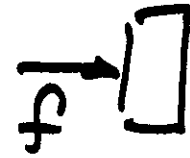
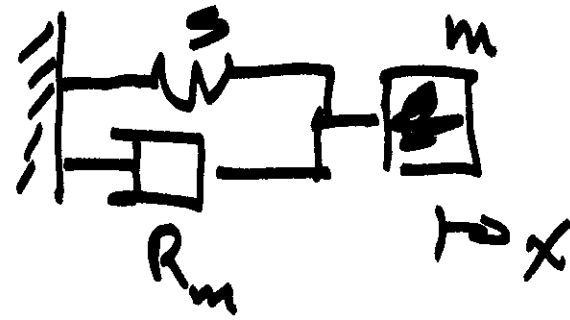


- governing equations
- wave equation
- possible solutions
- boundary solutions



Restoring force

$$f = -s x - R_m \frac{dx}{dt}$$

$$f = m \frac{d^2 x}{dt^2}$$

$$\omega_0 = \sqrt{\frac{s}{m}}$$

$$\frac{d^2 x}{dt^2} + \left(\frac{R_m}{m}\right) \frac{dx}{dt} + \omega_0^2 x = 0$$

$$\gamma = \frac{-\left(\frac{R_m}{m}\right) \pm \sqrt{\left(\frac{R_m}{m}\right)^2 - 4\omega_0^2}}{2}$$

$$\beta = \left(\frac{R_m}{2m}\right)$$

$$\gamma = -\beta \pm \sqrt{\beta^2 - \omega_0^2}$$

Usually - underdamped

$$\beta < \omega_0$$

$$\gamma = -\beta \pm j \sqrt{\omega_0^2 - \beta^2} \text{ positive}$$

$$\underline{\underline{x^2 = \tilde{A} e^{\gamma t}}}$$

$\omega_0^2 - \beta^2 = \omega_d^2$   
damped natural  
frequency.

$$\gamma = -\beta \pm j\omega_d$$

$$\tilde{A} e^{\gamma t}$$

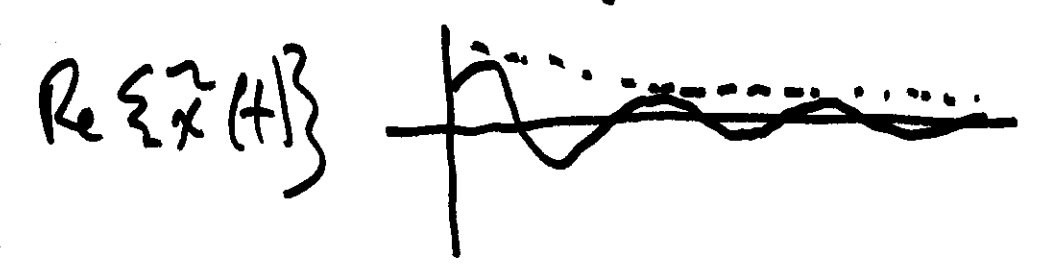
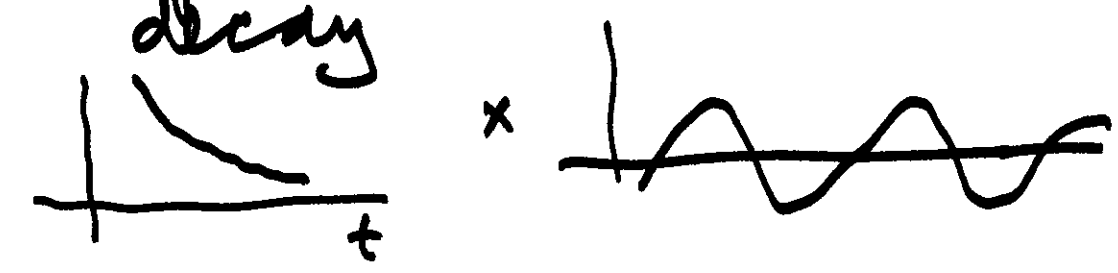
Complete solution

$$\tilde{x} = \tilde{A}_1 e^{-\beta t} e^{+j\omega_d t} + \tilde{A}_2 e^{-\beta t} e^{-j\omega_d t}$$

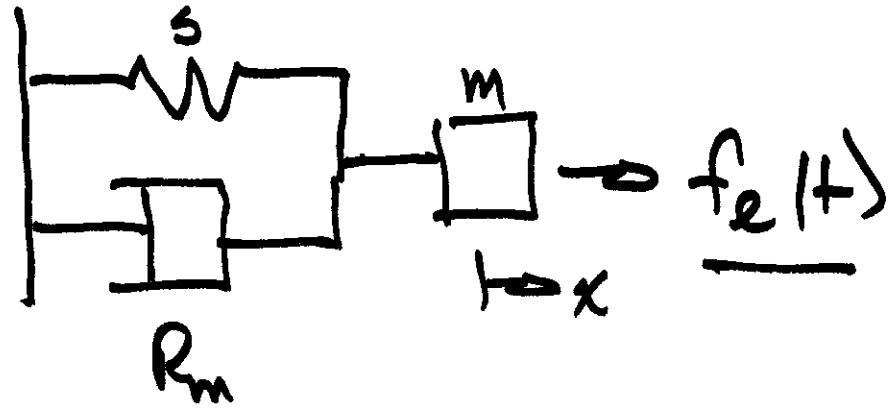
$$\beta = \frac{R_m}{2m}$$

$$= e^{-\beta t} \left[ \tilde{A}_1 e^{+j\omega_d t} + \tilde{A}_2 e^{-j\omega_d t} \right]$$

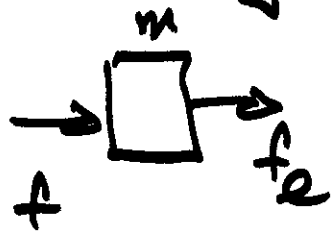
Exponential decay      oscillatory term



## 1.3 Forced Oscillator



### 1.3.1 Governing Equations



$$f = -sx - R_m \frac{dx}{dt} \quad (1)$$

ROM

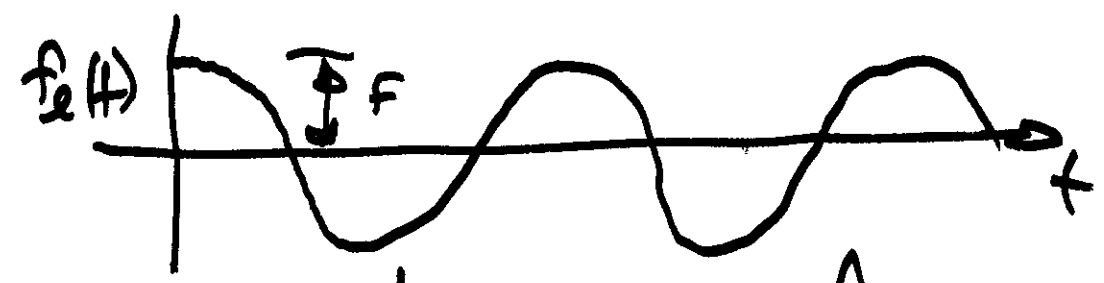
$$f + f_e = ma = m \frac{d^2x}{dt^2} \quad (2)$$

Combine (1) + (2)

$$\boxed{m \frac{d^2 x}{dt^2} + R_m \frac{dx}{dt} + s x = f_e(t)} \quad (3)$$

inhomogeneous ODE

Example  $f_e(t) = F \cos \omega t \quad t \geq 0$



Laplace Transform  
 ← soln of the homogenous eqn

$$x = x_{\text{transient}} + x_{\text{steady-state}}$$

6

Since  $R_m > 0$  always

as  $t$  increases transient sol'n decays  
and becomes negligible

$e^{-\beta t} \rightarrow 0$  when  $\beta t \Rightarrow \infty$

### 1.3.2 Steady State Solutions

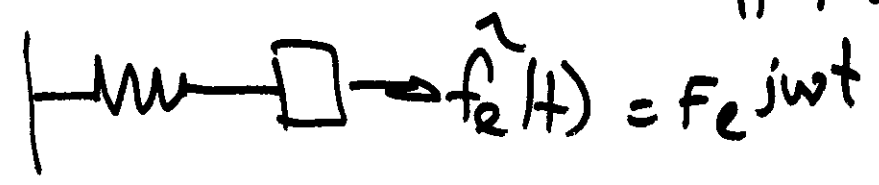
Assume complex form for the driving force

$$\hat{f}_e = F e^{j\omega t}$$

↑  
driving  
frequency

$$\text{Re} \{ \hat{f}_e(t) \} = \text{~~2 cos~~ } F \cos \omega t$$

if F is real.



Linear systems

- in steady state respond at the driving frequency

Assumed solution

$$\vec{P}_e = F e^{j\omega t}$$

$$\vec{x} = \hat{A} e^{j\omega t}$$

substitute into (3)

$$- \omega^2 m \hat{A} e^{j\omega t} + j\omega R_m \hat{A} e^{j\omega t} + s \hat{A} e^{j\omega t} = F e^{j\omega t}$$

$$\hat{A} = \frac{F}{-\omega^2 m + j\omega R_m + s}$$

$$= \frac{F}{j\omega \left( R_m + i \left( \omega m - \frac{s}{\omega} \right) \right)}$$



$$\tilde{x} = \tilde{A} e^{j\omega t}$$

$$= \frac{F e^{j\omega t}}{j\omega \left( R_m + j \left( \omega m - \frac{S}{\omega} \right) \right)}$$

Steady-state  
response

# Physical Displacement (steady-state)

$$x(t) = \text{Re} \{ \tilde{x}(t) \}$$

Velocity:  $\tilde{u} = \frac{d\tilde{x}}{dt} = j\omega \tilde{x}$

Acceleration:  $\tilde{a} = \frac{d^2\tilde{x}}{dt^2} = \frac{d\tilde{u}}{dt} = -\omega^2 \tilde{x} = j\omega \tilde{u}$

$\omega$   
↑  
driving  
freq

$\omega_0$   
↑  
natural  
freq

$$\tilde{u} = \frac{F e^{j\omega t}}{R_m + j(\omega m - \frac{s}{m})}$$

### 1.3.3 Mechanical Impedance

Define:

$$\hat{z}_m = \frac{\text{complex driving force}}{\text{complex velocity at the driven point}}$$

$$= \frac{F}{v}$$

Damped SDOF

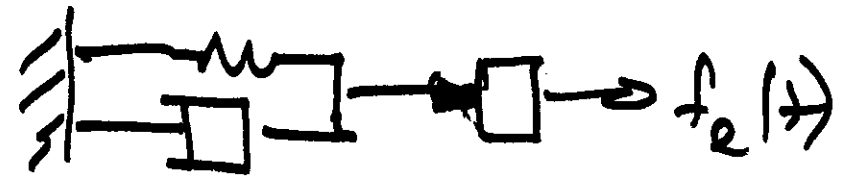
$$\hat{z}_m = \frac{F_e^{j\omega t}}{F_e^{j\omega t} / (R_m + i(\omega m - \frac{s}{\omega}))}$$

$$= \underbrace{R_m}_{\text{mechanical resistance}} + i \underbrace{(\omega m - \frac{s}{\omega})}_{\text{mechanical reactance}}$$

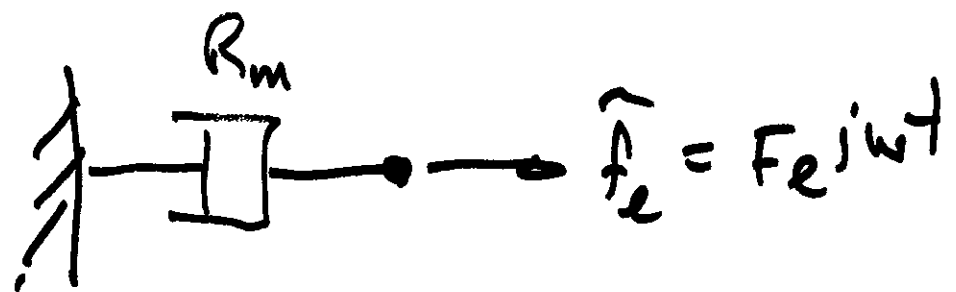
$$\hat{z}_m = R + iX$$

$$R = \text{Re} \{ \hat{z}_m \}$$

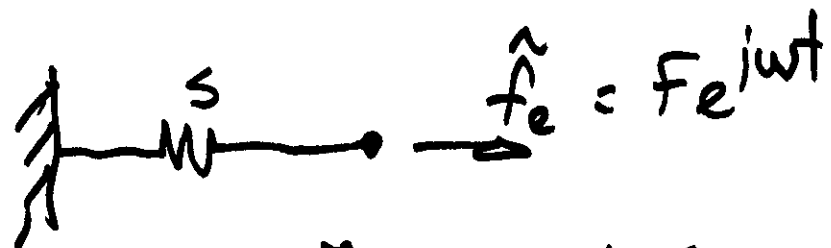
$$X = \text{Im} \{ \hat{z}_m \}$$



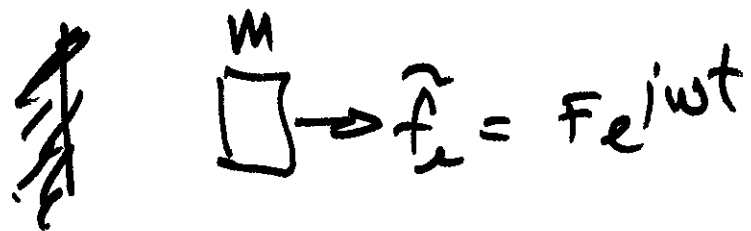
Impedance is a property of the system



$$\hat{z}_m = R_m \quad \text{pure resistance}$$



$$\hat{z}_m = -j \frac{s}{\omega} \quad \text{pure stiffness}$$



$$\hat{z}_m = j\omega m \quad \text{pure mass}$$

