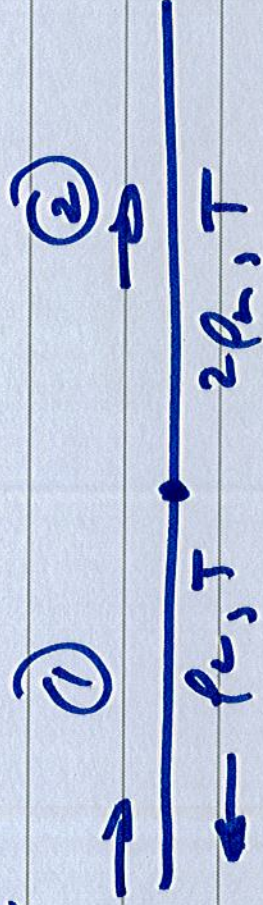


2.8.2



$$y_1 = \underline{A e^{j(\omega t - k_1 x)}} + B e^{j(\omega t + k_1 x)}$$

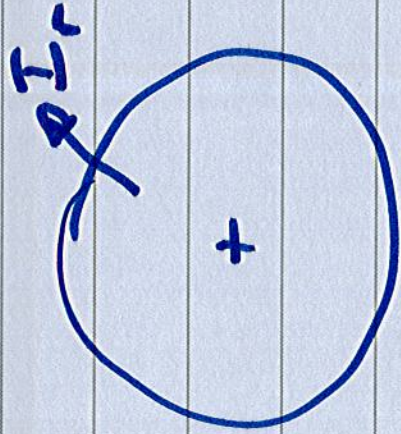
$$y_2 = C e^{j(\omega t - k_2 x)}$$

Acoustic Intensity

\bar{I} = sound power / unit area

$$= \frac{1}{2} \operatorname{Re} \{ \bar{p} \bar{u}^* \}$$

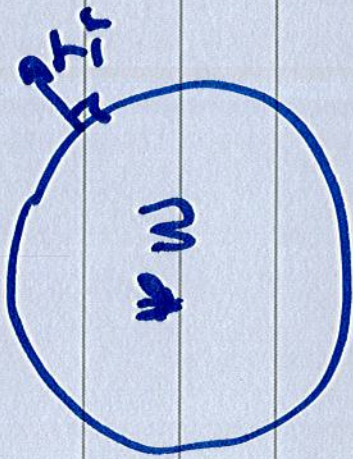
$$= \frac{\bar{P}_{rms}}{\rho_0 c} \quad \text{freely propagating plane wave}$$



$$I_r = \frac{|\bar{p}_+|^2}{2 \rho_0 c} \quad |\bar{p}_+| \propto \frac{1}{r}$$

$$I_r \propto \frac{1}{r^2}$$

no intensity nearfield



$$W = \int_S I_r \, dS$$

spher sym $\underline{W} = I_r (4\pi r^2)$

$$I_r = \frac{W}{4\pi r^2}$$

3.6 Decibels

- sound pressure cover an enormous range

Very loud sound 20 Pa $[120 \text{ dB}]$

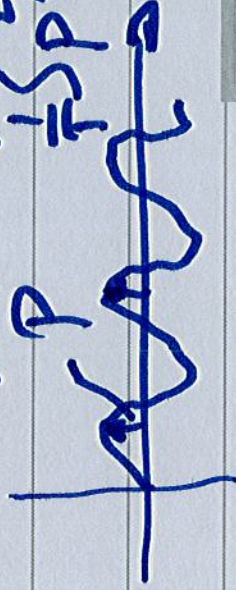
Very quiet $30 \text{ dB} \Rightarrow 6 \times 10^{-4} \text{ Pa}$

humans respond on a logarithmic scale.

Level - always refers to decibels

Sound Pressure Level

$$L_p = 10 \log \left(\frac{P_{rms}^2}{P_{ref}^2} \right) \sim \text{mean square pressure}$$



$$= 20 \log \frac{P_{rms}}{P_{ref}}$$

$$P_{rms} = 0.2 \text{ Pa}$$

$$P_{rms} = \frac{|P|^2}{2}$$

$$20 \mu\text{Pa}$$

$$= 2 \times 10^{-5} \text{ Pa}$$

minimum SPLs audible by healthy young adult $\approx 1 \text{ kHz}$

for sinusoidal signals

$$\beta_p = 87 \text{ dB re } 20 \mu\text{Pa}$$

Sound Intensity level

$$\beta_I = 10 \log \frac{I}{I_{\text{ref}}} \text{ dB re } I_{\text{ref}}$$

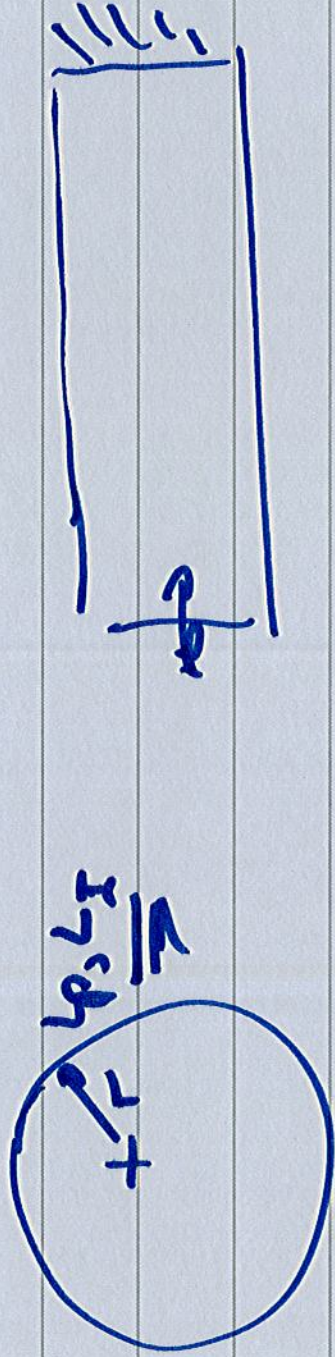
$$I_{\text{ref}} = 1 \times 10^{-12} \text{ W/m}^2$$

for a freely propagating
plane wave

$$I = \frac{P_{\text{rms}}^2}{\rho_0 c} = \frac{P_{\text{ref}}^2}{\rho_0 c}$$

with this choice
in a free field

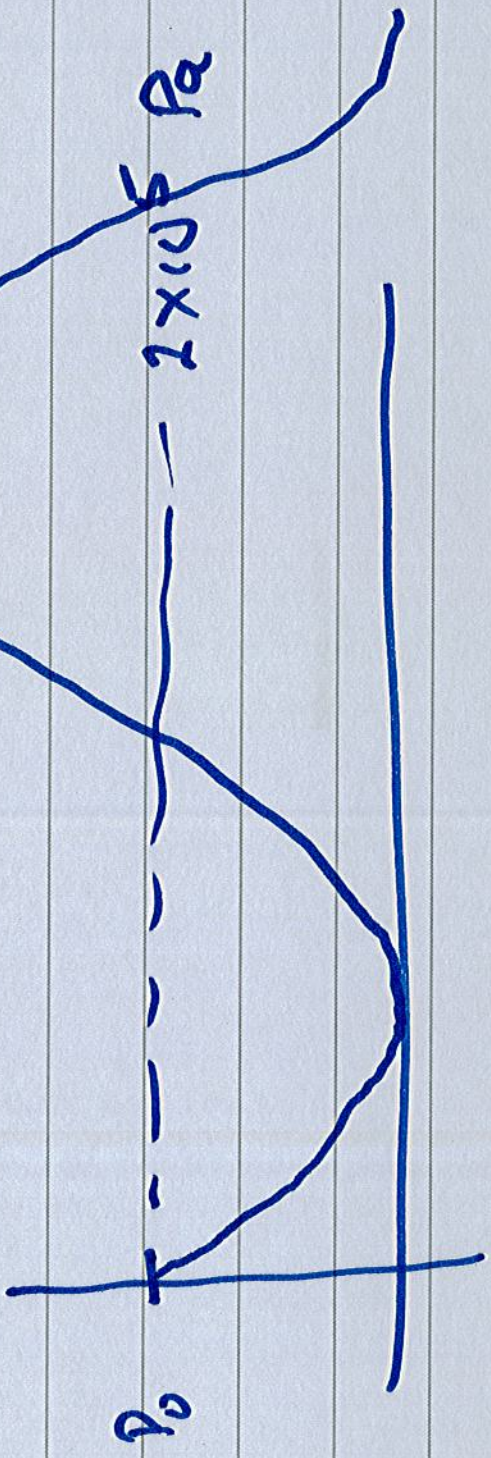
L_p & L_I are numerically the same



Sound Power Level

$$L_w = 10 \log \frac{W}{W_{ref}} \quad \text{dB re } W_{ref}$$

$$W = 1 \times 10^{-12} \text{ Watts}$$



$$L_p \approx 10 \log \frac{1 \times 10^{-10}}{4 \times 10^{-10}} \approx 194 \text{ dB}$$

re 20 μ Pa

$$0 \leq L_p \leq 194$$

above above 120 dB

non-linear effects become significant

$$70 \text{ dB} + 70 \text{ dB} \neq 140$$

└──────────┘

73 dB

Adding decibels \rightarrow add the corresponding
mean square pressures
(quadratic quantities)

$$L_p = 10 \log \frac{\overline{p_{rms}^2}}{p_{ref}^2}$$

$$L_{p_1} + L_{p_2} + \dots = L_{total} \rightarrow \text{since signals are statistically uncorrelated}$$

$$(\overline{p_{rms}^2})_1 = p_{ref}^2 10^{L_{p_1}/10}$$

$$(\overline{p_{rms}^2})_2 = p_{ref}^2 10^{L_{p_2}/10}$$

⋮

⋮

$$(P_{rms})_N = P_{ref}^2 \cdot 10^{L_N/10}$$

$$(P_{rms})_{total} = (P_{rms1}^2) + (P_{rms2}^2) + \dots$$

$$L_{total} = 10 \log_{10} \frac{(P_{rms})_{total}}{P_{ref}^2}$$

$$2 \times |P| \rightarrow 4 P_{rms}^2 \quad \left. \begin{array}{l} \text{pressure} \\ \text{doubling} \end{array} \right\} \underline{\underline{6 \text{ dB}}}$$

||

$L_{p1} \sim 70 \text{ dB}$

$L_{p2} \sim 55 \text{ dB}$

70

$\frac{73}{72}$
75

Section 3

Development of wave Equations

- assumptions

eqns of state
" " continuity
" " motion

2nd order
PDE

- linearized momentum equation

Pressure \rightarrow Particle velocity

One-Dimensional Solutions

- Plane $|p| \neq f(x)$

- cylindrical $|p| \propto r^{1/2}$

- spherical $|p| \propto \frac{1}{r}$

Specific Acoustic Impedance

$$\tilde{z} = \frac{\tilde{p}}{\tilde{u}} \quad \tilde{u} = \frac{\tilde{p}}{\tilde{z}}$$

Intensity - time-averaged
product of \vec{p} & \vec{u}

$$\frac{1}{2} \operatorname{Re} \{ \vec{p} \vec{u}^* \}$$

vector Intensity is a vector

Decibels - always give a
reference

Levels \rightarrow decibels