2} - \beta\right) \underline{a}_n + \beta \underline{a}_{n+1} \right] \quad \text{--- (2)}$$

Enforce equations of motion at t_{n+1} :



$$\left[\underline{\underline{M}} + \gamma \Delta t \underline{\underline{C}} + \frac{1}{2} \beta \Delta t^2 \underline{\underline{K}} \right] \underline{a}_{n+1} = \left\{ \begin{array}{l} \underline{f}_{n+1} - \left[\underline{v}_n + \Delta t (1-\gamma) \underline{a}_n \right] \underline{\underline{C}} \\ - \left[\underline{d}_n + \Delta t \underline{v}_n + \Delta t^2 \left(\frac{1}{2} - \beta\right) \underline{a}_n \right] \underline{\underline{K}} \end{array} \right\}$$

Solve for \underline{a}_{n+1} and then \underline{v}_{n+1} & \underline{d}_{n+1} from (1) & (2).

Common Methods:

γ	β	Method	Type	Stability	Accuracy
$1/2$	0	Central-difference	Explicit	$\Delta t < \frac{2}{\omega_{max}}$	$O(\Delta t^2)$
$1/2$	$1/4$	Trapezoidal Rule (Const. Avg Accel)	Implicit	<u>Unconditional</u>	$O(\Delta t^2)$
$1/2$	$1/6$	Linear Avg. Acceleration	Implicit	$\Delta t < \frac{2\sqrt{3}}{\omega_{max}}$	$O(\Delta t^2)$

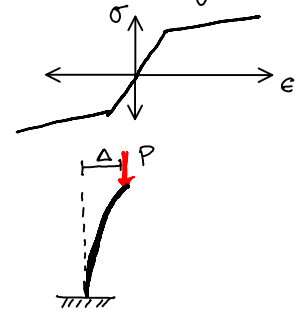
Non-Linear Problems

For structural / solid mechanics problems, nonlinearity arises from

- Material \rightarrow eg. Plasticity

- Geometric (P- Δ) \rightarrow

$$\begin{cases} \tilde{F} = \frac{\partial \tilde{x}}{\partial \underline{x}} & ; \tilde{C} = \tilde{F}^T \tilde{F} \\ \tilde{E} = \frac{1}{2} (\tilde{C} - \tilde{I}) = \frac{1}{2} (\nabla \underline{u} + \nabla \underline{u}^T + \nabla \underline{u}^T \nabla \underline{u}) \end{cases}$$

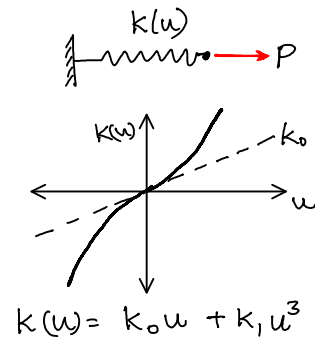
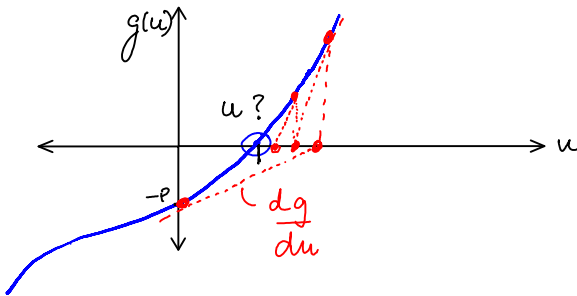


(or both)

Newton's Method for Numerical Solution of Non-linear Problems:

First, consider a Single-Degree-of-freedom (SDOF)

$$g(u) = k(u) - P = 0 \quad (\text{find } u)$$



Steps: 1) Assume a solution $u = u^0$

- 2) Find $g(u)$
- 3) If $g(u) \stackrel{*}{=} 0 \rightarrow$ return u
 - Else
 - Solve: $g(u + \Delta u) = 0$ for Δu

$$g(u + \Delta u) \approx g(u) + \underbrace{\left[\frac{dg}{du} \right]_u}_{K_T \text{ (Tangent)}} \Delta u = 0$$

$$[K_T] \Delta u = \{-g\}$$

$$\left\{ K_T = \frac{dk}{du} = k_0 + 3k_1 u^2 \right.$$
 - $u = u + \Delta u$
 - End
- 4) Repeat 2 & 3.

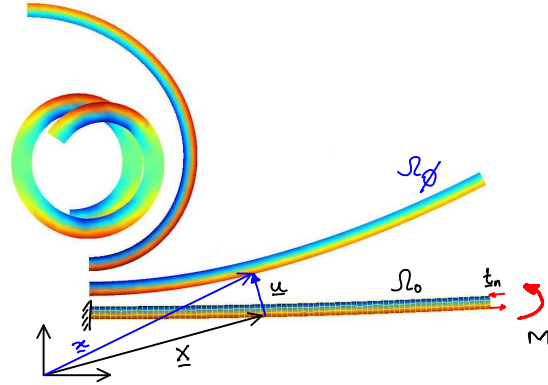
For System of Non-linear Equations

Strong form:

$$\frac{\partial}{\partial x} \underline{\underline{\sigma}} + \underline{\underline{b}} = \underline{\underline{0}} \quad \text{on } \Omega_\phi$$

or

$$\frac{\partial}{\partial X} \underline{\underline{P}} + \underline{\underline{B}} = \underline{\underline{0}} \quad \text{on } \Omega_0$$



Weak form:

$$G(\underline{u}, \underline{\bar{u}}) = \int_{\Omega_\phi} \underline{\bar{u}} (\text{div } \underline{\underline{\sigma}} + \underline{\underline{b}}) d\Omega_\phi$$

$$= \int_{\Omega_0} \underline{\bar{u}} (\text{div } \underline{\underline{\sigma}} + \underline{\underline{b}}) |J| d\Omega_0$$

$$= \int_{\Omega_0} \underline{\bar{u}} (\text{DIV } \underline{\underline{P}} + \underline{\underline{B}}) d\Omega_0$$

$$\underline{\underline{E}} = \underline{\underline{I}} + \nabla \underline{u} \quad ; \quad |J| = \det(\underline{\underline{E}})$$

$$\underline{\underline{P}} = |J| \underline{\underline{\sigma}} \underline{\underline{E}}^{-T} \quad ; \quad \underline{\underline{B}} = \underline{\underline{b}} / |J|$$

$$\underline{\underline{S}} = |J| \underline{\underline{E}}^{-1} \underline{\underline{\sigma}} \underline{\underline{E}}^{-T}$$

Integration by parts:

$$G(\underline{u}, \underline{\bar{u}}) = - \left[\int_{\Omega_0} \underbrace{(\underline{\nabla} \underline{\bar{u}} : \underline{\underline{P}})}_{(\underline{\underline{E}} : \underline{\underline{S}})} d\Omega_0 - \int_{\Omega_0} \underline{\bar{u}} \cdot \underline{\underline{B}} d\Omega_0 - \int_{\Gamma_{N_0}} \underline{\bar{u}} \cdot \underline{\underline{T}}_N d\Gamma \right]$$

Discretization + FE approximation:

$$\underline{u} = \underline{N}(\underline{x}) \underline{d} \quad ; \quad \underline{\underline{E}} = \underline{\underline{B}}(\underline{x}) \underline{d}$$

$$\tilde{G}^h(\underline{d}, \underline{\bar{d}}) = - \sum_{e=1}^M \underline{\bar{d}}^{eT} \left[\int_{\Omega_0^e} \underline{\underline{B}}^T \underline{\underline{S}} d\Omega_0 - \int_{\Omega_0^e} \underline{N}^T \underline{\underline{B}} d\Omega_0 - \int_{\Gamma_{N_0}^e} \underline{N}^T \underline{\underline{T}}_N d\Gamma \right]$$

↑
Voigt

$$\tilde{G}^h(\underline{d}, \underline{\bar{d}}) = - \underline{\bar{d}}^{eT} \left(\underline{f}^{int}(\underline{d}^e) - \underline{f}^{ext}(\underline{d}^e) \right) = 0 \quad \forall \underline{\bar{d}}^e$$

Thus to solve for \underline{d}^e :

$$\underline{g}(\underline{d}^e) = \underline{f}^{int}(\underline{d}^e) - \underline{f}^{ext}(\underline{d}^e) = 0$$

For Newton's method we need "tangent":

$$\underline{\underline{K}}_T^e = \frac{\partial \underline{g}(\underline{d})}{\partial \underline{d}} = \sum_{e=1}^M \underline{\underline{A}} \underline{\underline{K}}_T$$

$$\underline{\underline{K}}_T = \int_{\Omega_0^e} \underline{\underline{B}}^T \left(\frac{\partial \underline{\underline{S}}}{\partial \underline{\underline{E}}} \right) \underline{\underline{B}} d\Omega_0 + \int_{\Omega_0^e} \frac{\partial \underline{N}^T}{\partial \underline{x}} \underline{\underline{S}} \frac{\partial \underline{N}}{\partial \underline{x}} d\Omega_0 + \int_{\Omega_0^e} \underline{N}^T \left(\frac{\partial \underline{\underline{B}}}{\partial \underline{d}} \right) d\Omega_0 + \int_{\Gamma_{N_0}^e} \underline{N}^T \left(\frac{\partial \underline{\underline{T}}_N}{\partial \underline{d}} \right) d\Gamma$$

$$\underline{\underline{K}}_T = \underbrace{\underline{\underline{K}}_M}_{\text{Material}} + \underbrace{\underline{\underline{K}}_G}_{\text{Geometric}} + \underbrace{\underline{\underline{K}}_L}_{\text{Load}}$$

Newton's Method:

→ Load Steps

1) Assume a solution $\underline{d}_i^g = \underline{d}_0^g$

2) Calculate $\underline{g}(\underline{d}_i^g)$

3) If $\|\underline{g}(\underline{d}_i^g)\| < \text{tol}$ → return \underline{d}_i^g

Else

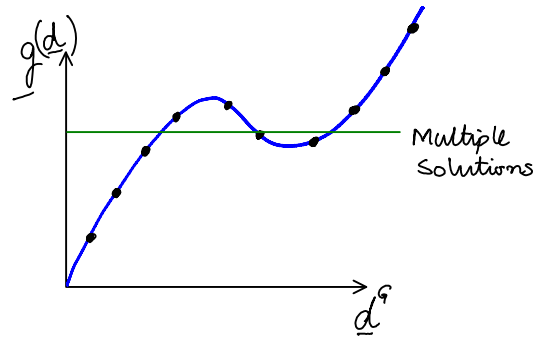
• solve: $[\underline{K}_T^g] \underline{\Delta d} = -\underline{g}(\underline{d}_i^g)$

• $\underline{d}_{i+1}^g = \underline{d}_i^g + \underline{\Delta d}$

End

4) Repeat 2 & 3

Repeat 1) thru 4).



In Class Assignment:

$$K(u) = K_0 u + K_1 u^3 \quad ; \quad K_0 = 1 \quad K_1 = 1$$

$$g(u) = K(u) - P = 0 \quad ; \quad P = 1 \quad \Rightarrow u = ?$$

(i) Assumption $u^0 = \frac{P}{K_0} = \frac{1}{1} = 1$

$$\Rightarrow g(u^0) = 1 + 1 - 1 = 1$$

Find u^1 ?

$$K_T = \frac{\partial g}{\partial u} = K_0 + 3K_1 u^2 = 4$$

$$\Rightarrow K_T \Delta u = -g \Rightarrow \Delta u = \frac{-1}{4}$$

$$u^1 = u^0 + \Delta u = 1 - \frac{1}{4} = \frac{3}{4}$$