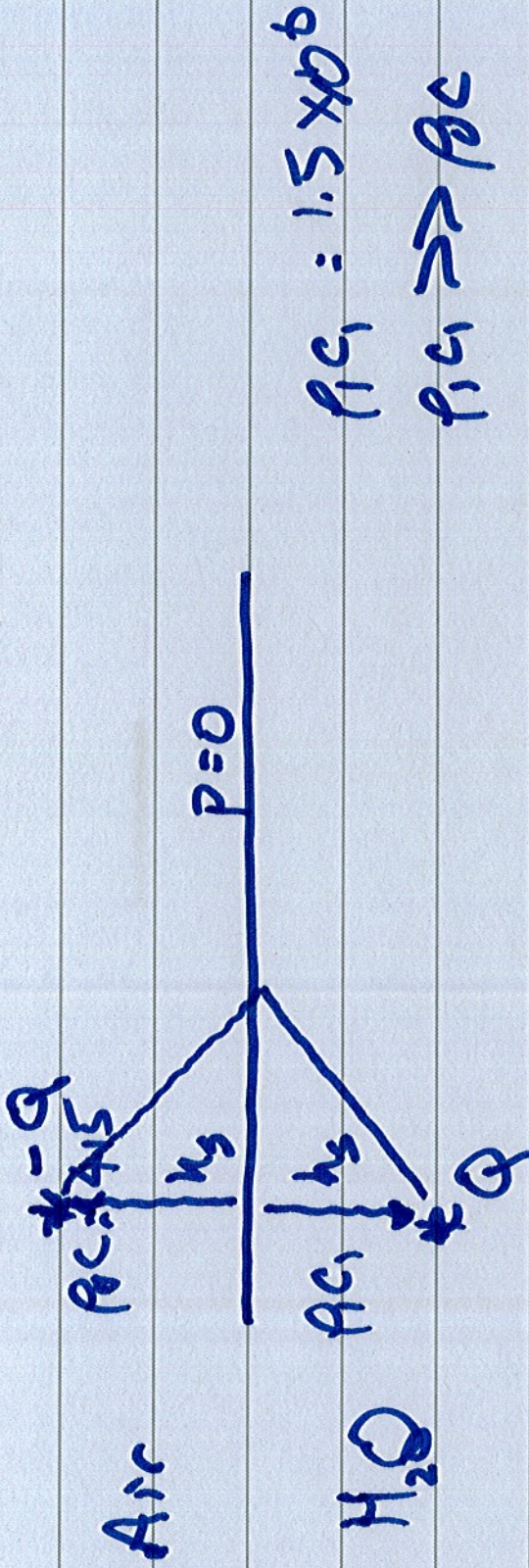


$$\psi(r) \approx A \left\{ \frac{e^{-ikr_1}}{r_1} + R(\theta) \frac{e^{-kr_2}}{r_2} \right\}$$

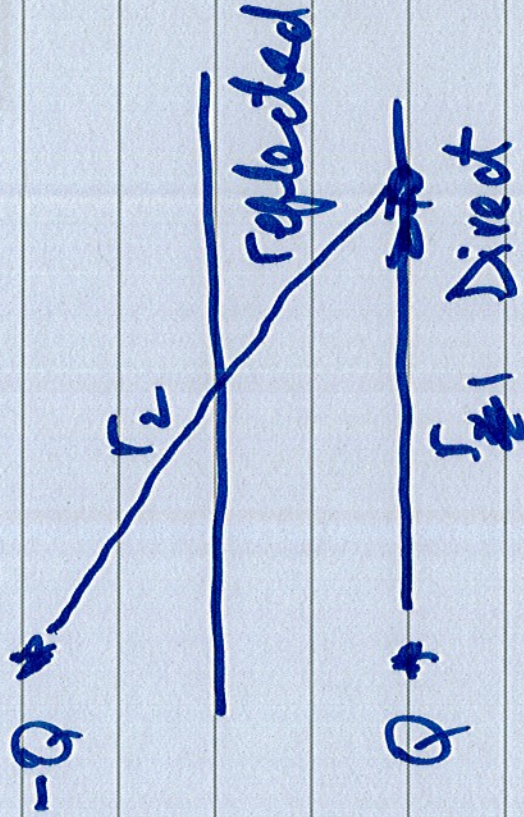
$$\theta < 75^\circ$$

s.3.4.2 Pressure Release Surface

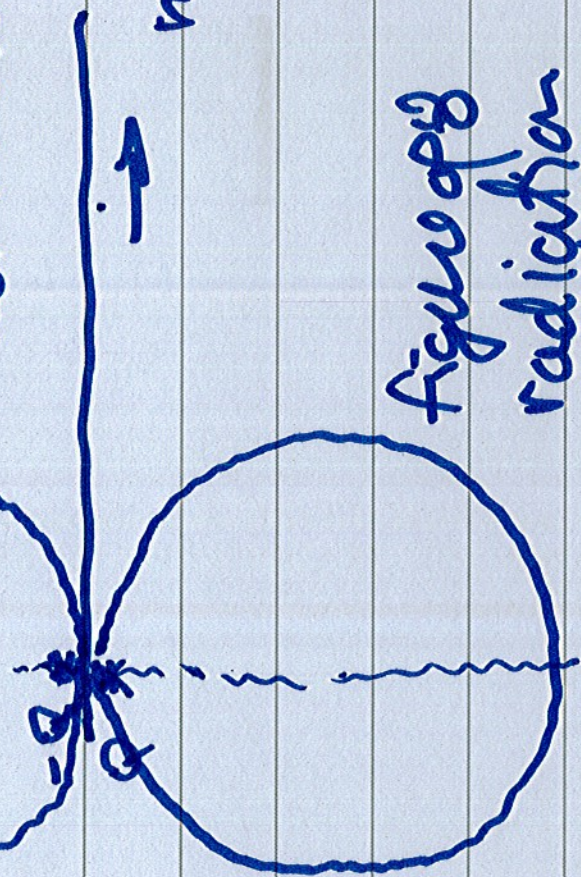


negative image required
to satisfy $P=0$ R.C.

$$\tilde{\psi}(r) = A \left\{ \frac{e^{-ik_1 r_1}}{r_1} - \frac{e^{-k_2 r_2}}{r_2} \right\}$$



let the source approach
the surface



minimum radiation
due to nearly
perfect cancellation

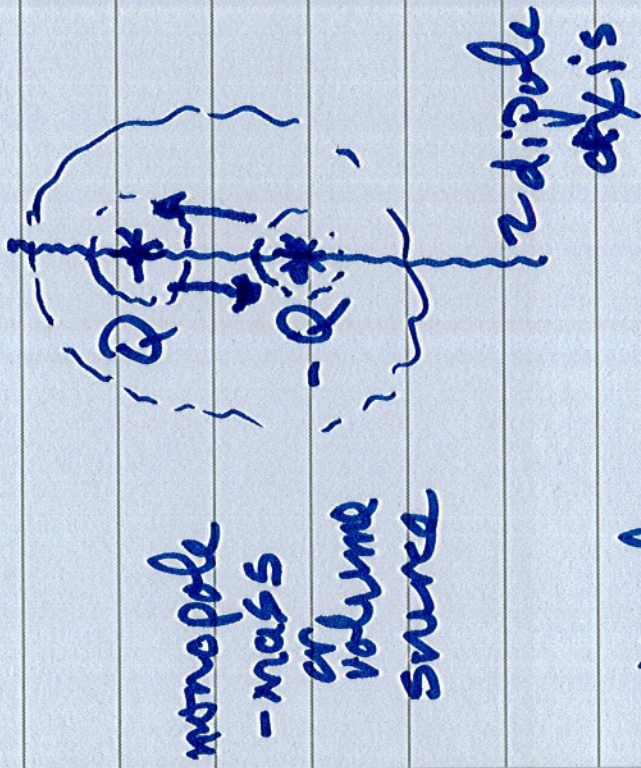
figure 8
radiation
pattern

H₂O

maximum radiation
⊥ to the surface

Dipole - two closely-spaced monopoles of equal strength running out of phase

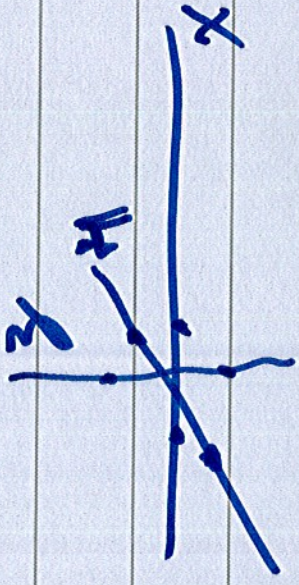
- two monopoles creating an oscillatory flow
- pt force acting on the fluid to accelerate it back & forth



- dipole
 - pt force
 - no change in volume

small axial fan

The diagram shows a dipole source (two small circles) with a dashed line representing a 'small axial fan' extending from the source. Below the fan, there is a horizontal line with a double-headed arrow, indicating the direction of the fan's spread.

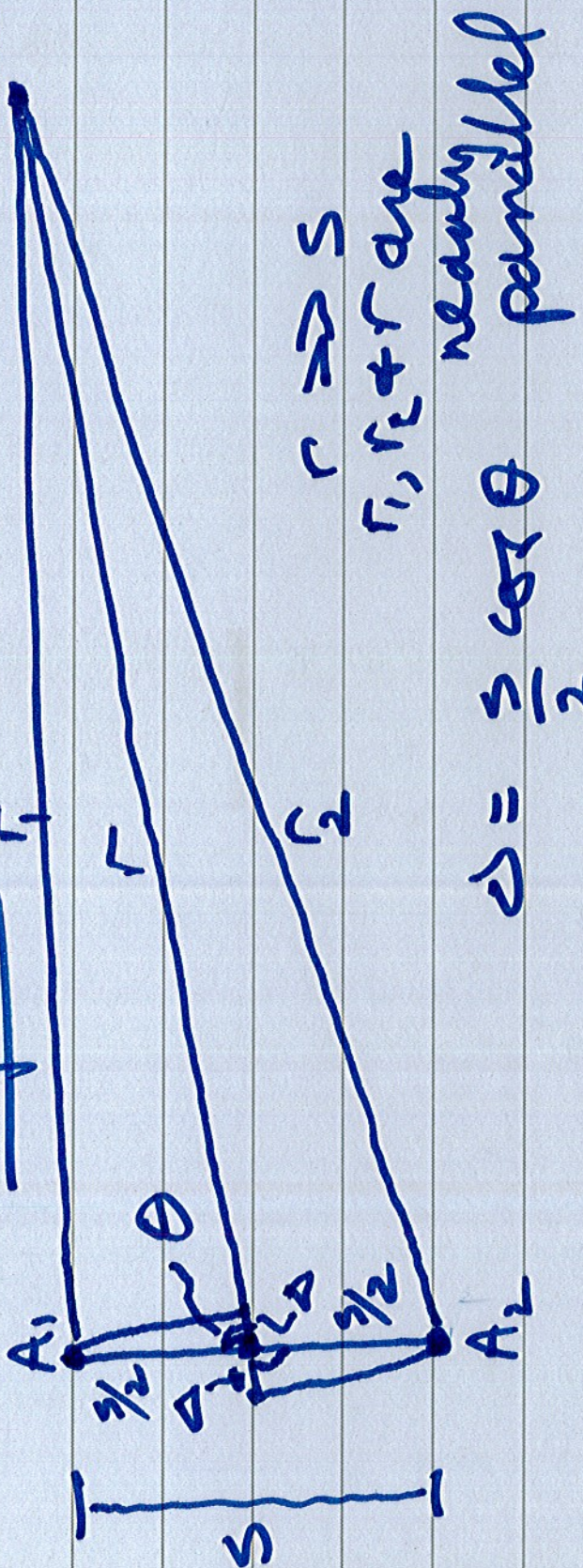


3 orthogonal dipoles can be used to represent a point force in an arbitrary direction

Dipole

5.3.5

receiver



$r \gg s$
 $r_1, r_2 \approx r$ are
 nearly parallel

$$\Delta = \frac{s}{2} \cos \theta$$

θ is the angle from the dipole axis directed towards the positive monopole

$$r_1 \approx r - \Delta = r - \frac{s}{2} \cos \theta$$

$$r_2 \approx r + \Delta = r + \frac{s}{2} \cos \theta$$

8

$$\psi(r) = \cancel{A_1 e^{ikr}} \quad A_1 \frac{e^{-ikr}}{r_1} + \frac{A_2 e^{-ikr}}{r_2}$$

assumed $r \gg r_2$

Next $A_2 = -A_1$

$$A_1 = A$$

$$\vec{p}(r) = A \left\{ \frac{e^{-jk(r-\Delta)}}{r-\Delta} - \frac{e^{-jk(r+\Delta)}}{r+\Delta} \right\}$$

$$= A e^{-jkr} \left\{ \frac{r(e^{+jks} - e^{-jks}) + \Delta(e^{+jks} + e^{-jks})}{r^2 - \Delta^2} \right\}$$

receiver is in the
far field

$$r \gg \Delta$$

$$\Delta \ll r \quad \vec{p}(r) = A \frac{e^{-jkr}}{r} (e^{+jks} - e^{-jks})$$

$$ks = \frac{k\Delta}{2} \cos \theta = 2jA \frac{e^{-jkr}}{r} \sin ks$$

compact source $ks \ll 1$
 $\frac{\Delta}{\lambda} \ll 1$

$$\sin kx = \sin\left(\frac{kx}{2}\right) \underbrace{\cos\left(\frac{kx}{2}\right)}_{\text{very small}}$$

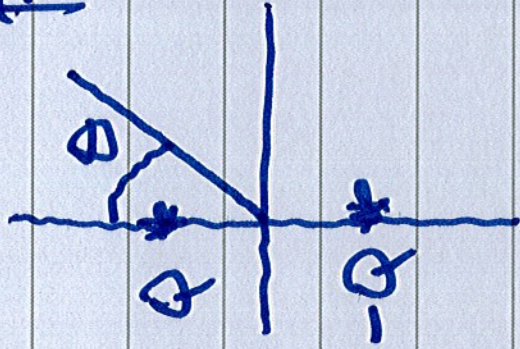
$$\sin \theta \approx \theta$$

when θ

$$\approx \frac{kx}{2}$$

is small

$$\tilde{p}(r) = zj A e^{-jkr} \frac{k^2}{r} \cos \theta$$



$$A = j \rho c \frac{kQ}{4\pi} \quad] \text{ appropriate for the monopoles}$$

Q is the volume source strength $\frac{m^3/s}{m^3/s}$

$$\tilde{p}(r) = -\rho c k^2 (Qs) \frac{e^{-jkr}}{4\pi r} \cos \theta$$

sound field radiated by a point dipole

$Qs =$ dipole source strength

$= D$ dipole moment

Assumptions

- compact source $ks \ll 1$
- receiver is in the farfield

$$s \ll r \quad \Delta \ll r$$

For a monopole

$$\tilde{p}(r) = j\omega c k Q \frac{e^{-jkr}}{4\pi r}$$

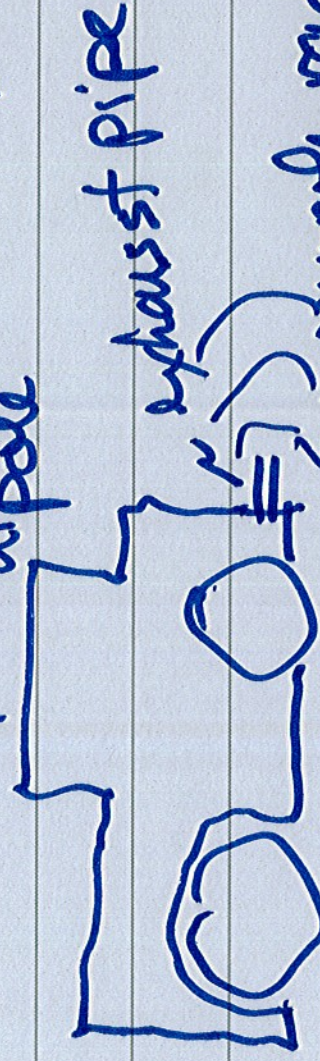
For a dipole

$$\tilde{p}(r) = -\rho c k^2 (Qs) \frac{e^{-jkr}}{4\pi r} \cos\theta$$

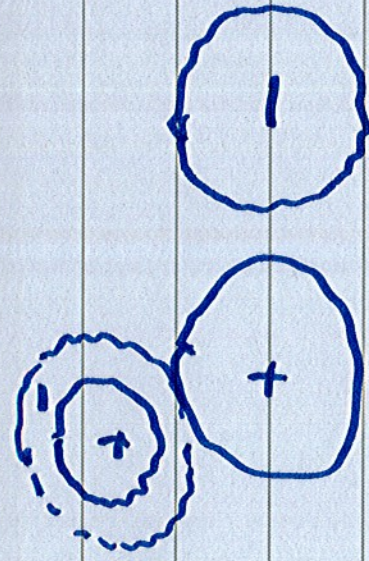
$$\frac{|\tilde{p}(r)|_{\text{dipole}}}{|\tilde{p}(r)|_{\text{monopole}}} = ks \cos\theta \quad ks \ll 1$$

for a given monopole source strength Q ,
dipole is a much weaker source

$$|\hat{P}(r)|_{\text{dipole}} \ll |\hat{P}(r)|_{\text{monopole}}$$



monopole radiation from
the opening of the
pipe



turn monopole
into a dipole

↑
L/S
- greatly reduce the
sound radiation by
reducing radiation efficiency