

3.2.4 Speed of Sound

$$c = \sqrt{\frac{\gamma P_0}{\rho_0}}$$

Isothermal Atmosphere

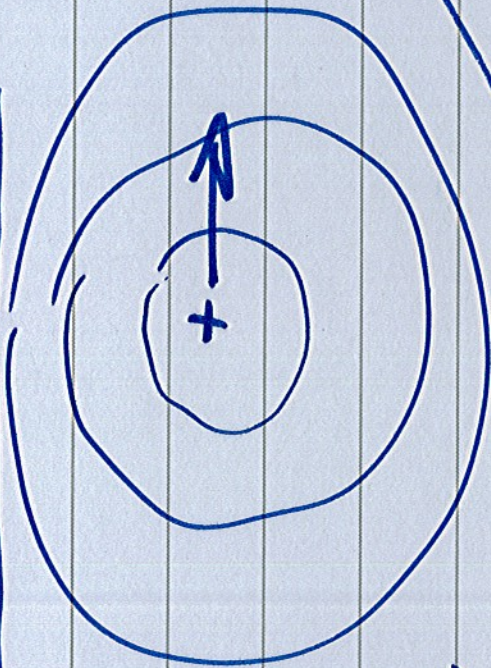
$$\left(\frac{P_0}{\rho_0}\right)$$

$c \neq$ function of height

3.3 One-Dimensional Solutions

- spherical

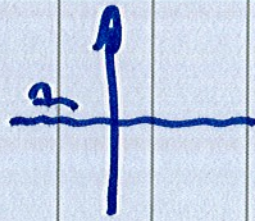
- spherically symmetric



- cylindrical waves

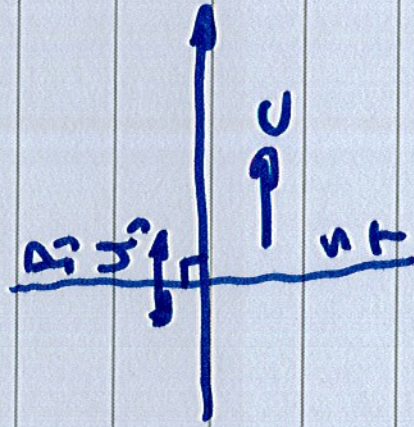


- plane waves



3.3.1 Plane waves

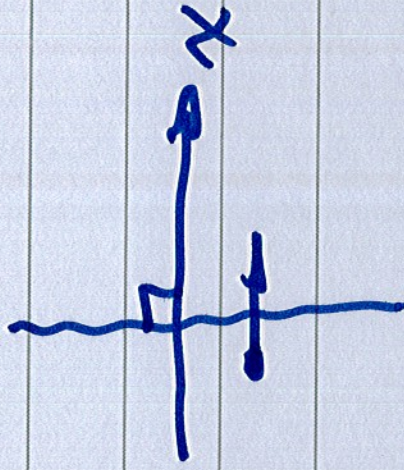
- properties are instantaneously uniform over an infinite plane \perp to the direction of wave propagation



longitudinal wave

- Particle motion is back and forth in the direction of wave propagation

3.3.1.1 Propagation in the x-direction 4



no variation of pressure
in the y or z directions

$$\frac{\partial}{\partial y} = 0$$

$$\nabla^2 p = \frac{\partial^2 p}{\partial x^2} = 0$$

$$\frac{\partial^2 p}{\partial x^2} - \frac{1}{c^2} \frac{\partial^2 p}{\partial t^2} = 0$$

$$p(x,t) = p_1(ct-x) + p_2(ct+x)$$

+ve -ve

complex harmonic wave

$$p(x,t) = \tilde{p}(x) e^{i\omega t}$$

separable form - sub into the

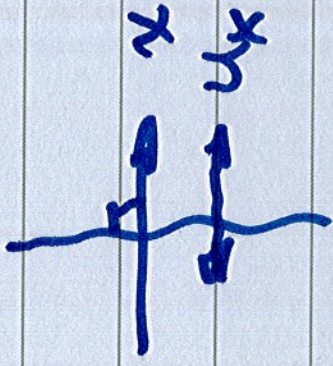
$$\frac{d^2 \tilde{p}}{dx^2} + k^2 \tilde{p} = 0$$

scalar Helmholtz Eqn

$$k = \frac{\omega}{c}$$

$$\tilde{p}(x) = A_1 e^{-ikx} + A_2 e^{+ikx}$$

$$p(x,t) = \tilde{p}(x) e^{i\omega t}$$



linearized momentum eqn

$$-\nabla p = \rho \frac{\partial \bar{u}}{\partial t} \quad \bar{u} = u_x \bar{i} + u_y \bar{j} + u_z \bar{k}$$

$$(1-1) \quad -\frac{\partial p}{\partial x} = \rho \frac{du_x}{dt}$$

$$\begin{aligned} p(x,t) &= \tilde{p}(x) e^{j\omega t} \\ u_x(x,t) &= \tilde{u}_x(x) e^{j\omega t} \end{aligned}$$

$$-\frac{\partial \tilde{p}}{\partial x} = \rho j\omega \tilde{u}_x$$

$$\tilde{u}_x = -\frac{1}{j\omega \rho} \frac{\partial \tilde{p}}{\partial x}$$

7

Given a pressure distribution
- can derive the particle velocity

$$\tilde{u}_x = -\frac{1}{j\omega\rho} \frac{dP}{dx}$$

$$\tilde{u}_y = -\frac{1}{j\omega\rho} \frac{dP}{dy}$$

$$\tilde{u}_z = -\frac{1}{j\omega\rho} \frac{dP}{dz}$$

Particle velocity associated with

$$\vec{P}_T = A_1 e^{-ikx} \quad \frac{P}{x}$$

$$\vec{u}_{xT} = -\frac{1}{j\omega\rho} \frac{\partial \vec{P}_T}{\partial x} = -\frac{1}{j\omega\rho} (-jk) A_1 e^{-ikx}$$

$$= \frac{k}{\omega\rho} A_1 e^{-ikx} \quad k = \frac{\omega}{c}$$

$$= \frac{A_1 e^{-ikx}}{\rho c} \text{ characteristic impedance}$$

$$\vec{u}_{xT} = \frac{P_T}{\rho c}$$

$$(\rho c) \approx 415 \text{ [Rayls]}$$

specific acoustic impedance

$\frac{\tilde{P}_+}{\tilde{u}_{x+}} = \rho c$ characteristic impedance & impedance

$$\tilde{P}_- = A_2 e^{+jkx} \quad \leftarrow x \quad \leftarrow \tilde{u}_{x-}$$

$$\tilde{u}_{x-} = -\frac{1}{j\omega\rho} \frac{d\tilde{P}_-}{dx} = -\frac{1}{j\omega\rho} (jk) A_2 e^{+jkx}$$

$$= -\frac{(A_2 e^{+jkx})}{\rho c} \tilde{P}_-$$

$$\tilde{u}_{x-} = -\frac{\tilde{P}_-}{\rho c} = -\rho c$$

ρc sign of the impedance depends on direction of prop

we use these results to calculate the particle velocity for any given sound field

$$\bar{u} = -\frac{1}{j\omega\rho_0} \nabla \hat{p}$$

harmonic motion

$$\hat{p}(x) = A_1 e^{-ikx} + A_2 e^{+ikx}$$

$$\begin{aligned} \hat{u}_x(x) &= \frac{\hat{p}_+}{\rho_0 c} - \frac{\hat{p}_-}{\rho_0 c} \\ &= \frac{A_1}{\rho_0 c} e^{-ikx} - \frac{A_2}{\rho_0 c} e^{+ikx} \end{aligned}$$

Typical Velocities

- plane wave prop in the +ve x-dir

$$|\vec{p}_+| = 1 \text{ Pa} \quad \approx 94 \text{ dB}$$

$$|\vec{u}_{x+}| = \frac{|\vec{p}_+|}{\rho_0 c} = 2.4 \text{ mm/s}$$

$$c = 340 \text{ m/s}$$

$$|\vec{u}_{x+}| \ll c$$

x_i

12

ξ displacement

for harmonic fields

$$j\omega \tilde{\xi}_{x_t} = \tilde{u}_{x_t}$$

$$|\tilde{\xi}_x| = \frac{|\tilde{u}_{x_t}|}{\omega}$$

ξ particle displacement

$$|\tilde{\xi}_{x+}| = \frac{|\tilde{u}_{x+}|}{\omega}$$

at 1 kHz $\omega \approx 6000$

$$94 \text{ dB} \quad |\tilde{\xi}_{x+}| = \frac{2.4 \times 10^{-3}}{6 \times 10^{+3}} \approx 0.25 \times 10^{-6} \quad \underline{\text{micron}}$$

$$74 \text{ dB} \quad |\tilde{\xi}_{x+}| \quad \underline{0(10^{-7})}$$

$$54 \text{ dB} \quad |\tilde{\xi}_{x+}| \quad \underline{0(10^{-8})}$$

$$34 \text{ dB} \quad |\tilde{\xi}_{x+}| \quad \underline{0(10^{-9})}$$

nanometers