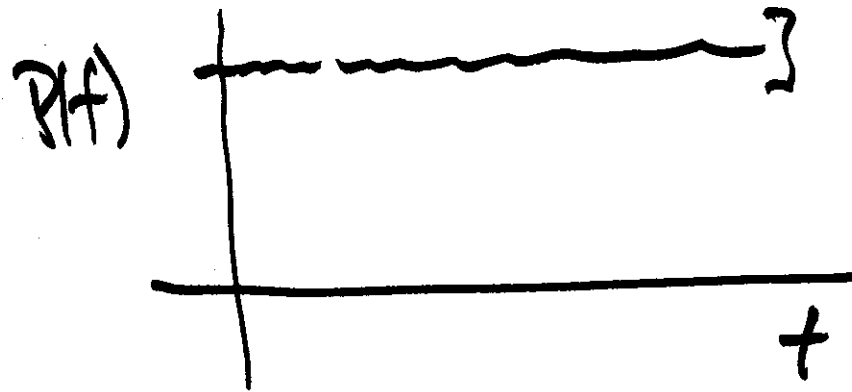


Sound Waves

- propagating pressure fluctuations in an elastic medium



$$\underline{p(t)} = P(t) - P_0$$

$P \propto u$

~~Eqn of state~~
Eqn of continuity

$$P = \beta \frac{\delta s}{\delta P_0}$$

$$\frac{\delta s}{\delta t} + \underline{\nabla \cdot \vec{u}} = 0$$

$$\nabla \cdot \bar{u}$$

$$\bar{u} = u_x \bar{i} + u_y \bar{j} + u_z \bar{k}$$

∇ grad operator

$$\nabla = \bar{i} \frac{d}{dx} + \bar{j} \frac{d}{dy} + \bar{k} \frac{d}{dz}$$

$$\nabla \cdot \bar{u} = \frac{du_x}{dx} + \frac{du_y}{dy} + \frac{du_z}{dz}$$

(ii i) pressure vs particle velocity

combine (1) $p = \beta s$ with (2)

$$\boxed{\frac{1}{\beta} \frac{dp}{dt} + \nabla \cdot \bar{u} = 0} \quad (3)$$

linearized continuity equation

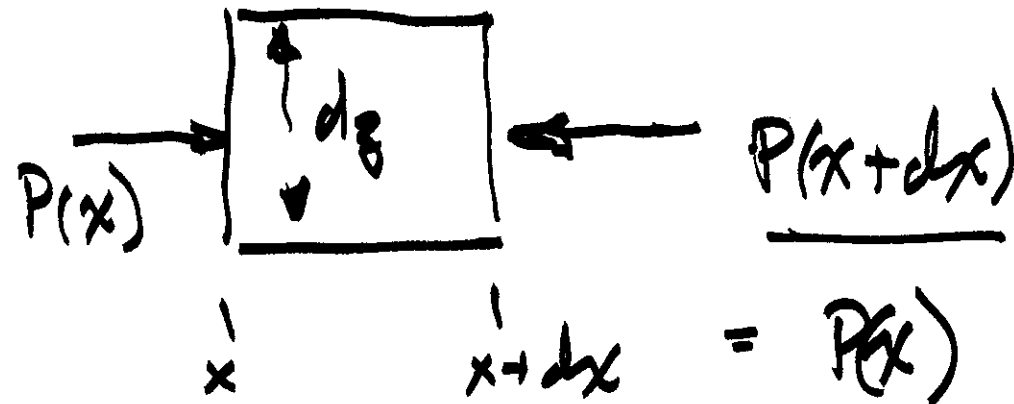
3.2.2 Pressure - Velocity Relation (II)

(i) Eqn of motion

Apply $f = ma$ to a fixed mass of fluid that is moving with the fluid



1-D



$$P(x+dx) = P(x) + \frac{dP}{dx} \Big|_x dx + \dots$$

Net force acting in the x-direction

$$df_x = P(x)dydz - P(x+dx)dydz$$

$$= - \frac{dP}{dx} \Big|_x \underbrace{dx dy dz}_{\text{volume } dV}$$

Force/unit volume

In 3-D $d\bar{f} = -\nabla\phi dV$

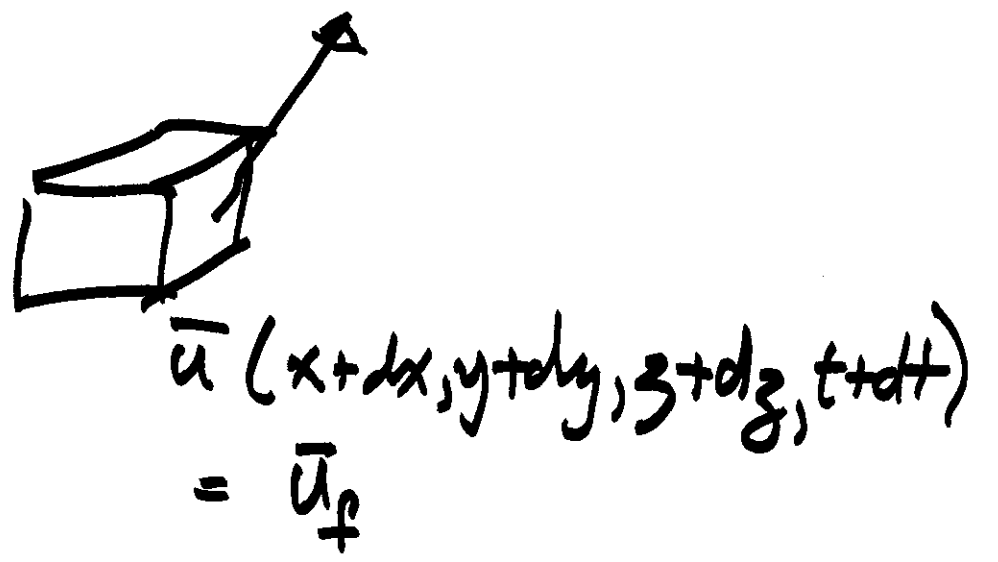
∇P = gradient of the pressure

$$= \frac{\partial P}{\partial x} \bar{i} + \frac{\partial P}{\partial y} \bar{j} + \frac{\partial P}{\partial z} \bar{k}$$

direction - steepest change
in pressure

(iii) Acceleration of a moving element of fluid

$$f = ma$$



$$\bar{u}_f = \bar{u}_i u(x, y, z, t) + \frac{d\bar{u}}{dx} \bigg|_{x, y, z, t} dx + \dots + dz$$

$$dx = u_x dt$$

$$dy = u_y dt$$

$$dz = u_z dt$$

$$+ \frac{d\bar{u}}{dt} \bigg|_{x, y, z, t} dt$$

$$\bar{a} = \frac{\bar{u}_f - \bar{u}_i}{dt} = \left[\underbrace{u_x \frac{\partial \bar{u}}{\partial x} + u_y \frac{\partial \bar{u}}{\partial y} + u_z \frac{\partial \bar{u}}{\partial z}}_{\text{convective acceleration}} + \frac{\partial \bar{u}}{\partial t} \right]$$

convective acceleration

- very small (in linear acoustics)

$$\bar{a} = \cancel{(\bar{u} \cdot \nabla) \bar{u}} + \frac{\partial \bar{u}}{\partial t}$$

neglect because terms are
non-linear

$$d\bar{f} = -\nabla P dV$$

$$\bar{a} = \frac{\partial \bar{u}}{\partial t}$$

mass of element

$$m = \rho dV$$

$$d\bar{f} = m \bar{a}$$

$$-\nabla P dV = \rho dV \frac{d\bar{u}}{dt}$$

(iv) Pressure-Velocity Relation

$$P = P_0 + p$$

$$\underline{\nabla P = \nabla p}$$

$$\rho = \rho_0 (1 + s)$$

$$s = \frac{\rho - \rho_0}{\rho_0}$$

$$s \ll 1$$

$$\boxed{-\nabla p = \rho_0 \frac{d\bar{u}}{dt}} \quad (4)$$

↑

linearized momentum equation
Euler Equation

3.2.3 Linear Wave Equation

$$\left. \begin{array}{l} (3) \quad \frac{1}{\beta} \frac{\partial^2 p}{\partial t^2} + \nabla \cdot \bar{u} = 0 \\ (4) \quad \nabla p + \rho_0 \frac{\partial \bar{u}}{\partial t} = 0 \end{array} \right\} \begin{array}{l} \frac{\partial}{\partial t} \\ \frac{1}{\rho_0} \nabla \cdot \end{array}$$

eliminate \bar{u}

~~and~~ subtract (3) from (4)

$$\nabla^2 p - \left(\frac{\rho_0}{\beta}\right) \frac{\partial^2 p}{\partial t^2} = 0$$

$$c = \sqrt{\frac{\beta}{\rho_0}}$$

$$\nabla^2 p - \frac{1}{c^2} \frac{d^2 p}{dt^2} = 0$$

$$c = \text{speed of sound}$$

$$= \sqrt{\frac{\beta}{\rho}} = \sqrt{\frac{\gamma p_0}{\rho}}$$

linearized wave equation

- governs propagation of small amplitude pressure fluctuations in a stationary, ideal fluid.

$$\angle u < \frac{1}{10} c$$

3.2.4 Sound speed

c = speed of wave propagation

$$= \sqrt{\frac{\gamma P_0}{\rho_0}} = 331.6 \text{ m/s at } 0^\circ\text{C}$$

$$\approx 340 \text{ m/s at room temp } 20^\circ\text{C}$$

in an isothermal atmosphere

$\left(\frac{P_0}{\rho_0}\right)$ is a constant

$c \neq$ function of height (if temp is constant)

c = directly proportional to
absolute temp to the $\frac{1}{2}$ power

In air $c = c_0 \sqrt{\frac{T_k}{273}}$ T_k [°K]

$= c_0 \sqrt{1 + \frac{T_c}{273}}$ T_c [°C]

3.3 One-Dimensional Solutions

- depends only on 1 space variable

- plane

- spherical

- cylindrical



spherically symmetric

cylindrically symmetric