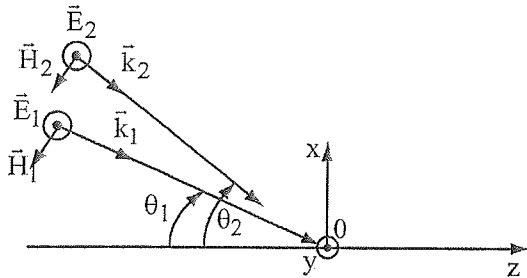


Two electromagnetic plane waves propagate through vacuum in directions defined by \vec{k}_1 and \vec{k}_2 as shown, where $\vec{k}_i = \frac{2\pi}{\lambda}(-\sin \theta_i \hat{x} + \cos \theta_i \hat{z})$ for $i = 1, 2$. For each, the electric wave can be expressed as $\vec{E}_i = \vec{E}_{i,0} e^{-j(\vec{k}_i \cdot \vec{r} - \omega_i t)}$.



A) (26 points) For $\omega_1 = \omega_2$, determine the distance Λ between the time-averaged power density maxima in the plane defined by $z = 0$. Express your result in terms of $\lambda (= 2\pi / |\vec{k}|)$, θ_1 and θ_2 . Determine Λ for $\lambda = 1.0 \mu\text{m}$, $\theta_1 = 0.01 \text{ rad}$ and $\theta_2 = -0.01 \text{ rad}$.

B) (27 points) Now let ω_1 and ω_2 be similar to, but slightly different from one another (i.e., $|\omega_1 - \omega_2| \ll \omega_1, \omega_2$), and let $\theta_1 = -\theta_2$. Determine the velocity v of the time-averaged power density peaks in the $z = 0$ plane in terms of $\lambda, \theta_1, \theta_2$, and $\Delta\omega = \omega_1 - \omega_2$. Estimate v for $\lambda = 1.0 \mu\text{m}$, $\theta_1 = -\theta_2 = 0.01 \text{ rad}$ and $\Delta\omega = 2\pi \times 10^7 \text{ rad/sec}$.

C) (27 points) Consider the two waves of frequency ω_1 and ω_2 (again with $|\omega_1 - \omega_2| \ll \omega_1, \omega_2$) as they propagate collinearly ($\theta_1 = \theta_2 = 0$) through a dispersive, non-absorbing, isotropic, non-magnetic medium. The relative permittivity of the medium is given by

$$\epsilon_r(\omega) = \epsilon_r(\omega_0) + \epsilon_r'(\omega_0)(\omega - \omega_0) + \dots, \text{ where } \epsilon_r'(\omega_0) = \left. \frac{d\epsilon_r(\omega)}{d\omega} \right|_{\omega=\omega_0} \text{ and } \omega_0 \text{ is the average}$$

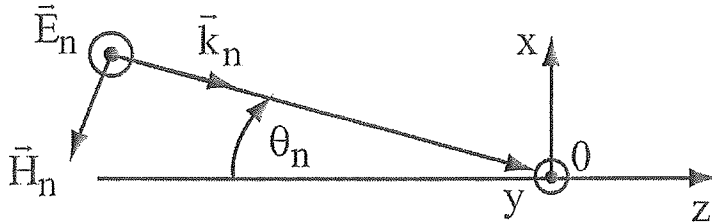
frequency given by $\omega_0 = (\omega_1 + \omega_2) / 2$. Derive an expression for the velocity of the time-averaged power density peak of this waveform, travelling in the $+z$ direction. Your answer should be in terms of $\omega_0, c (= (\epsilon_0 \mu_0)^{-1/2}), \epsilon_r(\omega_0)$ and $\epsilon_r'(\omega_0)$. Estimate this velocity for

$$\lambda_0 = 1.0 \mu\text{m} \text{ (wavelength in free space), } \epsilon_r(\omega_0) = 2, \epsilon_r'(\omega_0) = \frac{10^{-11}}{2\pi} (\text{rad/sec})^{-1} \text{ and}$$

$$\Delta\omega = 2\pi \times 10^7 \text{ rad/sec.}$$

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D) (20 points) Finally, we consider a large number N of plane waves, each propagating through vacuum (impedance $\eta \approx 377 \Omega$) at an angle $\theta_n = n\Delta\theta$ with respect to the z axis, where n is an integer between $-N/2$ and $+N/2$. Let $N\Delta\theta \ll 1$.



$$E_n = E_0 e^{-j(\vec{k}_n \cdot \vec{r} - \omega_0 t)}$$

$$\vec{k}_n = \frac{2\pi}{\lambda} (-\sin \theta_n \hat{x} + \cos \theta_n \hat{z}).$$

The amplitude E_0 and frequency ω_0 are the same for all waves.

- What is the peak time-averaged power density in the $z = 0$ plane in terms of N , E_0 and η ?
- Show that the distance from the axis to the first zero of the time-averaged power density is $x_0 = \frac{\lambda}{N\Delta\theta}$.
- Find the distance x_m from the axis to the next time-averaged power density maximum.
- What is the time-averaged power density of the peak at $x = x_m$, relative to that of the central peak?

[**Hint:** Represent the sum of these waves using a phasor diagram. The magnitude of each individual phasor is E_0 and the phase difference between phasors is δ . Express δ as a function of x . As N gets large, the phasor sum representing the total field is the arc of a circle. Draw this phasor diagram for the field at the central peak. Repeat for the field at the first zero, and the next maximum.]

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Maxwell's Equations:

$$\nabla \cdot \mathbf{D} = \rho_v$$

$$\oiint_S \mathbf{D} \cdot d\mathbf{S} = Q_{enc}$$

$$\nabla \cdot \mathbf{B} = 0$$

$$\oiint_S \mathbf{B} \cdot d\mathbf{S} = 0$$

$$\nabla \times \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t}$$

$$\oint_C \mathbf{E} \cdot d\mathbf{l} = -\frac{d}{dt} \oiint_S \mathbf{B} \cdot d\mathbf{S}$$

$$\nabla \times \mathbf{H} = \mathbf{J} + \frac{\partial \mathbf{D}}{\partial t}$$

$$\oint_C \mathbf{H} \cdot d\mathbf{l} = I_{enc} + \frac{d}{dt} \oiint_S \mathbf{D} \cdot d\mathbf{S}$$

Poynting's Theorem:

$$\nabla \cdot (\mathbf{E} \times \mathbf{H}) = -\frac{\partial}{\partial t} (\mathbf{B} \cdot \mathbf{H}) - \frac{\partial}{\partial t} (\mathbf{D} \cdot \mathbf{E}) - \mathbf{J} \cdot \mathbf{E}$$

Potentially Useful Vector Algebra

$$\nabla \cdot \mathbf{A} = \hat{x} \left(\frac{\partial A_x}{\partial x} + \frac{\partial A_y}{\partial y} + \frac{\partial A_z}{\partial z} \right) + \hat{y} \left(\frac{\partial A_x}{\partial z} - \frac{\partial A_z}{\partial x} \right) + \hat{z} \left(\frac{\partial A_y}{\partial x} - \frac{\partial A_x}{\partial y} \right)$$

$$\nabla \cdot \mathbf{A} = \frac{\partial A_x}{\partial x} + \frac{\partial A_y}{\partial y} + \frac{\partial A_z}{\partial z}$$

$$\nabla = \hat{x} \frac{\partial}{\partial x} + \hat{y} \frac{\partial}{\partial y} + \hat{z} \frac{\partial}{\partial z}$$

Potentially Useful Integral Identities

$$\int \frac{dx}{x} = \ln x$$

$$\int \frac{dx}{\sqrt{x^2 + a^2}} = \ln \left(x + \sqrt{x^2 + a^2} \right)$$

$$\int \frac{dx}{x^2 + a^2} = \frac{1}{a} \tan^{-1} \frac{x}{a}$$

$$\int \frac{dx}{(x^2 + a^2)^{3/2}} = \frac{x}{a^2 \sqrt{x^2 + a^2}}$$

$$\int \sin^2 x dx = \frac{x}{2} - \frac{1}{4} \sin 2x$$

$$\int \sin^3 x dx = \frac{\cos^3 x}{3} - \cos x$$

$$\int \sinh^2 x dx = \frac{1}{2} [-x + \sinh x \cosh x]$$

$$\int \frac{xdx}{\sqrt{x^2 + a^2}} = \sqrt{x^2 + a^2}$$

$$\int \frac{xdx}{x^2 + a^2} = \frac{1}{2} \ln(x^2 + a^2)$$

$$\int \frac{xdx}{(x^2 + a^2)^{3/2}} = -\frac{1}{\sqrt{x^2 + a^2}}$$

$$\int \cos^2 x dx = \frac{x}{2} + \frac{1}{4} \sin 2x$$

$$\int \cos^3 x dx = -\frac{\sin^3 x}{3} + \sin x$$

$$\int \cosh^2 x dx = \frac{1}{2} [x + \sinh x \cosh x]$$

Other Information

$$v_g = \frac{d\omega}{dk}, \quad k = \omega \sqrt{\mu\epsilon}$$