

# **An Explicit Inverse Based Fast Direct Volume Integral Equation Solver for Electromagnetic Analysis**

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# Outline

- **Introduction**
- **Explicit Inverse Based Fast VIE Solver**
- **Numerical Results**
- **Conclusions**



# Introduction

- **Unlike PDE, IE based methods lead to dense matrices**
- **State-of-the-art *iterative* solutions**
  - Examples: FMM-based, Fast QR-based, FFT-based methods
  - Require iterations for each right hand side
- **State-of-the-art *direct* solutions**
  - Example:  $\mathcal{H}^2$ -based methods
  - Avoid iterations, require numerical inversion/factorization



# Introduction

- Unlike Surface-IE, volume IE based methods offer flexibility of analyzing problems with complicated materials
- Resultant numerical system, however, is complicated by all four possible forms of double integrals
- Contribution of this work
  - Developed an explicit inverse for solving VIE system, which bypasses the need of numerical inversion and factorization
  - Further enhanced the efficiency of the proposed explicit inverse by error controlled  $\mathcal{H}^2$ -matrix-based fast computation



# Explicit Inverse Based Fast VIE Solver

- VIE System of Equations
- Explicit Inverse
- $\mathcal{H}^2$  -Matrix Representation and Its Error Bound
- $\mathcal{H}^2$  -Matrix Partition
- Rank Function
- Compact Storage
- Efficient Matrix-Vector Multiplication



# VIE System

- VIE based formulation

$$\mathbf{S}\mathbf{D} = \mathbf{E} \quad \text{where} \quad \mathbf{S} = (\mathbf{\Lambda} + \mathbf{G})$$

- $\mathbf{\Lambda}$  is sparse while  $\mathbf{G}$  is dense

$$\Lambda_{mn} = \frac{1}{\epsilon_n^+} \int_{T_m^+} \vec{f}_m \cdot \vec{f}_n dv + \frac{1}{\epsilon_n^-} \int_{T_m^-} \vec{f}_m \cdot \vec{f}_n dv ;$$

$$\mathbf{G}_{mn} = -\omega^2 \int_V \vec{f}_m \cdot \vec{A}_n(\vec{r}) dv - \int_V \phi_n(\vec{r}) \nabla \cdot \vec{f}_m dv + \int_S \phi_n(\vec{r}) \vec{f}_m \cdot \hat{n} ds ; \quad \text{where}$$

$$\vec{A}_n(\vec{r}) = \frac{\mu_o a_n}{3} \left( \frac{\kappa_n^+}{V_n^+} \int_{T_n^+} \rho_n^+ \vec{g}(\vec{r}, \vec{r}') dv' + \frac{\kappa_n^-}{V_n^-} \int_{T_n^-} \rho_n^- \vec{g}(\vec{r}, \vec{r}') dv' \right) ; \phi_n(\vec{r}) = \frac{-a_n}{\epsilon_o} \left( \frac{\kappa_n^+}{V_n^+} \int_{T_n^+} \vec{g}(\vec{r}, \vec{r}') dv' - \frac{\kappa_n^-}{V_n^-} \int_{T_n^-} \vec{g}(\vec{r}, \vec{r}') dv' - \frac{(\kappa_n^+ - \kappa_n^-)}{a_n} \int_{a_n} \vec{g}(\vec{r}, \vec{r}') ds' \right)$$

$$G_i = K \int_P f_m(\vec{r}) \int_{Q'} g(\vec{r}, \vec{r}') h(\vec{r}') dq' dp \quad \{P, Q\} \in \{Volume, Surface\}$$



# Explicit Inverse

- A straightforward approach

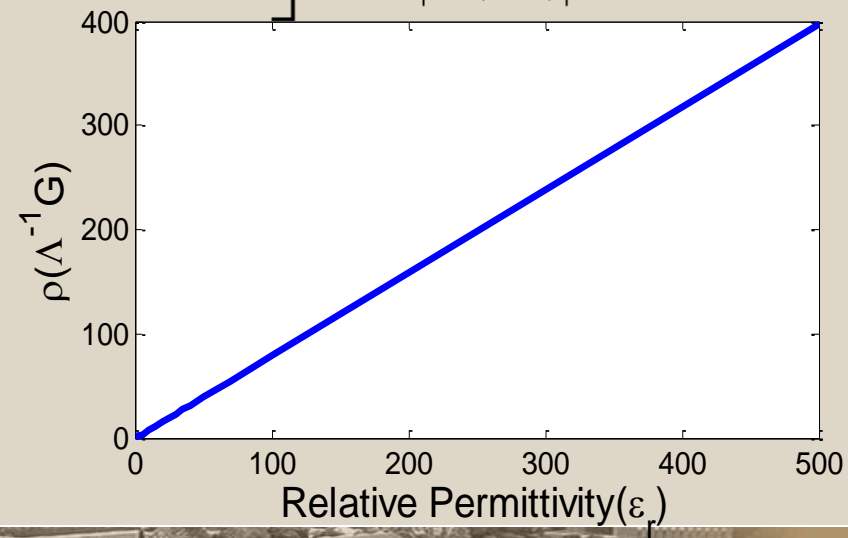
$$\mathbf{S} = \mathbf{\Lambda} \left[ \mathbf{I} + \mathbf{\Lambda}^{-1} \mathbf{G} \right]$$

Then using Newman Series

$$\left[ \mathbf{I} + \mathbf{W} \right]^{-1} = \left[ \mathbf{I} - \mathbf{W} + \mathbf{W}^2 - \mathbf{W}^3 + \mathbf{W}^4 - \mathbf{W}^5 + \mathbf{W}^6 \dots \right] \quad ; \quad |\rho(\mathbf{W})| \leq 1$$

**Problem:**

Resultant spectral radius depends on dielectric constant





# Explicit Inverse

- A much better approach

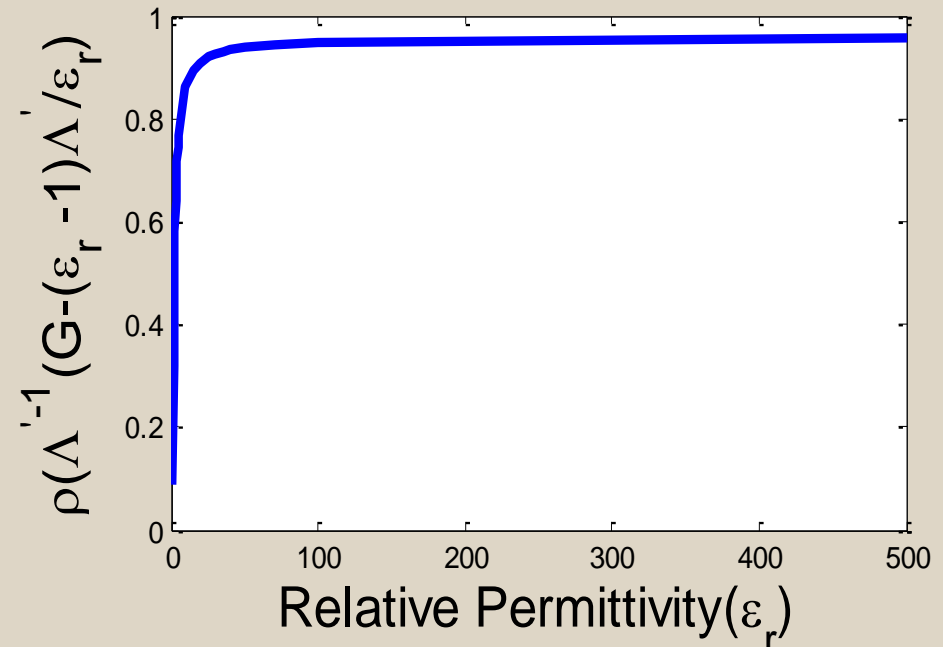
$$\mathbf{S} = \mathbf{\Lambda}' \left[ \mathbf{I} + \mathbf{\Lambda}'^{-1} (\mathbf{G} - \mathbf{Z}) \right]$$

where

$$\mathbf{\Lambda}' = \varepsilon_r \mathbf{\Lambda} \quad \mathbf{Z} = \mathbf{\Lambda}' - \mathbf{\Lambda}$$

**Advantage:**

**Resultant spectral radius has little dependence on dielectric constant.**



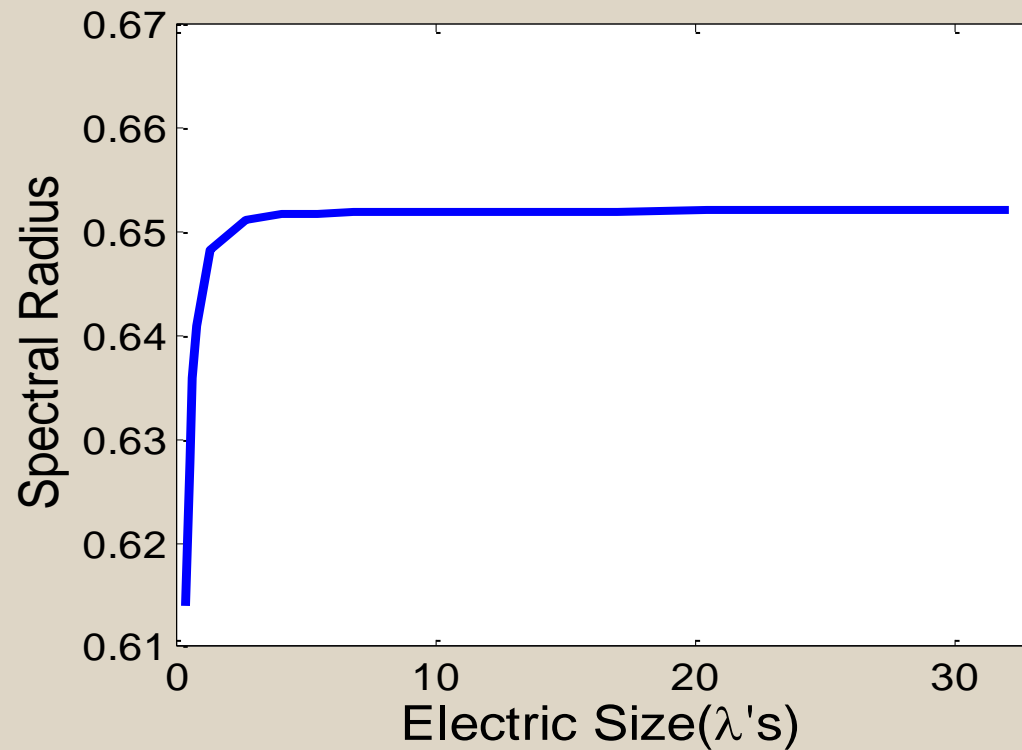
Comparison of Spectral Radius vs. Permittivity ( length =  $1\lambda$ )



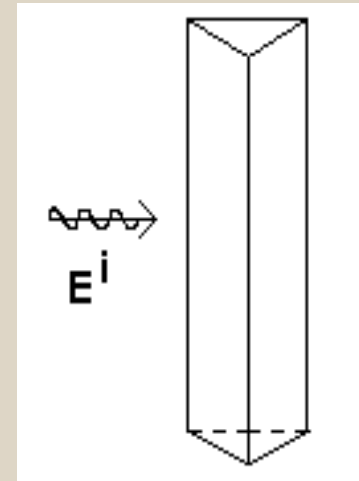


# Explicit Inverse

- Spectral radius vs. electric size for a dielectric rod



Spectral radius vs. electric size



# Explicit Inverse

- **Direct solution based on explicit inverse**

$$\mathbf{D} = \left[ \mathbf{I} + \Lambda'^{-1}(\mathbf{G} - \mathbf{Z}) \right]^{-1} \Lambda'^{-1} \mathbf{E}$$

$$\text{if } \left| \rho \left( \Lambda'^{-1}(\mathbf{G} - \mathbf{Z}) \right) \right| \leq 1$$

$$= \left[ \mathbf{I} - \Lambda'^{-1}(\mathbf{G} - \mathbf{Z}) + \Lambda'^{-1}(\mathbf{G} - \mathbf{Z}) \Lambda'^{-1}(\mathbf{G} - \mathbf{Z}) - \Lambda'^{-1}(\mathbf{G} - \mathbf{Z}) \Lambda'^{-1}(\mathbf{G} - \mathbf{Z}) \Lambda'^{-1}(\mathbf{G} - \mathbf{Z}) + \dots \right] \Lambda'^{-1} \mathbf{E}$$

- The  $(k+1)$ -th term is  $\Lambda'^{-1}(\mathbf{G} - \mathbf{Z})$  times the  $k$ -th term
- Total cost is a few matrix-vector multiplications

$$\Lambda'^{-1}(\mathbf{G} - \mathbf{Z}) \mathbf{f}^{(k)}, \quad k = 1, 2, \dots, p$$



# Explicit Inverse

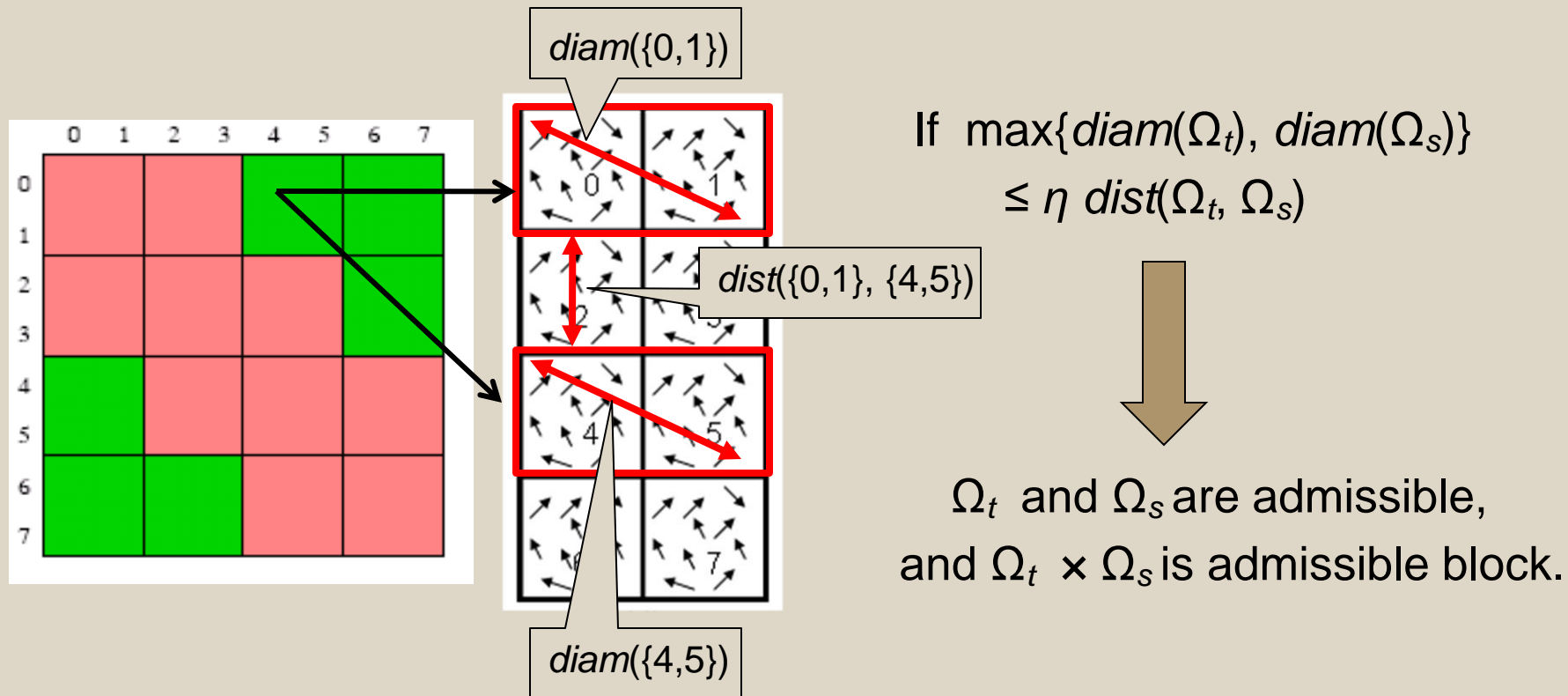
- **Accelerate**  $\Lambda'^{-1}(\mathbf{G}-\mathbf{Z})f^{(k)}$ ,  $k = 1, 2, \dots, p$ 
  - The sparse matrix inverse  $\Lambda'^{-1}$  can be performed in linear complexity by orthogonal bases [1]
  - The dense matrix-vector multiplication is accelerated by  $\mathcal{H}^2$ -matrix-based fast computation, the cost of which is  $O(N)$  from small to tens of wavelengths [2]
  - The overall cost of the proposed explicit inverse is  $O(N)$  from small to tens of wavelengths

[1] D. Jiao and J. M. Jin, "Three-dimensional orthogonal vector basis functions for time-domain finite element solution of vector wave equations," *IEEE Trans. Antennas Propagat.*, vol. 51, no. 1, pp. 59-66, Jan. 2003.

[2] W. Chai and D. Jiao, "An  $\mathcal{H}^2$ -Matrix-Based Integral-Equation Solver of Reduced Complexity and Controlled Accuracy for Solving Electrodynamical Problems," *IEEE Trans. Antennas Propag.*, vol. 57, no. 10, pp. 3147-3159, 2009.

# $\mathcal{H}^2$ -Matrix Representation and Its Error Bound

## Admissibility condition



- [1] W. Hackbusch, B. Khoromskij, and S. Sauter, "On  $\mathcal{H}^2$ -matrices," Lecture on Applied Mathematics, H. Bungartz, R. Hoppe, and C. Zenger, eds., pp. 9-29, 2000.
- [2] W. Chai and D. Jiao, "An  $\mathcal{H}^2$ -Matrix-Based Integral-Equation Solver of Reduced Complexity and Controlled Accuracy for Solving Electrodynamics Problems," *IEEE Trans. Antennas Propag.*, vol. 57, no. 10, pp. 3147-3159, 2009.

# $\mathcal{H}^2$ -Matrix Representation and Its Error Bound

## Degenerate Approximation

If the admissibility condition is satisfied,

$$g(\vec{r}, \vec{r}') = \exp(-jk|\vec{r} - \vec{r}'|) / 4\pi |\vec{r} - \vec{r}'|$$



$$\tilde{g}^{t,s}(\vec{r}, \vec{r}') = \sum_{\nu \in K^t} \sum_{\mu \in K^s} g(\xi_\nu^t, \xi_\mu^s) L_\nu^t(\vec{r}) L_\mu^s(\vec{r}')$$

where,  $K := \{\mathbf{v} \in \mathbb{N}^d : v_i \leq p \text{ for all } i \in \{1, \dots, d\}\} = \{1, \dots, p\}^d$   
 $d=1, 2, 3$ , for 1-, 2-, and 3- $D$  problems respectively  
 $p$  is the number of interpolation points in each dimension  
 $\xi$  is a family of interpolation points  
 $L$  are the corresponding Lagrange polynomials



# $\mathcal{H}^2$ -Matrix Representation and Its Error Bound

$$G_i = K \int_P f_m(\vec{r}) \int_{Q'} g(\vec{r}, \vec{r}') h(\vec{r}') dq' dp$$

$\{P, Q\} \in \{Volume, Surface\}$

$$\tilde{G}_i = K \sum_{v \in K^t} \sum_{\mu \in K^s} g(\xi_v^t, \xi_\mu^s) \int_P f_m(\vec{r}) L_v^t(\vec{r}) dp \int_{Q'} h(\vec{r}') L_\mu^s(\vec{r}') dq'$$

$$\tilde{G}_i := \mathbf{V}^t \mathbf{S}^{t,s} \mathbf{V}^{sT} \quad \text{where,}$$

$\mathbf{V}^t \in \mathbb{R}^{\#t \times \#K^t}$  is the integral over domain  $P$

$\mathbf{V}^s \in \mathbb{R}^{\#s \times \#K^s}$  is the integral over domain  $Q'$

$\mathbf{S}^{t,s} \in \mathbb{C}^{\#K^t \times \#K^s}$  represents the coupling matrix  $K g(\xi_v^t, \xi_\mu^s)$



# $\mathcal{H}^2$ -Matrix Representation and Its Error Bound

- If the admissibility condition is satisfied:

$$\tilde{\mathbf{G}}_i := \mathbf{V}^t \mathbf{S}^{t,s} \mathbf{V}^s{}^T \quad (\mathcal{H}^2 \text{ - matrix representation})$$

$$\begin{aligned} & \| \mathbf{g}(\mathbf{r}, \mathbf{r}') - \tilde{\mathbf{g}}^{(t,s)}(\mathbf{r}, \mathbf{r}') \|_{\infty, Q_t \times Q_s} \\ & \leq \frac{4ed}{\pi} (\Lambda_p)^{2d} p \frac{1}{\text{dist}(Q_t, Q_s)} [1 + \sqrt{2\kappa\eta} \text{dist}(Q_t, Q_s) + \sqrt{2\eta}] \left[ 1 + \frac{\sqrt{2}}{\kappa\eta \text{dist}(Q_t, Q_s) + \eta} \right]^{-p} \end{aligned}$$

where,

$\mathbf{K}$  is the wave number

$\Lambda_p$  is interpolation scheme dependent constant

- If not :

$$\tilde{\mathbf{G}}_i^{(t,s)} := \mathbf{G}_i^{(t,s)} \quad (\text{Full- matrix representation})$$

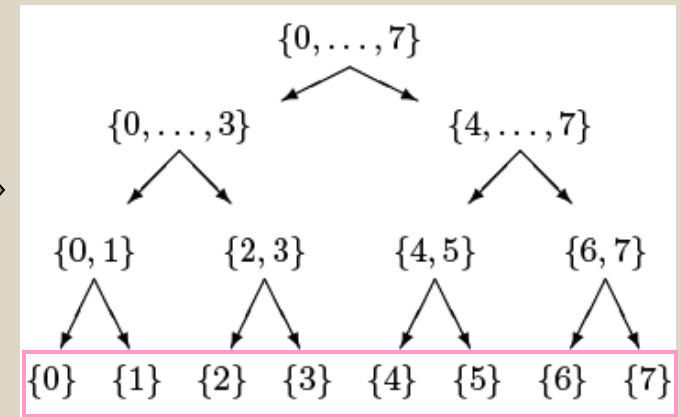
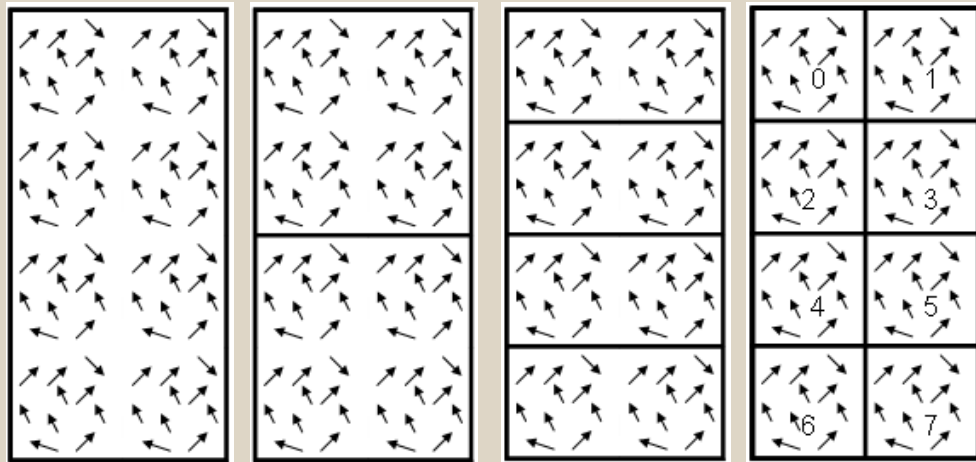




# $\mathcal{H}^2$ -Matrix Partition

## Build Cluster Tree

Geometric partitioning



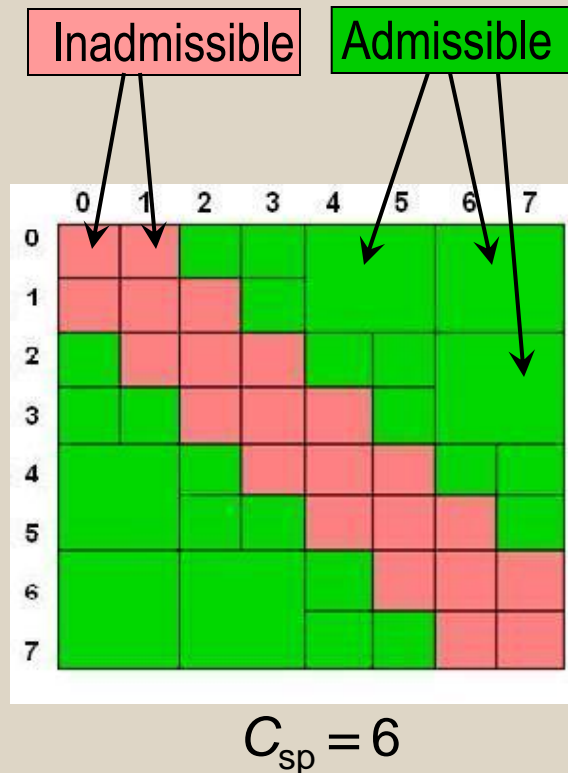
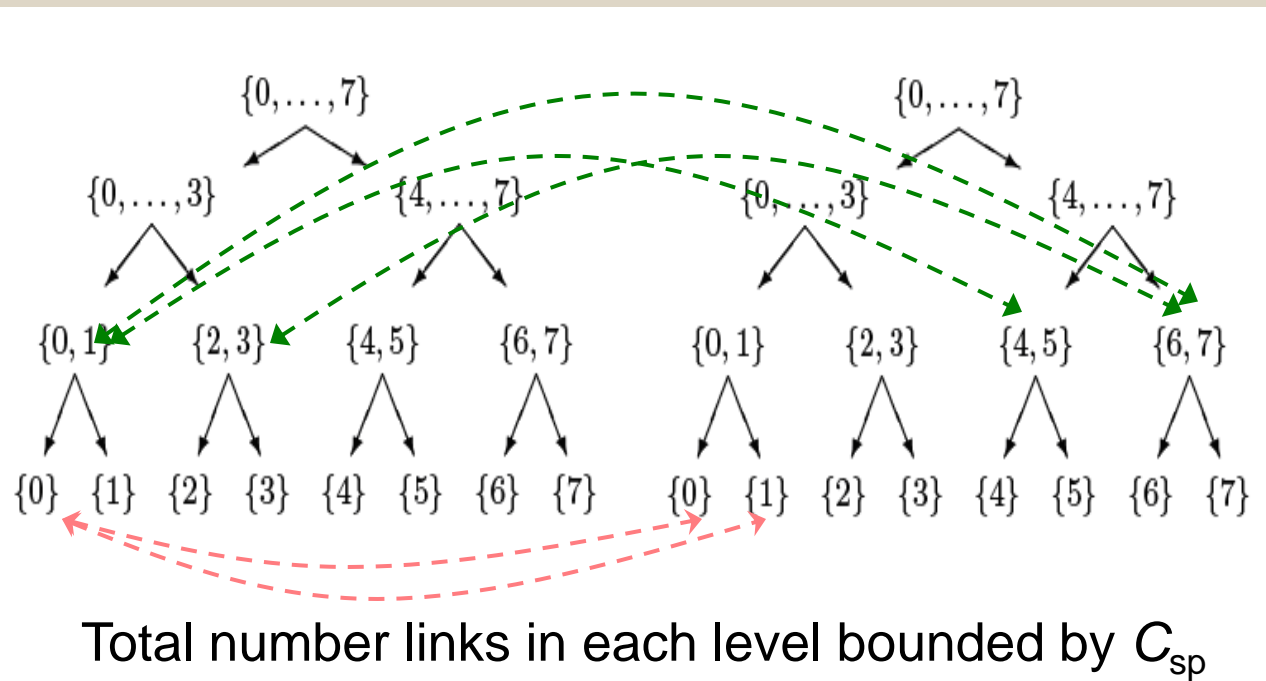
A cluster tree

$$\text{leafsize } (n_{\min}) = 10$$



# $\mathcal{H}^2$ -Matrix Partition

## Build Block Cluster Tree



# Rank Function

$$k_{\text{var}}(b) = p(b)^d$$

$$p(b) = \hat{a} + \hat{b}(L - l(b))$$

where

$$l(b) = \text{level}(t) = \text{level}(s)$$

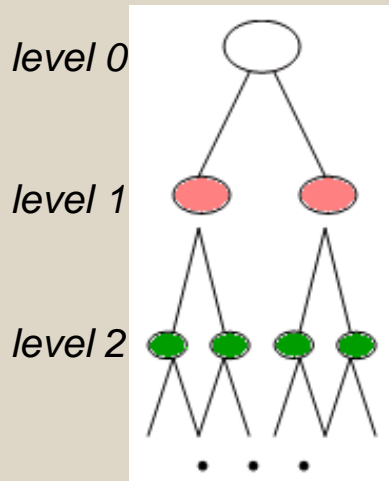
$$p(b) = \hat{a} \quad \text{if } L \leq l(b)$$

$\hat{a}, \hat{b}$  are constants

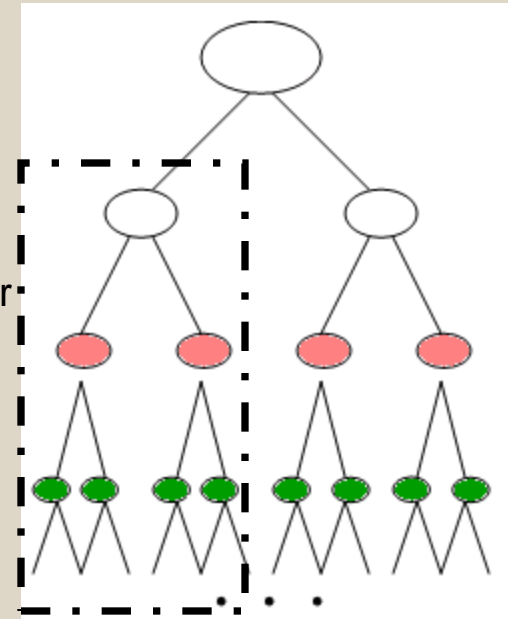
$$L = L_{\min} = \min\{\text{level}(\tau) : \tau \in \mathcal{L}_I\}$$

Frequency is increased

$$\left[1 + \frac{\sqrt{2}}{\kappa \text{diam}(Q_t \times Q_s) + \eta}\right]^{-p} \longrightarrow$$



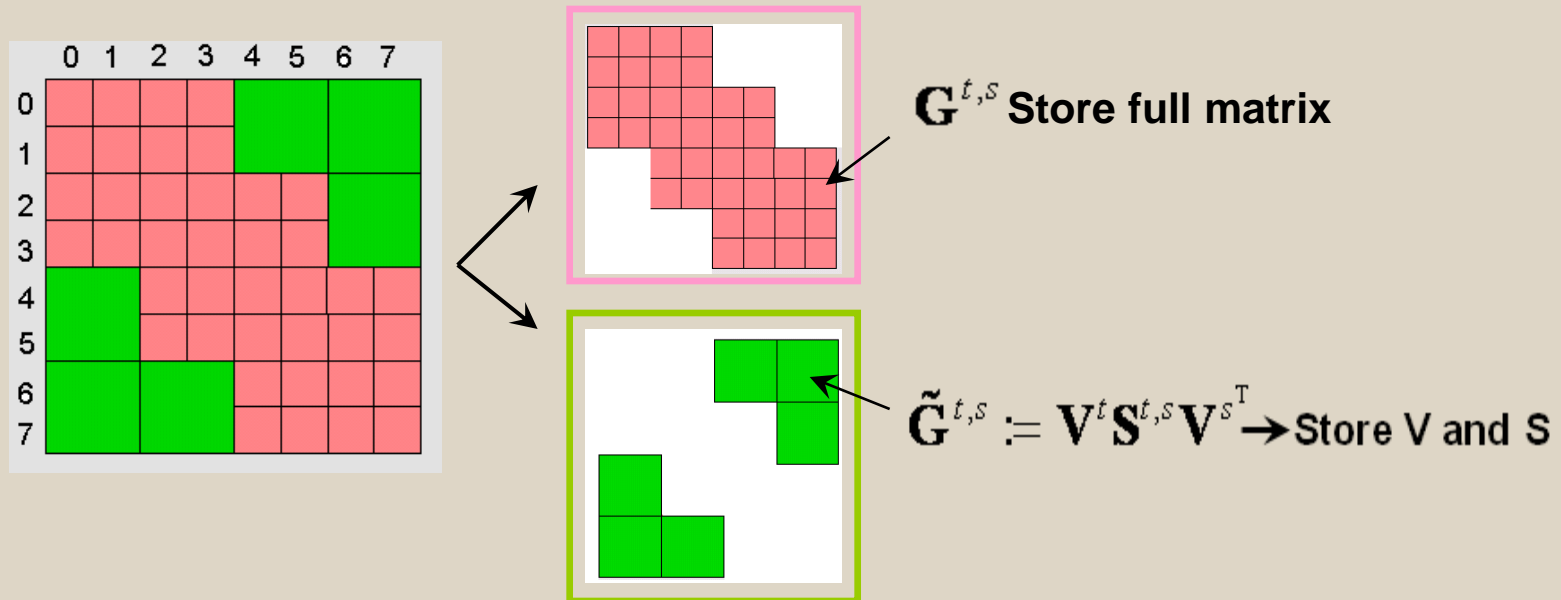
same diameter



☞ An effective means of increasing the rank for accuracy control in a finite electric size range (tens of wavelengths) without sacrificing the linear cost

# $O(N)$ Storage

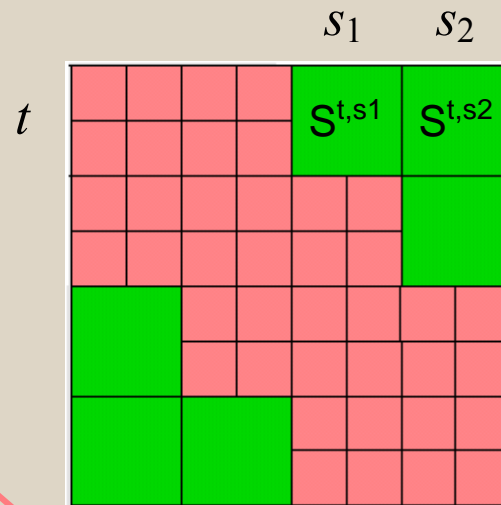
- $\mathcal{H}^2$ -matrix Representation



# $O(N)$ Storage

- Nested property

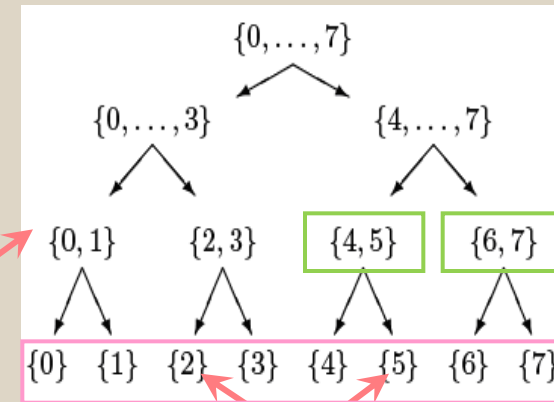
$$\mathbf{V}^t = \begin{pmatrix} \mathbf{V}^{t1} \mathbf{E}^{t1} \\ \mathbf{V}^{t2} \mathbf{E}^{t2} \end{pmatrix} = \begin{pmatrix} \mathbf{V}^{t1} & \\ & \mathbf{V}^{t2} \end{pmatrix} \begin{pmatrix} \mathbf{E}^{t1} \\ \mathbf{E}^{t2} \end{pmatrix}$$



Store  $\mathbf{E}$  for all non-leaf clusters

$$\mathbf{E}_{v',v}^{t'} = L_v^t(\xi_{v'}^{t'})$$

Store  $\mathbf{V}$  for leaf clusters



Store  $\mathbf{V}$  for leaf clusters



# $O(N)$ Storage

- Cost Analysis

$$\begin{aligned}
 St(\mathcal{H}^2\text{-matrix}) &= St(\text{all leaf clusters}) + St(\text{all nonleaf clusters}) + St(\text{all admissible blocks}) + St(\text{all inadmissible blocks}) \\
 &= St(\mathbf{V}^t) + St(\text{transfer matrix } \mathbf{E}^t) + St(\text{coupling matrix } \mathbf{S}^b) + St(\text{full matrix } \mathbf{G}^b) \\
 &= \sum_{t \in \mathcal{L}_{\mathcal{T}}} O(k_{\text{var}}(t)) \# \hat{t} + \sum_{t \in T_{\mathcal{T}} \setminus \mathcal{L}_{\mathcal{T}}} \sum_{t' \in \text{sons}(t)} O(k_{\text{var}}(t) k_{\text{var}}(t')) + \sum_{b=(t,s) \in \mathcal{L}_{\mathcal{T}}^+ \times \mathcal{T}} O(k_{\text{var}}(t) k_{\text{var}}(s)) + \sum_{b=(t,s) \in \mathcal{L}_{\mathcal{T}}^- \times \mathcal{T}} \# \hat{t} \# \hat{s} \\
 &\leq O(\hat{\alpha}^d) \cdot N + \sum_{t \in T_{\mathcal{T}}} \sum_{t' \in \text{sons}(t)} O(k_{\text{var}}(t) k_{\text{var}}(t')) + \sum_{t \in T_{\mathcal{T}}} \sum_{s \in \text{col}(t)} O(k_{\text{var}}^2(t)) + \sum_{t \in T_{\mathcal{T}}} \sum_{s \in \text{col}(t)} n_{\min}^2 \\
 &\leq O(\hat{\alpha}^d) \cdot N + 2 \sum_{t \in T_{\mathcal{T}}} O(k_{\text{var}}^2(t)) + C_{sp} \sum_{t \in T_{\mathcal{T}}} O(k_{\text{var}}^2(t)) + C_{sp} \sum_{t \in T_{\mathcal{T}}} n_{\min}^2 \\
 &\leq O(\hat{\alpha}^d) \cdot N + 2 \sum_{t \in T_{\mathcal{T}}} O(k_{\text{var}}^2(t)) + C_{sp} \sum_{t \in T_{\mathcal{T}}} O(k_{\text{var}}^2(t)) + 2C_{sp} n_{\min}^2 N
 \end{aligned}$$

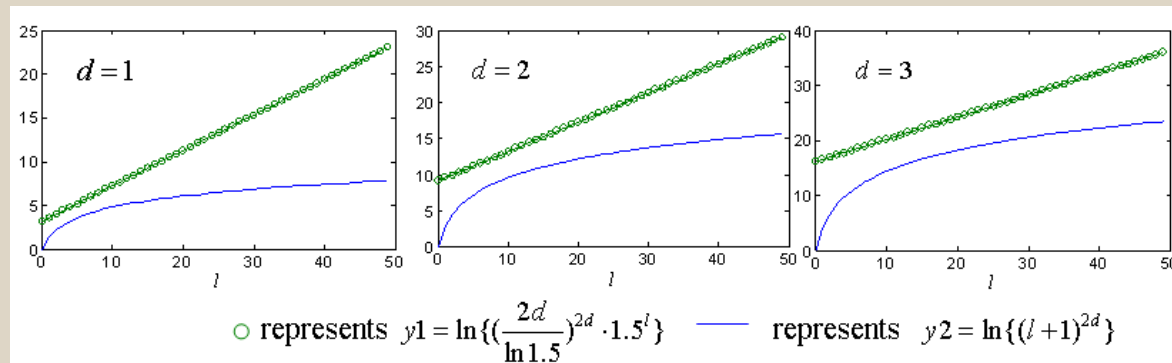




# $O(N)$ Storage

$$\begin{aligned}
 \sum_{t \in T_2} O(k_{\text{var}}^2(t)) &= O\left(\sum_{l=0}^L \sum_{t \in T_2^{(l)}} [\hat{a} + \hat{b}(L-l)]^{2d}\right) \leq O\left(\sum_{l=0}^L [\hat{a} + \hat{b}(L-l)]^{2d} \cdot 2^l\right) \\
 &\leq O((\hat{a} + \hat{b})^{2d}) \sum_{l=0}^L (1+L-l)^{2d} \cdot 2^l \leq O((\hat{a} + \hat{b})^{2d}) \cdot 2^L \cdot \sum_{l=0}^L (l+1)^{2d} \cdot 2^{-l} \leq O((\hat{a} + \hat{b})^{2d}) \cdot 2^L \cdot \sum_{l=0}^L \left(\frac{2d}{\ln 1.5}\right)^{2d} \cdot 1.5^l \cdot 2^{-l} \\
 &\leq 2^{L+1} \cdot O\left(\left(\frac{2d(\hat{a} + \hat{b})}{\ln 1.5}\right)^{2d}\right) \cdot \sum_{l=0}^L \left(\frac{3}{4}\right)^l \leq 2^{L+1} \cdot O\left(\left(\frac{2d(\hat{a} + \hat{b})}{\ln 1.5}\right)^{2d}\right) \cdot 4 \leq 4O\left(\left(\frac{2d(\hat{a} + \hat{b})}{\ln 1.5}\right)^{2d}\right) \cdot N
 \end{aligned}$$

$(l+1)^{2d} \leq \left(\frac{2d}{\ln 1.5}\right)^{2d} \cdot 1.5^l$



$$\text{St}(\mathcal{H}^2\text{-matrix}) \Leftarrow O(\tilde{a}^2) \cdot N + 8O\left(\left(\frac{2d(\hat{a} + \hat{b})}{\ln 1.5}\right)^{2d}\right) \cdot N + 4C_{\text{sp}} O\left(\left(\frac{2d(\hat{a} + \hat{b})}{\ln 1.5}\right)^{2d}\right) \cdot N + 2C_{\text{sp}} O(n_{\text{fin}}^2) \cdot N$$





# $O(N)$ Matrix-Vector Multiplication

- If  $\mathbf{G}^{t,s}$  is an **admissible** block

Matrix-vector multiplication can be performed in **three** steps

$$\tilde{\mathbf{G}}_i \cdot \mathbf{x} := \mathbf{V}^t \mathbf{S}^{t,s} \mathbf{V}^{sT} \cdot \mathbf{x}$$

(a) Forward transformation

$$\mathbf{x}^s = \mathbf{V}^{sT} \cdot \mathbf{x} \longrightarrow O(N)$$

(b) Coupling-matrix Multiplication

$$\mathbf{y}^t = \sum_{s \in R^t} \mathbf{S}^{t,s} \mathbf{x}^s \longrightarrow O(N)$$

(c) Backward transformation

$$\mathbf{y}_i = \sum_{t \in \hat{t}} (\mathbf{V}^t \mathbf{y}^t)_i \longrightarrow O(N)$$

- If  $\mathbf{G}^{t,s}$  is an **inadmissible** block

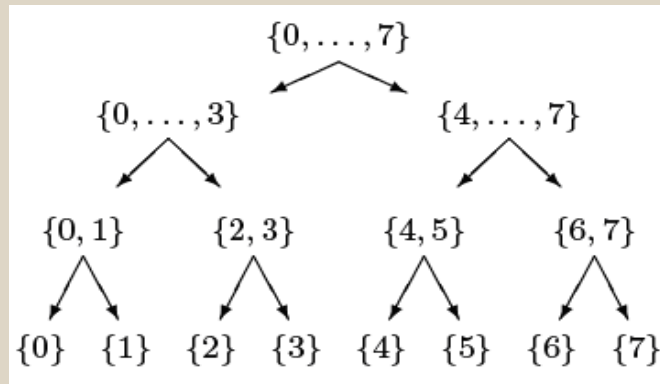
(d) Direct multiplication

$$\longrightarrow O(N)$$



$$\tilde{\mathbf{G}}^{t,s} \cdot x := \mathbf{V}^t \mathbf{S}^{t,s} \mathbf{V}^{s^T} \cdot x$$

(a) Forward transformation  $\mathbf{x}^s = \mathbf{V}^{s^T} \cdot \mathbf{x}$



$$\mathbf{V}^s = \begin{pmatrix} \mathbf{V}^{s1} \mathbf{E}^{s1} \\ \mathbf{V}^{s2} \mathbf{E}^{s2} \end{pmatrix} = \begin{pmatrix} \mathbf{V}^{s1} & \\ & \mathbf{V}^{s2} \end{pmatrix} \begin{pmatrix} \mathbf{E}^{s1} \\ \mathbf{E}^{s2} \end{pmatrix}$$



Non-leaf cluster



$$\mathbf{x}^s = \mathbf{V}^{s^T} \mathbf{x}|_s = \sum_{s' \in \text{sons}(s)} (\mathbf{V}^{s'} \mathbf{E}^{s'})^T \mathbf{x}|_{s'} = \sum_{s' \in \text{sons}(s)} (\mathbf{E}^{s'})^T \mathbf{x}^{s'}$$

Leaf cluster



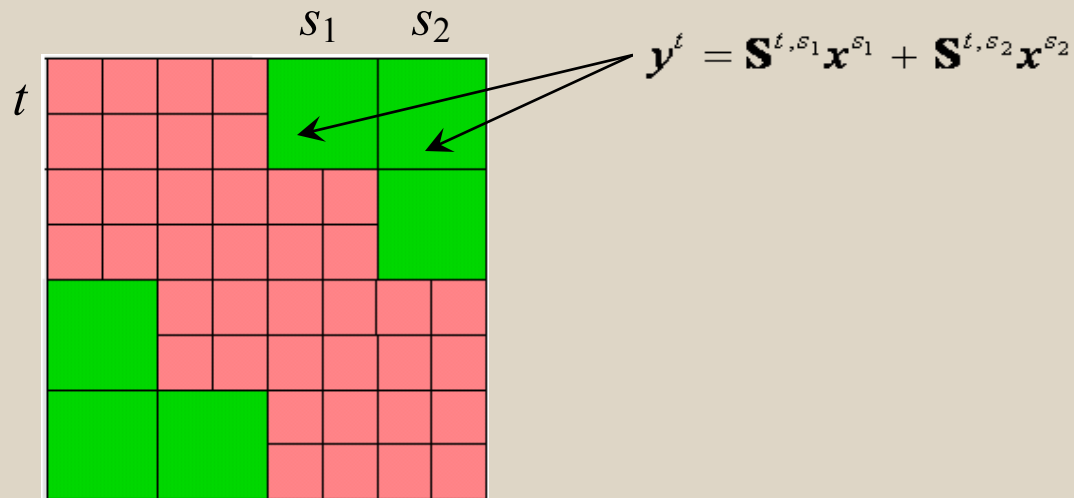
$$\mathbf{x}^s = \mathbf{V}^{s^T} \mathbf{x}|_s$$

$$\begin{aligned} \text{Comp}(\sum_{s \in T_{\mathcal{X}}} \mathbf{V}^{s^T} \mathbf{x}|_s) &= \sum_{s \in \mathcal{L}(T_{\mathcal{X}})} O(k_{\text{var}}(s)) \# \hat{s} + \sum_{\substack{s \in T_{\mathcal{X}} \setminus \mathcal{L}_{\mathcal{X}} \\ s' \in \text{sons}(s)}} O(k_{\text{var}}(s) k_{\text{var}}(s')) \leq O(\hat{a}^d) N + \sum_{s \in T_{\mathcal{X}}} O(k_{\text{var}}^2(s)) \\ &\leq O(\hat{a}^d) N + 2O\left(\left(\frac{2d(\hat{a} + \hat{b})}{\ln 1.5}\right)^{2d}\right) N \end{aligned}$$



$$\tilde{\mathbf{G}}^{t,s} \cdot x := \mathbf{V}^t \mathbf{S}^{t,s} \mathbf{V}^{s^T} \cdot x$$

**(b) Coupling-matrix Multiplication**  $y^t = \sum_{s \in R^t} \mathbf{S}^{t,s} X^s$

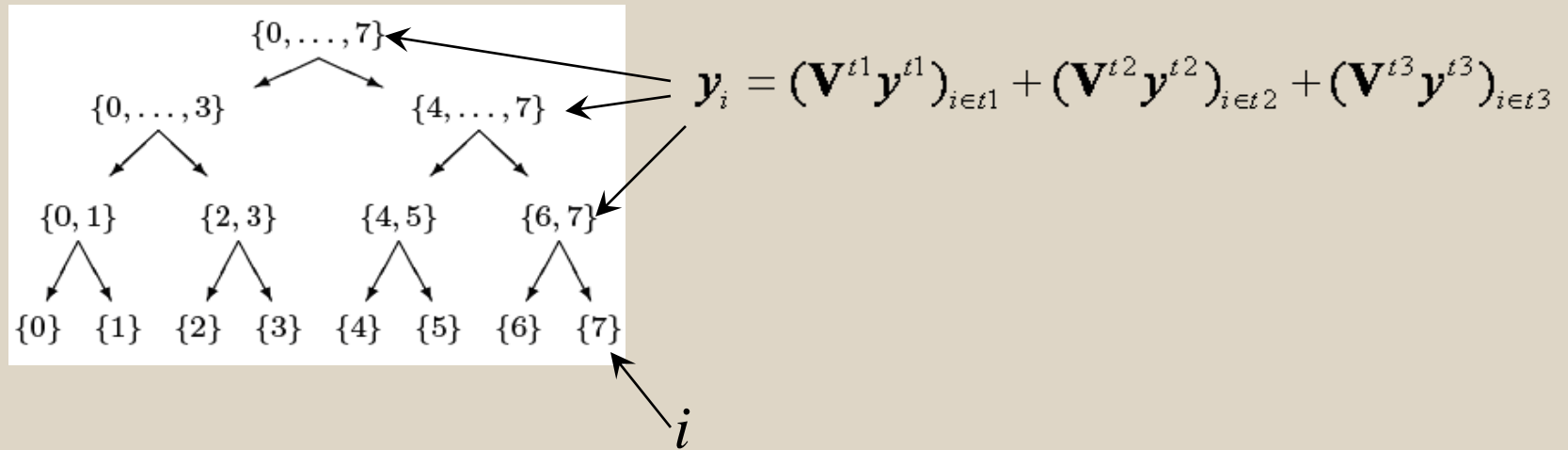


$$\begin{aligned} \text{Comp}(\sum_{b=(t,s) \in \mathcal{L}_{\mathcal{I} \times \mathcal{I}}^+} \mathbf{S}^{t,s} \mathbf{x}^s) &= \sum_{b=(t,s) \in \mathcal{L}_{\mathcal{I} \times \mathcal{I}}^+} O(k_{\text{var}}^2(b)) \leq \sum_{b=(t,s) \in \mathcal{L}_{\mathcal{I} \times \mathcal{I}}} O(k_{\text{var}}^2(b)) \leq \sum_{t \in \mathcal{I}_{\mathcal{I}}} \sum_{s \in \text{col}(t)} O(k_{\text{var}}^2(b)) \\ &\leq C_{\mathcal{I}} \sum_{t \in \mathcal{I}_{\mathcal{I}}} O(k_{\text{var}}^2(b)) \leq 2C_{\mathcal{I}} O\left(\frac{2d(\hat{a} + \hat{b})}{\ln 1.5}\right)^{2d} \cdot N \end{aligned}$$



$$\tilde{\mathbf{G}}^{t,s} \cdot x := \mathbf{V}^t \mathbf{S}^{t,s} \mathbf{V}^{s^T} \cdot x$$

**(c) Backward transformation**  $y_i := \sum_{t, i \in \hat{t}} (\mathbf{V}^t \mathbf{y}^t)_i$



$$\text{Comp}(\sum_{t \in T_x} \mathbf{V}^t \mathbf{y}^t) = \sum_{t \in \mathcal{L}(T_x)} O(k_{\text{var}}(t)) \# \hat{t} + \sum_{\substack{t \in T_x \setminus \mathcal{L}_x \\ t' \in \text{sons}(t)}} O(k_{\text{var}}(t) k_{\text{var}}(t')) \leq O(\hat{a}^d) N + 2O((\frac{2d(\hat{a} + \hat{b})}{\ln 1.5})^{2d}) N$$



# $O(N)$ Matrix-Vector Multiplication

If  $\mathbf{G}^{t,s}$  is **inadmissible** block

Since they are full matrices, multiplication can be done directly.

$$Comp(\sum_{b=(t,s) \in \mathcal{L}_{\mathcal{I} \times \mathcal{I}}} \mathbf{G}^{t,s} x|_{\hat{s}}) = \sum_{b=(t,s) \in \mathcal{L}_{\mathcal{I} \times \mathcal{I}}} O(n_{\min}^2) \leq 2C_{sp} O(n_{\min}^2) N$$



# $O(N)$ Matrix-Vector Multiplication

Add four steps

$$\text{Comp}(\mathcal{H}^2 - \text{Matrix-vector mult.}) \rightarrow O(N)$$



# Numerical Results

- **Dielectric Sphere**
- **Triangular Dielectric Rod**





# Case 1: Dielectric Sphere ( $\epsilon_r = 36$ )

Parameters:  $n_{\min} = 18$ ,  $\eta = 0.95$ ,  $\hat{a} = 4$ ,  $\hat{b} = 2$

Electric size:  $k_0 a = 0.408$

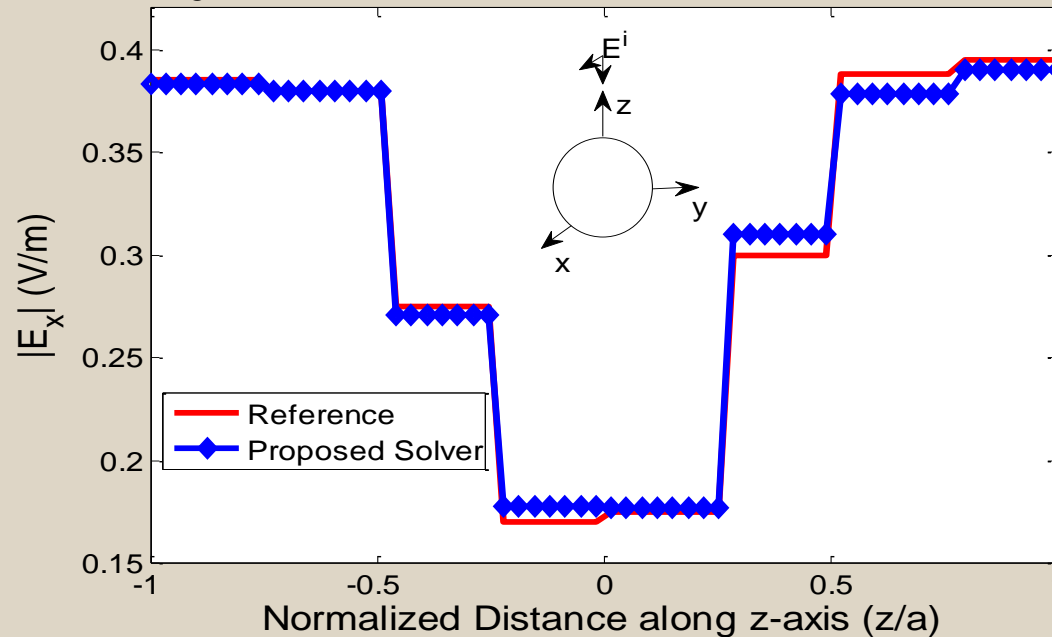
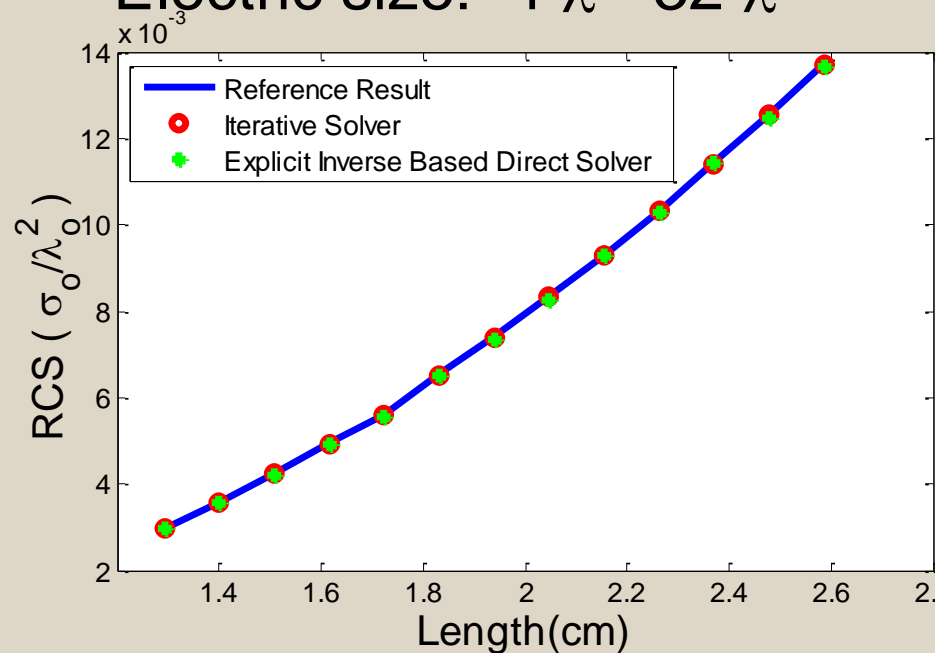


Fig. 1. Field along z-axis of a dielectric sphere with  $\epsilon_r = 36$  (Reference data is from [1]).

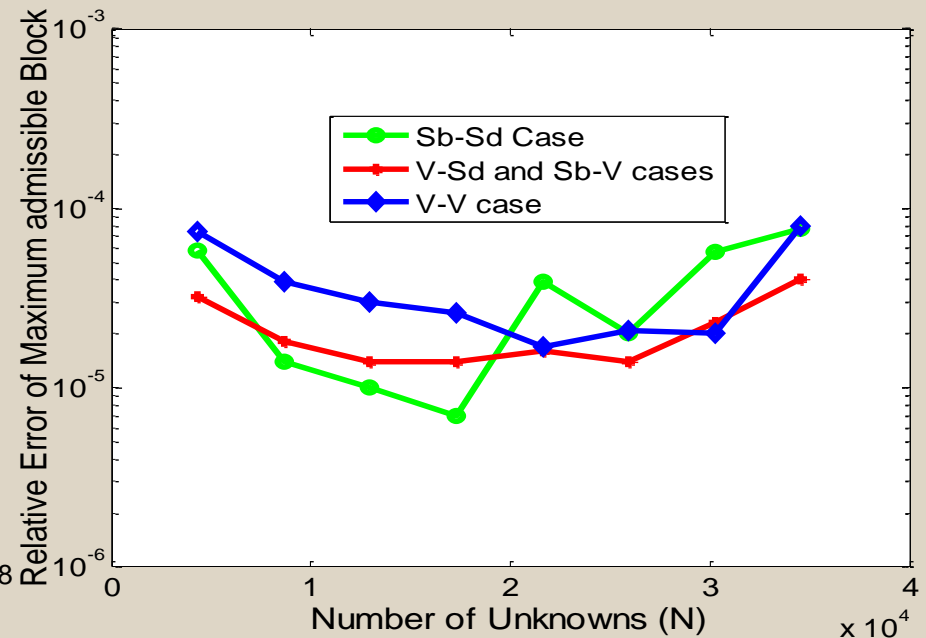
# Case 2: Triangular Dielectric Rod ( $\epsilon_r = 2.54$ )

Parameters:  $n_{\min} = 80$ ,  $\eta = 0.95$ ,  $\hat{a} = 4(5)$ ,  $\hat{b} = 4(5)$

Electric size:  $4\lambda - 32\lambda$



(a) RCS comparison

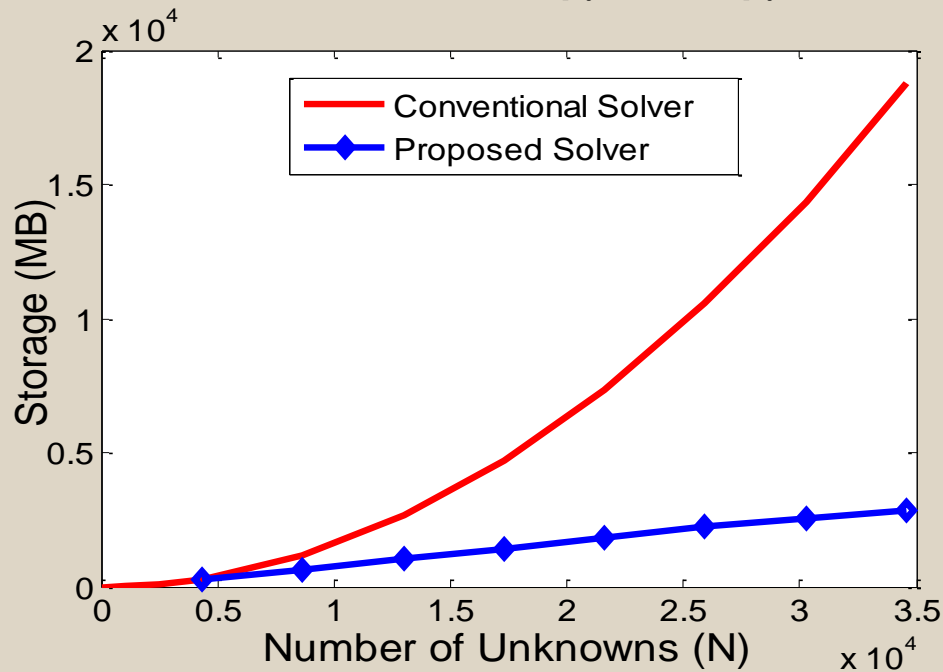


(b) Relative error of the maximal admissible block

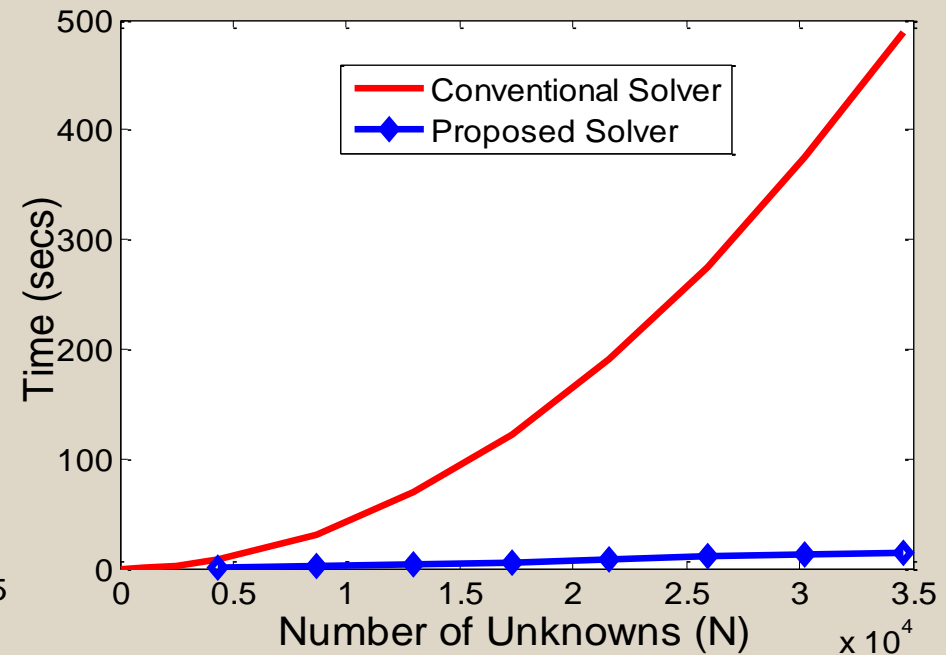
# Case 2: Triangular Dielectric Rod ( $\epsilon_r = 2.54$ )

Parameters:  $n_{\min} = 80$ ,  $\eta = 0.95$ ,  $\hat{a} = 4(5)$ ,  $\hat{b} = 4(5)$

Electric size:  $4\lambda - 32\lambda$



(c) Storage comparison



(d) CPU MVM time comparison

Fig. 2. Simulation of a dielectric rod with  $\epsilon_r = 2.54$ .

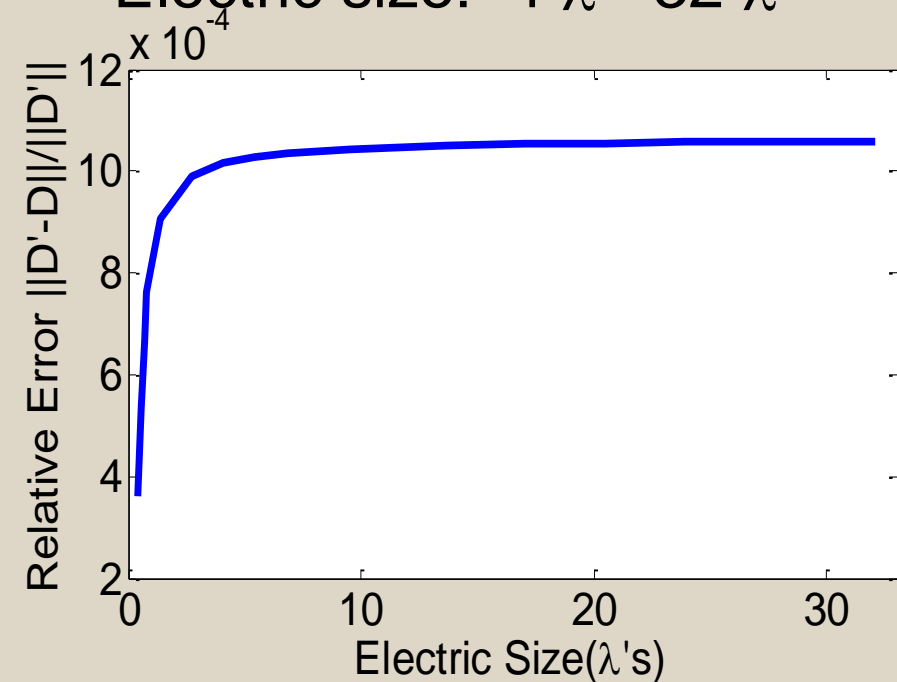


# Case 2: Triangular Dielectric Rod ( $\epsilon_r = 2.54$ )

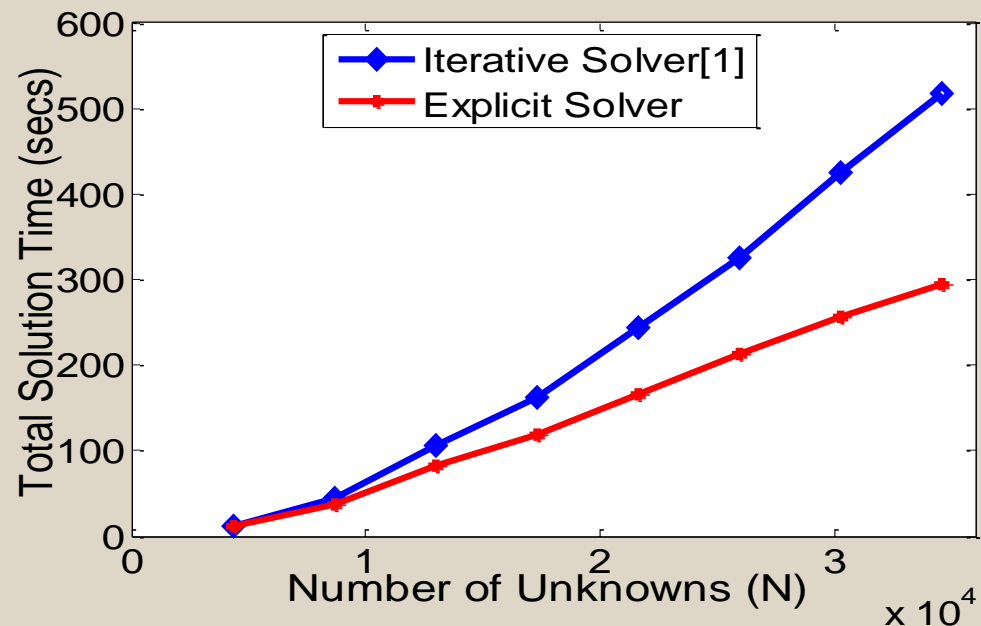
Parameters:  $n_{\min} = 80$ ,  $\eta = 0.95$ ,  $\hat{a} = 4(5)$ ,  $\hat{b} = 4(5)$

Electric size:  $4\lambda - 32\lambda$

No. of terms: 20



(e) Relative Error



(f) Solution time comparison with iterative solver[1]

**Fig. 2. Simulation of a dielectric rod with  $\epsilon_r = 2.54$ .**

## Case 2: Triangular Dielectric Rod ( $\epsilon_r = 2.54$ )

Parameters:  $n_{\min} = 80$ ,  $\eta = 0.95$ ,  $\hat{a} = 7(8)$ ,  $\hat{b} = 7(8)$

Electric size:  $80 \lambda$

No. of terms: 20

Accuracy :  $10^{-4}$

Memory : 9 GB (112 GB for conventional solver)

Solution Time : 1675.5 s (61096.3 s for conventional solver)



# Conclusions

## Explicit Inverse Based Fast VIE Solver Developed

- Avoid numerical inversion and troubles of iterations
- Accelerated by  $\mathcal{H}^2$ -based fast matrix-vector multiplications
- Applicable to problems with large dielectric contrast ratio
- Numerical results demonstrate its accuracy and efficiency
- Future work: To explore its application to general problems and larger electric sizes.

