# Quantum Phase Estimation: **High-Dimensional Photonic Advancements**

## What is a quantum phase estimation algorithm (PEA)?

A quantum phase estimation algorithm is a means of determining an unknown eigenphase of a unitary matrix  $U_{r}$ 

$$U|\nu_k\rangle = e^{i2\pi\theta_k}|\nu_k\rangle.$$

where  $|v_k\rangle$  is any eigenstate of U,  $e^{i2\pi\theta_k}$  is its corresponding eigenvalue and  $\theta_k$  is called the eigenphase. While U can be *applied* to quantum states, it is treated as an unknown black box<sup>[1].</sup>

### Why PEA?

Quantum phase estimation algorithms (PEAs) are is key subroutine in many quantum algorithms, including:

- Shor's factoring algorithm<sup>[2]</sup>
- Quantum principal component analysis<sup>[3]</sup>
- Generalized Grover's search algorithm<sup>[4]</sup>
- Quantum simulations<sup>[5,6]</sup>

[1] M. A. Nielsen and I. L. Chuang. Quantum Computation and Quantum Information: 10<sup>th</sup> Anniversary Edition. Cambridge

[2] P. W. Shor. "Algorithms for quantum computation: discrete logarithms and factoring". Proceedings 35th annual symposium on foundations of computer science. IEEE, 1994, pp. 124–134 [3] S. Lloyd, M. Mohseni, and P. Rebentrost. "Quantum principal component analysis". *Nature Physics* 10.9 (2014), pp. 631–633 [4] T. Byrnes, G. Forster, and L. Tessler. "Generalized Grover's Algorithm for Multiple Phase Inversion States". Phys. Rev. Lett.

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*Chemical Physics* 514 (2018), pp. 87–94 [6] A. Aspuru-Guzik, A. D. Dutoi, P. J. Love, and M. Head-Gordon. "Simulated guantum computation of molecular energies". Science 309.5741 (2005), pp. 1704–1707

# Photonic Realization of *high-dimensional* PEA for diagonal unitary systems <sup>(Lu, 2020)</sup>

**Experimental setup: frequency-controlled time-target PEA** 

**PM/IM**: Electro-optic phase/intensity modulator **PS**: Fourier-transform pulse shaper **CFBG**: Chirped fiber Bragg grating **SNSPD**: Superconducting nanowire single-photon detector **AWG**: Arbitrary waveform generator.

Both radio-frequency oscillators (18 and 27 GHz) are synchronized to the 10 MHz reference clock of the AWG.

#### A CW laser is carved into 3

frequency bins, spaced at 54GHz. Three narrow (~.2ns) spaced by 6ns are carved. The CFBG separates the frequencies within each time bin, allowing the PM to apply a different phase to each bin. This realizes a time-controlled (diagonal) unitary gate, illustrated in the lower figure. The next CFBG closes the time bin. The final PM and PS apply a probabilistic inverse (discrete) Fourier transform operation on the frequency degree of freedom.



# **Description of controlled-phase gate** Time bin ( 1 ns



(Lu, 2020) and (Moore, 2021) are collaborative work between Weiner and Kais Labs at Purdue University. Weiner Group: https://engineering.purdue.edu/~fsoptics Kais Group: https://www.chem.purdue.edu/kais/

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## **Core Concepts**

## **Traditional implementation of qubit-PEA** control register QFT $|0\rangle$ – $|0\rangle$ – $\theta_k = (0.m_1...m_{(n-2)}m_{(n-1)}m_n)_{\text{base 2}}$ register

n qubits all initialized to  $|0\rangle$  make up the control register. If the unitary of interest U is  $2^{b}$ -by- $2^{b}$  dimensional then b qubits comprise the target register which can be initialized to any state  $|\Phi\rangle$ . In the diagrammed circuit,  $|\Phi\rangle$  is initialized as eigenstate  $|\nu_k\rangle$ . Each control qubit controls whether a quantum gate will be applied to the target qubits; the gate is applied when the control is in  $|1\rangle$  and not applied when the control is  $|0\rangle$ . H represents the Hadamard gate, which converts  $|0\rangle$  to  $(|0\rangle + |1\rangle)/\sqrt{2}$ .  $QFT^{-1}$  is the inverse of the  $2^n$ -dimensional quantum Fourier transform gate. Upon measurement, each control qubit collapses to either 0 or 1. The eigenstate  $|v_k\rangle$  corresponds to eigenvalue  $e^{i2\pi\theta_k}$  and eigenphase  $\theta_k$ . If  $\theta_k \in [0, 1)$  can be represented with a base-2 number of n or fewer digits, the measured values of the control qubits  $\{m_x\}$  will represent this value:  $\theta_k = (0.m_1m_2...m_{n-1}m_n)_{\text{base 2}}.$ 

Otherwise, the measured values will represent an approximation  $\tilde{\theta}_k$  with precision  $\pm 2^n$  with high probability. The solid horizontal lines represent qubits in the circuit. Double lines represent classical information, typically generated by a measurement gate. The purple region outlines the n control qubits; the red region outlines the b target qubits.

Results



#### **Novel statistical analysis**

L	l <sub>1</sub>	
<b> 0</b> ⟩ <sub>t</sub>	1) <sub>t</sub>	$ 2\rangle_t$
.9948 ± .0004	.0101 ± .0004	.0122 ± .0005
.0023 ± .0002	.9805 ± .0009	.0120 ± .0005
.0029 ± .0002	.0094 ± .0004	.9758 ± .0010
0	$2\pi/3$	$4\pi/3$
<b>1.972</b> π	.612 <i>π</i>	<b>1.394</b> π
1.4%	2.7%	3.0%
Û	l <sub>2</sub>	
$ 0\rangle_t$	$ 1\rangle_t$	$ 2\rangle_t$
.878 ± .002	.316 ± .003	.143 ± .002
.032 ± .001	.530 ± .003	.318 ± .003
.090 ± .002	.154 ± .002	.539 ± .003
0	.3511π	<b>1.045</b> π
1.859π	.377π	<b>1.045</b> π
7.1%	1.3%	0.0%
	$ 0\rangle_{t}$ $.9948 \pm .0004$ $.0023 \pm .0002$ $.0029 \pm .0002$ $0$ $1.972\pi$ $1.4\%$ $\hat{U}$ $ 0\rangle_{t}$ $.878 \pm .002$ $.032 \pm .001$ $.090 \pm .002$ $0$ $1.859\pi$ $7.1\%$	$U_1$ $ 0\rangle_t$ $ 1\rangle_t$ .9948 ± .0004.0101 ± .0004.0023 ± .0002.9805 ± .0009.0029 ± .0002.0094 ± .00040 $2\pi/3$ 1.972 $\pi$ .612 $\pi$ 1.4%2.7% $\hat{U}_2$ $ 0\rangle_t$ $ 1\rangle_t$ .878 ± .002.316 ± .003.032 ± .001.530 ± .003.090 ± .002.154 ± .0020.3511 $\pi$ 1.859 $\pi$ .377 $\pi$ 7.1%1.3%

As many trials of the photon PEA have effectively been run, the statistics of the outputs can be used to determine a phase anywhere on  $\theta \in [0, 2\pi)$ . A visual representation of this fitting is made on the right. Results and error tabulated above.

 $2/_{3}$  $5/_{9}$  $4/_{9}$ -/3 2/9

### **Reported Works**

(Lu, 2020) Hsuan-Hao Lu, Zixuan Hu, Mohammed Saleh Alshaykh, Alexandria Moore, Yuchen Wang, Poolad Imany, Andrew Weiner, and Sabre Kais. "Quantum phase estimation with time-frequency qudits in a single photon". Advanced Quantum Technologies 3.2 (2020), p. 1900074. (Moore, 2021) Alexandria J Moore, Yuchen Wang, Zixuan Hu, Sabre Kais, and Andrew M Weiner. "Statistical approach to quantum phase estimation". New Journal of Physics 23.11 (Nov. 2021), p. 113027. doi: 10.1088/1367-2630/ac320d

Each single photon counted is equivalent to one run of the traditional PEA circuit. As both the control and target degrees of freedom are realized on a single photon, a CW laser source can be used, effectively running the PEA many times. On the left, figure (a) shows the results for a diagonal unitary with eigenphases  $(0, 2\pi/3, 4\pi/3)$ . Data (b) shows the results for a diagonal unitary with eigenphases (0,  $.3511\pi$ ,  $1.045\pi$ ). The implementation is considered successful, as the data allows each eigenphase to be determined to the nearest trit (0  $2\pi/3$ , or  $4\pi/3$ ). As we have special knowledge that the input states are eigenstates, additional analysis is possible. See below.



### Presenter: Alexandria J Moore



#### Simulated Results

The results from the quantum portion of the SPEA were simulated on both a classical computer and on IBM's Qiskit platform. The classical controller was implemented using a modified gradient search algorithm. The SPEA successfully arrives at valid eigenstate-phase pairs regardless starting conditions, as shown in the table (right). The SPEA was tested on two 3-dimensional unitaries and one 4-dimensional unitary.



# **Future Work Arbitrary unitary high-dimensional SPEA**



PM: Electro-optic phase/intensity modulator; PS: pulse shaper; IM: intensity modulator; QFP: Quantum frequency processor; C: 2-by-2 optical couplers; PC: phase controller (i.e. fiber shifter); AWG: Arbitrary waveform generator

A proposed setup for realizing a high-dimensional time-controlled frequency-target SPEA. A state preparation setup like (Lu, 2020) carves the CW laser into three frequency bins. The IM then carves the signal into four time bins prior to entering a loop. A QFP placed in a loop implements an arbitrary unitary operation on the frequency bins. The first (second, third) time-bin to enter the loop will make three (two, one) round trips prior to the last time bin entering the loop. This realizes a time-controlled unitary operation. The overlapping of the time bins realizes a pseudo-inverse Fourier transform which allows the output in time bin  $|0\rangle$  to be measured, satisfying the needs of the SPEA.

## Novel Statistical PEA (SPEA)<sup>(Moore, 2021)</sup>

Expanding on the work from (Lu, 2020), and integrating ideas from the iterative PEA (IPEA), the statistical PEA (SPEA) was proposed and studied. The SPEA utilizes a variational scheme: a classical controller controls the input state and the rotation gate of an IPEA circuit. The statistics of the measured state serve as an objective function: the prepared state and applied rotation are the eigen-state and –phase when the control qudit deterministically is measured in bin  $|0\rangle$ .

	Abs	Iteration			
y Input state	inner product	Mean	S.D.	Mean phase error	
(0.1951, 0.9808)	0.98	6.20	2.82	$1.099 \times 10^{-2}$	
(0.3827, 0.9239)	0.92	8.15	3.41	$1.005 \times 10^{-2}$	
(0.7071, 0.7071)	0.71	8.90	3.34	$1.005 \times 10^{-2}$	
(0, 0, 0.7432, 0.6690)	0.99	5.85	8.14	$2.083 imes10^{-2}$	
(0, 0, 0.6690, 0.7432)	0.99	6.7	10.42	$2.168 imes10^{-2}$	
(0, 0, 1, 0)	0.71	17.7	6.06	$1.663 \times 10^{-2}$	
(1, 0, 0, 0)	0.71	23.05	11.22	$2.167 imes10^{-2}$	
(0.7071, 0, 0.7071, 0)	0.50	21.3	10.71	$2.262  imes 10^{-2}$	
(-0.1379, 0, 0, 0.9904)	0.99	1.15	0.36	$1.885 imes10^{-2}$	
(0, 0.7807, 0.6247, 0)	0.99	1.1	0.3	$1.508 \times 10^{-2}$	
(0, 1, 0, 0)	0.71	4.35	4.17	$1.414  imes 10^{-2}$	
(0.7071, 0, 0, 0.7071)	0.62	4.15	1.01	$1.570 \times 10^{-2}$	
(0.5774, 0.5774, 0, 0.5774)	0.51	21.5	11.06	$2.199  imes 10^{-2}$	
	(0.1951, 0.9808) (0.3827,0.9239) (0.7071, 0.7071) (0, 0, 0.7432, 0.6690) (0, 0, 0.6690, 0.7432) (0, 0, 1, 0) (1, 0, 0, 0) (0.7071, 0, 0.7071, 0) (-0.1379, 0, 0, 0.9904) (0, 0.7807, 0.6247, 0) (0, 1, 0, 0) (0.7071, 0, 0, 0.7071) (0.5774, 0.5774, 0, 0.5774)	Abs. inner product           (0.1951, 0.9808)         0.98           (0.3827,0.9239)         0.92           (0.7071, 0.7071)         0.71           (0, 0, 0.7432, 0.6690)         0.99           (0, 0, 0.6690, 0.7432)         0.99           (0, 0, 1, 0)         0.71           (1, 0, 0, 0)         0.71           (0, 7071, 0, 0.7071, 0)         0.50           (-0.1379, 0, 0, 0.9904)         0.99           (0, 1, 0, 0)         0.71           (0, 7071, 0, 0.7071, 0)         0.50           (-0.1379, 0, 0, 0.9904)         0.99           (0, 1, 0, 0)         0.71           (0, 7071, 0, 0.7071)         0.62           (0, 7071, 0, 0, 0.7071)         0.62	Abs. Abs.Itera Mean $(0.1951, 0.9808)$ $0.98$ $6.20$ $(0.3827, 0.9239)$ $0.92$ $8.15$ $(0.7071, 0.7071)$ $0.71$ $8.90$ $(0, 0, 0.7432, 0.6690)$ $0.99$ $5.85$ $(0, 0, 0.6690, 0.7432)$ $0.99$ $6.7$ $(0, 0, 1, 0)$ $0.71$ $17.7$ $(1, 0, 0, 0)$ $0.71$ $23.05$ $(0.7071, 0, 0.7071, 0)$ $0.50$ $21.3$ $(-0.1379, 0, 0, 0.9904)$ $0.99$ $1.15$ $(0, 7071, 0, 0.7071)$ $0.62$ $4.15$ $(0.7071, 0, 0, 0.7071)$ $0.62$ $4.15$ $(0.5774, 0.5774, 0, 0.5774)$ $0.51$ $21.5$	Abs. (0.1951, 0.9808)Input stateIteration(0.1951, 0.9808)0.986.202.82(0.3827, 0.9239)0.928.153.41(0.7071, 0.7071)0.718.903.34(0, 0, 0.7432, 0.6690)0.995.858.14(0, 0, 0.6690, 0.7432)0.996.710.42(0, 0, 1, 0)0.7117.76.06(1, 0, 0, 0)0.7123.0511.22(0.7071, 0, 0.7071, 0)0.5021.310.71(-0.1379, 0, 0, 0.9904)0.991.150.36(0, 0, 7807, 0.6247, 0)0.991.10.3(0, 1, 0, 0)0.714.354.17(0.7071, 0, 0.7071)0.624.151.01(0.5774, 0.5774, 0, 0.5774)0.5121.511.06	

The SPEA can be iterated to perform a full spectral decomposition, recovering all eigenstate-phase pairs. This functionality was tested on the (quantum chemistry relevant) 16by-16 representation of the water molecule. The quality of the decomposition (both in overall matrix fidelity and minimum average eigenphase error) improved as higherdimensional control circuits were used. This illustrates not only the strength of the SPEA itself, but also indicates some unique advantages may be offered by high-dimensional algorithms.