

- Reflection and Transmission
 - o Oblique Incidence
 - Perpendicular Pol, Parallel Pol
 - Brewster Angle, Critical Angle

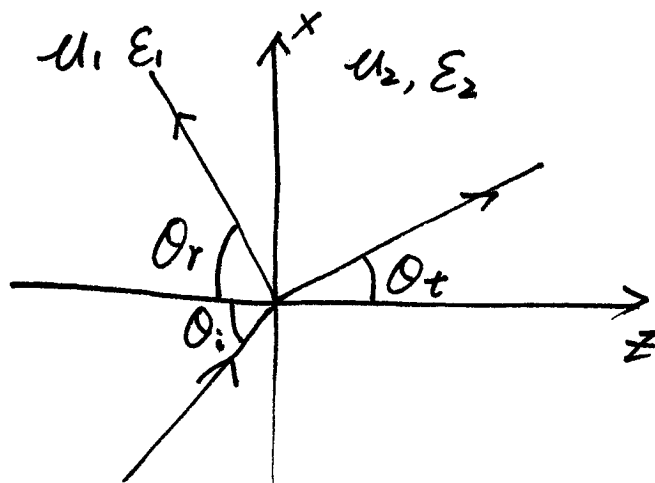
Interested Readers:

- Textbook 6.9, Balanis 5.3.1, 5.3.2
- Textbook 6.12-6.13, Balanis 5.3.3, 5.3.4

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2. Oblique Incidence

[A] Perpendicular polarization or E polarization



$$\vec{E}^i = \hat{y} E_0 e^{-j\vec{\beta}^i \cdot \vec{r}} = \hat{y} E_0 e^{-j\beta_1 (x \sin \theta_i + z \cos \theta_i)}$$

$$\vec{E}^r = \hat{y} R_{\perp} E_0 e^{-j\vec{\beta}^r \cdot \vec{r}} = \hat{y} R_{\perp} E_0 e^{-j\beta_1 (x \sin \theta_r - z \cos \theta_r)}$$

$$\vec{E}^t = \hat{y} T_{\perp} E_0 e^{-j\vec{\beta}^t \cdot \vec{r}} = \hat{y} T_{\perp} E_0 e^{-j\beta_2 (x \sin \theta_t + z \cos \theta_t)}$$

$$\nabla \times \vec{E} = -j\omega\mu \vec{H} \Rightarrow$$

$$\vec{H}^i = (-\hat{x} \cos \theta_i + \hat{z} \sin \theta_i) \frac{E_0}{\eta_1} e^{-j\beta_1 (x \sin \theta_i + z \cos \theta_i)}$$

$$\vec{H}^r = (\hat{x} \cos \theta_r + \hat{z} \sin \theta_r) \frac{R_{\perp} E_0}{\eta_1} e^{-j\beta_1 (x \sin \theta_r - z \cos \theta_r)}$$

$$\vec{H}^t = (-\hat{x} \cos \theta_t + \hat{z} \sin \theta_t) \frac{T_{\perp} E_0}{\eta_2} e^{-j\beta_2 (x \sin \theta_t + z \cos \theta_t)}$$

$$\hat{n} \times (\vec{E}_1 - \vec{E}_2) = 0 \quad \text{or} \quad E_{1tan} = E_{2tan} \Rightarrow$$

$$E_0 e^{-i\beta_1 \sin\theta_i x} + R_\perp E_0 e^{-i\beta_1 \sin\theta_r x}$$

$$= T_\perp E_0 e^{-i\beta_2 \sin\theta_t x}$$

$$\hat{n} \times (\vec{H}_1 - \vec{H}_2) = 0, \quad \text{or} \quad H_{1tan} = H_{2tan} \Rightarrow$$

$$-\cos\theta_i \frac{E_0}{\eta_1} e^{-i\beta_1 \sin\theta_i x} + \cos\theta_r \frac{R_\perp E_0}{\eta_1} e^{-i\beta_1 \sin\theta_r x}$$

$$= -\cos\theta_t \frac{T_\perp E_0}{\eta_2} e^{-i\beta_2 \sin\theta_t x}$$

Phase matching:

$$\boxed{\beta_1 \sin\theta_i = \beta_1 \sin\theta_r = \beta_2 \sin\theta_t}$$

$$\sin\theta_i = \sin\theta_r \Rightarrow \boxed{\theta_i = \theta_r} \quad \underline{\underline{\text{Snell's law of reflection}}}$$

$$\beta_1 \sin\theta_i = \beta_2 \sin\theta_t \Rightarrow \boxed{\frac{\sin\theta_i}{\sin\theta_t} = \sqrt{\frac{\epsilon_2 \epsilon_1}{\epsilon_1 \epsilon_2}}}$$

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$$1 + R_L = T_L$$

$$-\cos \theta_i \frac{1}{\eta_1} + \cos \theta_r \frac{R_L}{\eta_1} = -\cos \theta_t \frac{T_L}{\eta_2}$$

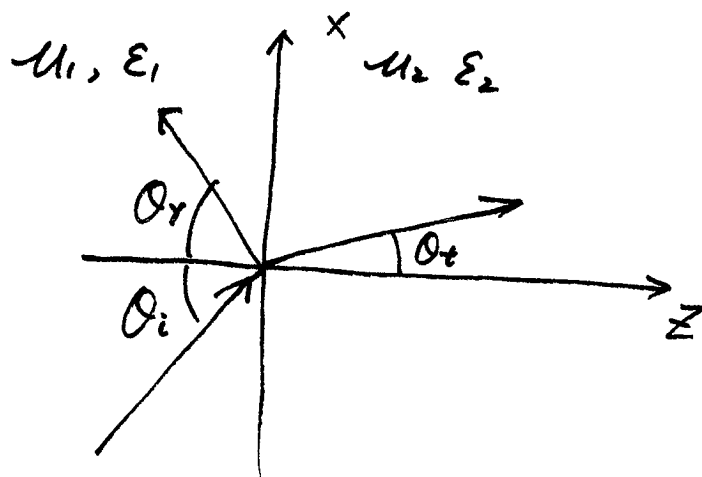
Solution:

$$R_L = \frac{\eta_2 \cos \theta_i - \eta_1 \cos \theta_t}{\eta_2 \cos \theta_i + \eta_1 \cos \theta_t}$$
$$T_L = \frac{2\eta_2 \cos \theta_i}{\eta_2 \cos \theta_i + \eta_1 \cos \theta_t}$$

$$R_L = \frac{\frac{\eta_2}{\cos \theta_t} - \frac{\eta_1}{\cos \theta_i}}{\frac{\eta_2}{\cos \theta_t} + \frac{\eta_1}{\cos \theta_i}} = \frac{Z_{22} - Z_{21}}{Z_{22} + Z_{21}}$$

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B Parallel polarization or H polarization



$$\vec{H}^i = \hat{y} e^{-i\beta_1 (x \sin \theta_i + z \cos \theta_i)}$$

$$\vec{H}^r = -\hat{y} R_{11} H_0 e^{-i\beta_1 (x \sin \theta_r - z \cos \theta_r)}$$

$$\vec{H}^t = \hat{y} T_{11} H_0 e^{-i\beta_2 (x \sin \theta_t + z \cos \theta_t)} \cdot \frac{Z_{z1}}{Z_{z2}}$$

Result

$$\theta_r = \theta_i$$

$$\frac{\sin \theta_i}{\sin \theta_t} = \sqrt{\frac{\mu_2 \epsilon_2}{\mu_1 \epsilon_1}}$$

$$R_{11} = \frac{\eta_2 \cos \theta_t - \eta_1 \cos \theta_i}{\eta_2 \cos \theta_t + \eta_1 \cos \theta_i}$$

$$T_{11} = \frac{2\eta_2 \cos \theta_t}{\eta_2 \cos \theta_t + \eta_1 \cos \theta_i}$$

$$R_{11} = \frac{Z_{z2} - Z_{z1}}{Z_{z2} + Z_{z1}} \quad Z_z = \eta \cos \theta$$

$$= \frac{\eta_1 \cos \theta_i}{\eta_2 \cos \theta_t}$$

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Brewster Angle (No reflection)

For (E polarization) $R_{\perp} = \frac{\eta_2 \cos \theta_i - \eta_1 \cos \theta_t}{\eta_2 \cos \theta_i + \eta_1 \cos \theta_t}$

If $\eta_2 \cos \theta_{iB} = \eta_1 \cos \theta_t \Rightarrow R_{\perp} = 0$

$$\sqrt{\frac{\mu_2}{\epsilon_2}} \cos \theta_{iB} = \sqrt{\frac{\mu_1}{\epsilon_1}} \cos \theta_t = \sqrt{\frac{\mu_1}{\epsilon_1}} \sqrt{1 - \sin^2 \theta_t}$$

From Snell's Law:

$$\frac{\sin \theta_{iB}}{\sin \theta_t} = \sqrt{\frac{\mu_2 \epsilon_2}{\mu_1 \epsilon_1}}$$

$$\sin^2 \theta_t = \frac{\mu_1 \epsilon_1}{\mu_2 \epsilon_2} \sin^2 \theta_{iB}$$

Solution: $\sin^2 \theta_{iB} = \frac{\frac{\epsilon_2}{\epsilon_1} - \frac{\mu_2}{\mu_1}}{\frac{\mu_1}{\mu_2} - \frac{\mu_2}{\mu_1}}$

$$\sin \theta_{iB} = \pm \sqrt{\frac{\frac{\epsilon_2}{\epsilon_1} - \frac{\mu_2}{\mu_1}}{\frac{\mu_1}{\mu_2} - \frac{\mu_2}{\mu_1}}}$$

For $\mu_1 = \mu_2$, no total transmission unless $\epsilon_1 = \epsilon_2$

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For H pto polarization:
$$R_{11} = \frac{\eta_2 \cos \theta_t - \eta_1 \cos \theta_i}{\eta_2 \cos \theta_t + \eta_1 \cos \theta_i}$$

If $\eta_2 \cos \theta_t = \eta_1 \cos \theta_{iB}$, $R_{11} = 0$

$$\sin \theta_{iB} = \pm \sqrt{\frac{\epsilon_2/\epsilon_1 - \mu_2/\mu_1}{\epsilon_2/\epsilon_1 - \epsilon_1/\epsilon_2}}$$

If $\mu_1 = \mu_2$, $\sin \theta_{iB} = \pm \sqrt{\frac{\epsilon_2}{\epsilon_1 + \epsilon_2}}$

Critical Angle (Total reflection)

Snell's law of refraction:

$$\frac{\sin \theta_i}{\sin \theta_t} = \sqrt{\frac{\mu_2 \epsilon_2}{\epsilon_1 \mu_1}} \quad \sin \theta_t = \sqrt{\frac{\mu_1 \epsilon_1}{\mu_2 \epsilon_2}} \sin \theta_i$$

$$\cos \theta_t = \sqrt{1 - \sin^2 \theta_t} = \sqrt{1 - \frac{\mu_1 \epsilon_1}{\mu_2 \epsilon_2} \sin^2 \theta_i}$$

When $\sin \theta_i = \sqrt{\frac{\mu_2 \epsilon_2}{\mu_1 \epsilon_1}}$, $\sin \theta_t = 1 \Rightarrow \theta_t = \frac{\pi}{2}$

$$\boxed{\theta_c = \sin^{-1} \sqrt{\frac{\mu_2 \epsilon_2}{\mu_1 \epsilon_1}}} \leftarrow \text{critical angle}$$

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when $\theta_i = \theta_c$:

$$R_{\perp} = 1, \quad T_{\perp} = 2$$

$$\begin{aligned}\vec{E}^t &= \hat{y} T_{\perp} E_0 e^{-i\beta_2(x \sin \theta_t + z \cos \theta_t)} \\ &= \hat{y} 2 E_0 e^{-i\beta_2 x}\end{aligned}$$

$$\begin{aligned}\vec{H}^t &= (-\hat{x} \cos \theta_t + \hat{z} \sin \theta_t) \frac{T_{\perp} E_0}{\eta_2} e^{-i\beta_2(x \sin \theta_t + z \cos \theta_t)} \\ &= \hat{z} \frac{2 E_0}{\eta_2} e^{-i\beta_2 x}\end{aligned}$$

$$\vec{S}_{av}^t \Big|_{\theta_i = \theta_c} = \frac{1}{2} \operatorname{Re} (\vec{E}^t \times \vec{H}^{t*}) = \hat{x} \frac{2|E_0|^2}{\eta_2}$$

$$|\vec{S}_{av}^i|_{\theta_i = \theta_c} = \frac{|E_0|^2}{2\eta_1}$$

$$|\vec{S}_{av}^r|_{\theta_i = \theta_c} = \frac{|E_0|^2}{2\eta_1} \quad \Leftarrow \text{All power is reflected back!}$$

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Point of discussion: Where does the power in medium 2 come from?

What happens when $\theta_i > \theta_c$?

$$\vec{E}^t = \hat{y} T_{\perp} E_0 e^{-i\beta_2 (x \sin \theta_t + z \cos \theta_t)}$$

When $\theta_i > \theta_c$, $\sin \theta_t > 1$

$$\cos \theta_t = \sqrt{1 - \sin^2 \theta_t} = \pm j \sqrt{\sin^2 \theta_t - 1}$$

$$\vec{E}^t = \hat{y} T_{\perp} E_0 e^{-i\beta_2 x \sin \theta_t} e^{-\beta_2 \sqrt{\sin^2 \theta_t - 1} z}$$

Attenuation constant:

$$\alpha_e = \beta_2 \sqrt{\sin^2 \theta_t - 1}$$

Phase velocity: $v_{pe} = \frac{\omega}{\beta_2 \sin \theta_t} < v_{p2}$

Application: Optical fiber

